JOURNAL OF MECHANICAL ENGINEERING

STROJNIŠKI VESTNIK



nd level(all): 56.9 dB(A)



no. 10 year **2008** volume 54

> Sound level(all): 52.2 dB(A) Sound Presurre Level: 52.2 [dB(A)]

Strojniški vestnik - Journal of Mechanical Engineering ISSN 0039-2480

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University of Ljubljana Faculty of Mechanical Engineering Journal of Mechanical Engineering Aškerčeva 6 SI-1000 Ljubljana, Slovenia Phone: 386-(0)1-4771 137 Fax: 386-(0)1-2518 567 E-mail: info@sv-jme.eu http://www.sv-jme.eu

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Cover plate: Figures represent sound patterns of a larger industrial plant, recorded with an acoustical camera. Figures show identified main noise source (in total - above) and two noise sources at a particular frequency range (below). Visualisation of noise sources in different frequency ranges is necessary for a successful noise control.

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Littera Picta, Medvode, printed in 600 copies

Yearly subscription

companies	100,00 EUR
individuals	25,00 EUR
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abroad	100,00 EUR
single issue	5,00 EUR

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Strojniški vestnik - Journal of Mechanical Engineering is also available on http://www.sv-jme.eu, where you access also to papers' supplements, such as simulations, etc.

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Instructions for Authors

Simple Feedback Structure of Active Noise Control in a Duct

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An active noise control is usually constructed with the use of electronic filters. For sufficient noise attenuation, electronic filters are not always needed. In this case, the electronic controller and appropriate software is not required so the system can be much easier. This paper deals with the use of a feedback structure of active noise control in an experimental ventilation duct. Simulation was performed to investigate the efficiency of a simple analoguous system for active noise control without incorporating electronic filters. The transfer function of the entire analogous system can be used to predict the maximum attenuation level. Tests were made to verify the simulation and to show what noise attenuation level can be achieved in an experimental duct. It has been shown that in a specific frequency range this kind of a system is efficient enough for use in some ventilation ducts.

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Keywords: active noise control, feedback control, ANC simulation, phase shift

0 INTRODUCTION

The beginning of active noise control (ANC) is in year 1936, when Lueg patented his idea about realization of the active noise control [1]. It uses the principle of interference and absorption. Because electronics in those years was not enough advanced to meet the needs of a controller, no real system was produced.

Today, two basic methods are used for active noise control; feed-back and feed-forward. They are using all capabilities of contemporary computers and other electronics. These systems are used mainly in conjunction with adaptive filters, which make them capable to cope with bad system response. But there is a question, if it is possible (in some cases) to use a simpler system, which is efficient enough and at the same time cheaper and more reliable because of less electronic components. This is the purpose of the experimental ventilation duct that was made in this research.

The characteristic of ventilation ducts' fans is a constant rotating speed. This means that the emitting noise is constant and tonal in lower frequency band. The intention of this test is to achieve good noise attenuation with feedback structure of an ANC system without incorporating adaptive or other electronic filters. These filters are capable to compensate bad system response, and at the same time, they contribute to considerable phase shift of the original signal. This part of phase shift can be eliminated if the filter is not included in the feedback loop. It should be aware of the fact that, for example, an amplifier itself also contains filters. In this research, no additional filter was used.

1 FEEDBACK STRUCTURE OF ACTIVE NOISE CONTROL

Feedback structure was found in 1953, when feedback loop was constructed of a microphone, an amplifier and a loudspeaker [2]. Authors have named this scheme as the »sound absorber«. In this scheme, microphone is used as an error microphone, where error means deviation from the theoretical attenuation. Signal is travelling through the amplifier and the controller, where the amplitude is adjusted and phase is shifted for 180°. Modified signal goes to the loudspeaker, where the attenuation of noise occurs. Schematic representation of a feedback structure is shown in Figure 1.



Fig. 1. Schematic representation of a feedback structure of ANC

Today, a feedback structure is sometimes used as a way to avoid acoustic feedback in feed-

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forward structure. A feedback ANC configuration is also very suitable when there is no available reference signal [3].

To achieve maximum level of attenuation. signal from the microphone should be amplified as much as possible, but should not exceed the stability range. Block diagram, corresponding to a feedback control system in the complex frequency s-plane, is shown in Figure 2, where E(s) is a transfer function of acoustic path between the primary source and the microphone, F(s) is a transfer function of acoustic path between the secondary source and the microphone, M(s) a transfer function of the microphone, C(s) a transfer function of the controller, N(s) a transfer function of the electronics (except the controller) and L(s) a transfer function of the secondary source.

Typical characteristics of this method is a feedback loop, which leads signal S(s) from the secondary source back to the microphone through the acoustic path. There it is added to the reference signal R(s) and the error signal D(s) is therefore increased. Signal from the error microphone D(s) can be written as:

$$D(s) = R(s) + F(s)S(s)$$
(1)

or

$$D(s) = E(s)P(s) + F(s)M(s)C(s)N(s)L(s)D(s)$$
(2)

It leads to a simplified form of a transfer function between a microphone signal D(s) and a signal of primary source P(s):

$$\frac{D(s)}{P(s)} = \frac{E(s)}{1 - M(s)C(s)N(s)L(s)F(s)}$$
(3)

Because the error signal should be as low as possible (theoretically vanish), it is D(s) = 0.

Further derivation shows that for the best attenuation, the transfer function of the controller should have an infinite amplitude. Because of the stability problem, this is not possible. When a feedback loop is amplified too much, other elements may cause the system to become unstable. Therefore an optimum operational point should be found, which is normally near the stability limit.

2 SYSTEM STABILITY

When constructing the ANC system with a feedback loop, the stability problem should always be observed. Position of the microphone and the secondary source has strong influence on stability. In theory, it is considered that both elements are infinitely close to each other. But this is not possible to reach because of the following reasons. First, it should be taken into account that both elements can not be placed close to each other because a near-field influence would be too high to get good accuracy. Second, dimensions of a loudspeaker and a microphone are preventing to define their optimum placement. For this reason, there is always some gap between the microphone and the loudspeaker, which affects the efficiency of the whole system.

If the microphone and the loudspeaker would be so close to each other that the sound waves around a microphone would be defined solely by the movement of the loudspeaker, the microphone would get the same signal as the loudspeaker produces (no phase shift). This is not possible and the gap causes some time delay in the signal, which results in phase shift. If the wavelength of the signal is significantly greater than a gap between the microphone and the loudspeaker (intermediary length, *l*), phase shift is relatively small.



Fig. 2. Block diagram of a feedback control system of ANC

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But if the wavelength is similar to the intermediary length (this occurs at the higher frequencies), phase shift becomes relatively large and efficiency of the system drastically decreases. If the frequency is increasing (at some fixed intermediary length), the phase shift is therefore also increasing, while the efficiency of the system is decreasing. Besides, it should be also taken into account the property of the real acoustical and other components, that each of them causes some phase shift in the signal. This is more significant for loudspeakers and filters.

When (for a specific sound frequency) the intermediary length corresponds to the equation (4), the phase shift between the sound, produced by the loudspeaker, and the sound, captured by the microphone, is equal to 180° :

$$l = \frac{\lambda}{2} \tag{4}$$

This means that the work of the controller, which reverses the phase for 180°, is totally ineffective and instead of attenuation it causes amplification of the primary noise.

Real noise, which is wanted to be attenuated, consists of many narrow frequency bands. Therefore, at every single intermediary length, one narrow frequency band exists, which causes the noise at this frequency to be amplified, and not attenuated. When the amplification is increasing, the amplitude is also increasing and at the specific point it exceeds the stability limit. Unstable working conditions occur and the system is not working appropriately anymore. The longer the intermediary length, the lower the limit frequency for stable operation.

Many systems of the ANC are constructed with the use of electronic adaptive filters [4] and [5]. They are able to partly compensate non-ideal response of each component and other causes of reduced efficiency. One of the filter problems is, that it causes phase shift in the signal, which results in lower efficiency. Therefore it is interesting to find out how these filters can be avoided.

3 MEASUREMENTS

The purpose of this experiment was to construct an ANC system with a feedback loop in an experimental ventilation duct without any additional compensation filter, and to get the noise attenuation, which is good enough. The setup was constructed of plywood with thickness of 2 cm (Figure 3). For simulating a fan, the loudspeaker with generated broadband pink noise was used. Narrowband or tonal noise seems to be a better choice for simulating a fan, but the experiment was intended to examine the efficiency of the system through the wider frequency spectrum.



Fig. 3. Experimental ventilation duct scheme

A feedback loop, used for generating antinoise, was constructed of a microphone Bruel&Kjaer (BK) 4165 (M2 in Fig.3), a measuring amplifier BK 2636 and a loudspeaker Visaton W170S8 with reversed polarity (L2 in Fig. 3). For primary noise generation (pink noise) the following equipment was used: a sine random generator BK 1027, a power amplifier BK 2706 and a loudspeaker JVC CS-HX621 (L1 in Figure 3).

3.1. Frequency Response of a Loudspeaker

The loudspeaker for anti-noise generation was built in an open box. According to its casing, theoretical response of the loudspeaker was calculated by the program WinISD 0.44. From the theoretical frequency response in Figure 4 it can be evident that because of resonant frequency at around 30 Hz, the loudspeaker is suitable for the use in a frequency range above 40 Hz.



Fig. 4. Theoretical frequency response of the loudspeaker



Fig. 5. Theoretical group delay of a loudspeaker

Frequency range, appropriate for noise attenuation, is seen in Figure 5, which shows theoretical group delay of the loudspeaker. At 60 Hz and lower, a group delay is too long for the system to work properly.

3.2. Frequency Response of a System

Many components of the ANC system have influence on the effectiveness of noise attenuation and each of them has its own transfer function. The main reason for delayed and modified signal in a simple ANC system (without electronic filters) is the loudspeaker, other components (microphone, amplifier, cables, ect) contribute less.

Measurement of the input and output was performed in the open loop (Figures 6 and 7), from which the transfer function can be calculated. It is obvious that the signal is very modified at the end of the loop, after it passes all the components of the ANC system. Such a simple system is very hard to compensate, because some frequencies are filtered out by the transfer function of the system. For appropriate compensation, the amplification of these particular frequencies with zero amplitude should theoretically be infinite, which is not possible.



Fig. 6. Input signal into the system

Fig. 7. Output signal from the open loop

3.3. Noise Attenuation Measurement

Then the noise attenuation with the ANC system was measured. The microphone, which is used for measurement of the attenuation level (M1 in Figure 3), was placed near the exit of the experimental duct, 45 cm from the end. Noise spectrum with and without use of an ANC, measured with the microphone M1, is seen in Figure 8 and the attenuation level in Figure 9. It is obvious that the useful frequency range of the system is approximately between 40 and 140 Hz.



Fig. 8. Attenuated (black line) and non-attenuated (grey line) noise in the duct



Fig. 9. Attenuation level in the duct

With regard to the fact that some ventilation fans produce noise mainly in lower frequency ranges, this ANC system could be very useful in that cases, despite of its narrow working range. Better attenuation could be measured if the primary noise would be narrowband or tonal noise, which is in fact more significant for ventilation fans.

4 SIMULATION

When constructing an ANC system, it is important to know its working limits. Simulation of the ANC was performed, using Matlab 7.1, to predict the effectiveness of the system. It shows the possibilities of a particular ANC system in a sense of how much the noise can be attenuated.

For a simulation to perform, the input and output of the system with the open loop were measured (as mentioned before). These two signals were used to calculate the coefficients of the impulse response, with the help of an LMS filter (Figure 10). From the impulse response coefficients, the response of a closed loop can be calculated.



Fig. 10. Impulse response coefficients

According to Figure 11, p(m) is measured input and x(m) is output of a closed loop. Filter coefficients are marked as a_i , where index *i* means its instantaneous number.



Fig. 11. Block diagram of a feedback loop

For a closed loop, the following equation can be written:

$$x(m) = p(m) + f(m), \tag{5}$$

where

$$f(m) = K[x(m)a_1 + x(m-1)a_2 + + x(m - (L-1))a_L]$$
(6)

K is gain of a feedback loop, m goes from L to the end of the signal and L is the number of filter coefficients. Extended equation is:

$$x(m) = p(m) + K[x(m)a_1 + x(m-1)a_2 + ... + x(m-(L-1))a_L]$$
(7)

The final equation of the output (attenuated) signal can be written as:

$$x(m) = \frac{1}{1 - Ka_1} \left[p(m) + Kx(m-1)a_2 + Kx(m-2)a_3 + \dots + Kx(m-(L-1))a_L \right]$$
(8)

In this way the transfer function of the whole system was simulated by the computer program. To represent the efficiency of the system, the input signal is compared with the output signal of the closed loop, which comes from the computer algorithm.

The results of the simulation are represented in Figure 12 (attenuated and original signal) and in Figure 13 (attenuation level). As expected, the simulated noise attenuation level is higher than the real measurement, because this is the best possible (ideal) attenuation for this specific ANC system and the equipment used.



Fig. 12. Original input (grey line) and simulated attenuated output signal (black line)

But the useful frequency range is more similar to that in measurement, which confirms that the simulation is correct.



Fig. 13. Attenuation level of the simulation

When evaluating the results, some other influences should also be taken into consideration. The result is in strong correlation with the intermediary length in a feedback loop and also with the position of the microphone with regard to the loudspeaker of the secondary source. This is because sound reflections are present in the duct and the sound is travelling in different directions. The more the microphone is far from the near-field, the more the reflections and other phenomena influence on the signal, consequently lower the efficiency of the system is.

Besides, the result also depends on the place, where the attenuation is measured (microphone M1 in Figure 3). Because of the impedance mismatch between the duct and the surroundings, a part of the sound waves is reflected back to the primary source, which leads to the standing waves phenomena [6] and [7]. On the specific points of the duct at the specific frequency, nodes are formed, which mean that at that points the measured noise level would be very low.

5 CONCLUSIONS

The experimental ventilation duct with a feedback method of an ANC was constructed. The possibilities of the ANC without additional electronic or other filters were investigated, because an important part of time delay of the signal may be caused by the filter. The LMS filter and the measured input and output signal were used to do the simulation of an ANC system. It shows the maximum possible attenuation for a

specific ANC system. Then the measurement of the attenuation level was performed. The experiment showed that the additional filters can be avoided, if only a specific frequency range of noise (about 40 to 140 Hz) must be attenuated. Some ventilation fans can meet this requirement. In this case, the ANC system becomes very simple and robust, because the feedback loop is constructed of fewer components. It contains just a microphone, an amplifier and a loudspeaker. During a design process, an engineer should be aware of the fact that the efficiency of the system depends on the position of the error microphone with regard to the loudspeaker for the anti-noise generation. They should be located close to each other, but not too close to become influenced too much by the near-field.

The noise, used for measurements and simulations, was broadband pink noise. Better attenuation is expected in case of using narrowband or tonal noise, which are more similar to the real ventilation fan noise.

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Dynamic Analysis of the Load Lifting Mechanisms

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This paper deals with the problem of dynamic behaviour of load lifting mechanism (such as elevators). In the case of considerable lifting heights, high velocity devices are applied, with the purpose of shortening cycle duration and increasing the capacity. In such case, the standard procedure of dynamic analysis is not applicable. In the paper, the procedure of establishing the appropriate dynamic model and corresponding equations is proposed. It enables the analysis of the relevant influences, such as variation of the rope free length, slipping of the elastic rope over the drum or pulley and damping due to the rope internal friction.

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Keywords: elevators, pulleys, dynamic models, rope length variation, motion stability analysis

0 INTRODUCTION

In the most types of load lifting mechanisms with considerable lifting heights operating under usual conditions, the steel wirerope and the driving drum (in cranes) or the driving pulley (in elevators) are applied.

The paper deals with the influence of these components upon the dynamic behaviour of the lifting mechanism, while other influences were left out of the scope of the research. These influences have already been investigated in details, e.g. dynamic phenomena in the gear box, in [5], the dynamic characteristics of the driving electromotor, in [12], etc. The research is based on the system with many degrees of freedom, also used in [13]. The analysis of the lifting mechanism with relatively low lifting velocity is usually performed by using dynamic models based on the longitudinal vibrations of elastic, homogeneous bar of a constant length, with or without mass, corresponding to certain boundary conditions, e.g. in [10]. The stability of beam acceleration and the effect of various stiffness values have been analysed in [6] and [11].

Boundary condition at the lower rope end is defined by the load mass and at the upper rope end by the moment of inertia of all driving mechanism components and by the excitation force corresponding to the mechanical characteristic of the driving electromotor, both reduced to the shaft of the driving drum or pulley. But, in the case of higher lifting velocity, (e.g. in the case of fast passenger elevator, mine shaft winding system, harbour crane, etc.) the numerous influences upon the dynamic behaviour of the lifting mechanism are also to be taken into account, such as:

- rope length variation during load lifting and lowering,
- steel wire-rope slipping over the driving pulley or drum,
- mechanical properties of the rope (such as elasticity modulus, rope internal friction, rope design, etc.),
- influence of dynamic processes on the incoming rope side upon the dynamic processes on the outgoing rope side in the systems with driving pulley,
- mechanical characteristic of the driving motor,
- influence of friction processes between the elevator car guiding shoes and rails.

Basic influence of the rope length variation appears directly through the variation of the basic dynamic parameter – stiffness of the rope free length. There is a considerable possibility that in the case of realistic devices with higher velocities and minor energy losses due to internal friction in elastic elements, the variable stiffness of the rope free length could induce the occurrence of parametric vibrations.

1 EQUATIONS OF LOAD LIFTING SYSTEM

A load lifting system with the driving pulley is presented in Fig. 1. It represents a more complex system for dynamic analysis in comparison with the system which contains a drum as a driving element.

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Fig.1. Elevator driving mechanism

The influence of the elasticity of shafts, gears, etc. can be neglected in comparison with the steel wire-rope elasticity. This enables the establishment of a simplified dynamic model of the driving mechanism, as shown in Fig. 2. The model takes into account the influence of the rope free length variations on both the incoming and outgoing pulley sides.



Fig.2. Dynamic model

The variations in rope free length (shortening or lengthening) directly affect the rope stiffness, and therefore also the dynamic behaviour of the rope, which is of significant importance in the systems with high lifting velocities. The influence of elastic rope slipping over the driving pulley is considered through the adequate boundary conditions.

According to [9] and Fig. 2, on the basis of the equilibrium conditions for the elementary rope length and for the driving pulley, the following system of equations can be established:

$$\frac{q}{g} \cdot \frac{\partial^2 u_1(x,t)}{\partial t^2} = E \cdot A \cdot \frac{\partial^2}{\partial x^2} \cdot \left[u_1(x,t) + b \cdot \frac{\partial u_1(x,t)}{\partial t} \right] + q \cdot (1 \pm \frac{a}{g})$$
(1)

$$\frac{q}{g} \cdot \frac{\partial^2 u_2(x,t)}{\partial t^2} = E \cdot A \cdot \frac{\partial^2}{\partial x^2} \cdot \left[u_2(x,t) + b \cdot \frac{\partial u_2(x,t)}{\partial t} \right] + q \cdot \left(1 \pm \frac{a}{g}\right)$$
(2)

$$M_{\rm m} = \frac{R}{i \cdot \eta} \cdot E \cdot A \cdot \frac{\partial}{\partial x} \\ \left\{ u_1(l_1, t) - u_2(l_2, t) + \frac{\partial}{\partial t} [u_1(l_1, t) - u_2(l_2, t)] - J_r \cdot \frac{a \cdot i}{\eta} \right\}$$
(3)

with:

A

g

 u_1, u_2 - rope elastic deformations, m

E - elasticity modulus, Pa

- rope cross-section, m²

a - driving mechanism acceleration, ms⁻²

- M_m driving motor torque, Nm
- *i* gear ratio,
- η driving mechanism efficiency,
- *J_r* moment of inertia of rotating masses, reduced to the pulley shaft, kgm²
- *R* driving pulley radius, m
- q rope weight per meter, Nm⁻¹
 - gravity acceleration, ms⁻²

In the presented Equations (1) to (3), the wire-rope is considered to be a viscous-elastic body (Kelvin's model), so the internal force can be defined as:

$$S(x,t) = E \cdot A \cdot \frac{\partial}{\partial x} \left[u(x,t) + b \cdot \frac{\partial u(x,t)}{\partial t} \right]$$
(4)

According to [2], dependence of the damping coefficient upon the rope stress can be adopted as:

$$b = \left(0.5 + \frac{2300}{350 + \sigma}\right) \cdot 10^{-4}, \, \text{s}$$

with:

 σ - rope stress, MPa

In the course of solving the Equation system (1) to (3), it is necessary to define the electro-mechanical model of the driving motor, or its mechanical characteristic in the form $M_{\rm m}=f(v)$.



Fig. 3. Diagram of velocity and acceleration for the elevator with double-speed motor

For the passenger elevators, the parameters defining the "driving comfort" are specified, limiting the values of maximum acceleration to: $a = 1.4 \text{ ms}^{-2}$ and its time derivate to: $(da/dt)_{\text{max}} = 1.0 \text{ ms}^{-3}$, [7]. This makes it possible to simplify the analysis by excluding the Equation (3) from the previous system. Diagram of velocity and acceleration for the elevator with double-speed motor is presented in Fig. 3.

When considering the whole driving, the analysis becomes very complex. Therefore, the paper presents only the analysis of dynamic behaviour concerning the rope incoming side of the driving pulley. In this case, the further analysis is to be carried out only with respect to the Equation (1).

2 BOUNDARY CONDITIONS FOR THE ROPE INCOMING SIDE OF THE DRIVING PULLEY

In Fig. 4, the characteristic parameters satisfying the boundary conditions for the rope

incoming side on the driving pulley or drum are presented.

Boundary conditions at point C, where the rope makes the first contact with the pulley, are dependent on whether the winding is with or without rope slipping.

The change of force in the wound rope part can be maintained only if this change is smaller than the adhesive force making it possible. Fig. 4.d shows different cases of force distribution over the wound rope length as a function of elevator velocity, [8].



Fig. 4. Boundary conditions for:
a) the pulley without slipping,
b) the pulley with slipping,
c) the elevator car (or counterweight),
d) the distribution of forces and slipping over the wound rope length v₁>v₂>v₃

The limiting value of the rope winding velocity is determined from the condition:

$$\frac{\partial S(l_1,t)}{\partial x} \le \frac{\partial S_{gr}}{\partial x}$$

with:

$$S(l_1,t) = S_{st} \cdot \left(1 + k \cdot \sin \frac{2 \cdot \pi}{v \cdot T} \cdot x\right)$$

- force in the wound rope part, no slipping,

$$S_{\rm gr} = S_{\rm st} \cdot e^{\frac{\mu}{R} \cdot (x - l_1)}$$

- force limit value on the basis of friction,

On the basis of the previous expressions, the limit slipping velocity is:

$$v_s = \frac{R \cdot \omega}{\mu} \cdot k \tag{5}$$

with:

 ω - vibration frequency of rope incoming side, s⁻¹ k - dynamic factor for elevators, k \approx 0.1 to 0.3 μ - friction coefficient between rope and pulley, $S_{\rm st}$ - rope static load at the incoming point C, N

T - period of vibration cycle, s⁻¹

If the elevator speed v exceeds the limit value v_s (the case v_1), Fig. 4.d, there occurs no rope slipping in the point C, so the force distribution over wound rope part remains the same. If the winding velocity is lower than v_{gr} , the force distribution over wound rope section has the form according to the dashed line v_2 . For very low lifting velocities, the force over the wound rope part is practically equal to the static load (line v_3).

Average values for rope slipping displacement and velocity on the pulley incoming side are as follows:

$$\Delta \bar{l}_1 \cong \frac{2 \cdot R \cdot k}{\mu \cdot \pi} \tag{6}$$

$$\Delta \bar{l}_{1} \cong \frac{2 \cdot R \cdot k \cdot \omega}{\mu \cdot \pi} \tag{7}$$

Boundary condition for rope incoming point on the pulley without slipping, Fig. 4.a (point C), is given in the form:

$$u_1(l_1,t) = \int_0^t \frac{\partial u_1(l_1,t)}{\partial x} \left(\frac{dl_1}{dt}\right) \cdot dt \tag{8}$$

with:

 l_1 - wound rope length, m

 $dl_1/dt = v$ - winding velocity, ms⁻¹

If the rope slipping is taken into account too, Fig. 4.b, the boundary condition is given in the form:

$$u_{1}(l_{1},t) = \int_{0}^{t} \frac{\partial u_{1}(l_{1},t)}{\partial x} \left(\frac{dl_{1}}{dt} - \frac{d}{dt} (\Delta l_{1}) \right) \cdot dt + \Delta l_{1} \cdot \frac{\partial u_{1}(l_{1},t)}{\partial x}$$
(9)

Boundary condition for the connection point of the rope with the elevator car or counterweight, when the friction forces between the sliding shoes and guiding rails are neglected, Fig. 4.c, is given as:

$$Q = E \cdot A \cdot \frac{\partial}{\partial x} \left(u_1(L_1, t) + b \cdot \frac{\partial u_1(L_1, t)}{\partial t} \right) + \left(\frac{Q}{g} \cdot \frac{\partial^2 u_1(L_1, t)}{\partial t^2} - a \right)$$
(10)

with:

Q - weight of the car with load passengers, N

The non-integral boundary condition (8) or (9) makes it impossible to solve partial differential Equation (1). Hence, the solution can be sought through by establishing the integral equations, which comprise both the differential equations and the appropriate boundary conditions.

3 THE INTEGRAL EQUATIONS AND ESTIMATION OF CRITICAL VELOCITY

A weightless string loaded at point M with the force F_i is shown in Fig. 5. The magnitude of rope point displacement (without slipping over the driving pulley) within the range y < x shows linear growth from zero to the boundary value u(y), while the displacement beneath the point M is of a steady value and equals to the boundary value u(y).



Fig.5. Deformation of the rope

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According to this, for an elementary force ($F_i=1$), deformation can be defined, e.g. [3], as:

$$K(x, y) = \begin{cases} \frac{x-l}{E \cdot A} & \text{for } y > x \\ \frac{y-l}{E \cdot A} & \text{for } y > x \end{cases}$$

and displacement as $u(x) = K(x, y, l) \cdot F_i$.

In case of several forces acting in points at $x = y_i$, the displacement has the following form:

$$u(x) = \sum_{i=1}^{n} K(x, y_i, l) \cdot F_i$$

For the load case with evenly distributed rope load mass, the expression takes the form:

$$u(x) = \int_{l}^{L} K(x, y, l) \cdot q \cdot dy$$

After applying this procedure to the differential Equation (1), replacing floating argument with (*y*), multiplying it with the function $K(x, y, l_1) = K$ and performing necessary mathematical transformations, e.g. [1] and [4], the rope deformation becomes as follows:

$$u_{1}(x,t) = -\int_{l_{1}}^{L_{1}} K \cdot q(y) \cdot \left[\frac{\partial^{2} u_{1}}{\partial t^{2}} - g - a \right] \cdot dy +$$

$$+ \int_{0}^{t} \frac{\partial u_{1}(l_{1},t)}{\partial x} \left[\left(\frac{dl_{1}}{dt} \right) - \frac{d}{dt} (\Delta l_{1}) \right] \cdot dt +$$

$$+ \Delta l \cdot \frac{\partial u_{1}(l_{1},t)}{\partial x} + E \cdot A \int_{l_{1}}^{L_{1}} K \cdot \frac{\partial^{3} u_{1}}{\partial s^{2} \cdot \partial t} \cdot ds -$$

$$- b \cdot E \cdot A \cdot \frac{\partial^{2} u_{1}(L_{1},t)}{\partial x \cdot \partial t} \cdot K(x,L_{1},l_{1})$$

$$(11)$$

Applying the method of particular integrals, e.g. [10], given as:

$$U_1(x,t) = X_1(x) \cdot T_1(t)$$

and considering only the first mode of vibrations, form of which - in the case of a rope with the weight at its lower end - can be adopted as the straight line, i.e. $X_1(x) = x - l_1$, the simple differential equation in the following form is obtained:

$$\ddot{T}_1 + \omega_1^2 \cdot T_1 = \varepsilon \cdot R(l_1, \nu) \cdot \dot{T}_1$$
(12)

with:

$$\omega_1^2(l_1) = \frac{g \cdot E \cdot A}{\left[Q + \frac{q \cdot (L_1 - l_1)}{3}\right] \cdot (L_1 - l_1)}$$

rope fundamental frequency, s⁻¹

$$R(l_1, v) = \frac{1}{L_1 - l_1} \cdot \left\{ v - b \cdot \frac{E \cdot A \cdot g}{\varepsilon \cdot \left[Q + \frac{q \cdot (L_1 - l_1)}{3}\right]} \right\}$$

$$\varepsilon = \frac{v}{L_1 \cdot \omega_1 \left(l_1 \right)} \,.$$

Differential Equation (12) describes nonlinear vibrations with viscous friction when the parameters are "slow-time functions", with the solution, according to [10], in the form:

$$T_1(t) = a_1(t) \cdot \cos\theta(t)$$

whereby, after some mathematical transformations, the obtained values are:

$$a_{1}(t) = h_{10} \cdot \left(\frac{L_{1} - I_{10}}{L_{1} - I_{1}}\right)^{0.25} \cdot e^{m}$$
$$\theta_{1}(t) = \int_{0}^{t} \omega_{1}(t) \cdot dt + \theta_{10}$$

with:

$$h_{10}^0 \cong -\frac{Q}{E \cdot A} \cdot \frac{a}{g}$$

amplitude value of the function, l_{10} - length of wound rope at t=0, m

$$m = \left[\frac{1}{2} \cdot \frac{d}{dt} \cdot (\Delta l_1) \cdot \int_0^t \frac{dt}{L_1 - l_1} - \frac{b}{2} \cdot \int_0^t \omega_1^2(l) \cdot dt\right]$$

After substituting previously mentioned values into Equation (9), and after performing mathematical transformations, the equation obtaines the form:

$$u_{1}(x,t) = (x-l_{1}) \cdot a_{10} \cdot \left(\frac{L_{1}-l_{10}}{L_{1}-l_{1}}\right)^{0.25} \cdot e^{m} \cdot \cos\left[\int_{0}^{t} \omega_{1}(l)dt\right] + \left[Q + \frac{q \cdot (2 \cdot L_{1} + x + l_{1})}{2}\right] \cdot \left(1 + \frac{a}{g}\right) \cdot \frac{x-l_{1}}{E \cdot A} + u_{1}(l_{1},t)$$
(13)

On the basis of Equations (4) and (13), the force in the rope is expressed as:

$$S_{1}(x,t) = E \cdot A \cdot a_{1}(t) \cdot \cos \theta_{1}(t) + \left[Q + q \cdot (L_{1} - x) \right] \cdot \left(1 + \frac{a}{g} \right)$$
(14)

Analyzing Equation (14), it can be concluded that, during the rope free length decreasing phase (load lifting), under certain conditions, its deformation can be increased, causing a permanent growth of rope dynamic load. Such phenomenon, usually described as unstable lifting, appears in the case that $\frac{da_1(t)}{dt} \ge 0$. On the basis of this condition, the critical lifting can be determined as

critical lifting velocity can be determined as:

$$v_c = \begin{cases} v_c^o & v \ge v_S \\ v_c^o + v_S & v < v_S \end{cases}$$
(15)

The critical lifting velocity in the case without rope slipping over the pulley at the incoming point C is:

$$v_c^o = \frac{2 \cdot b \cdot g \cdot E \cdot A}{Q + \frac{q \cdot (L_l - l_l)}{3}} \tag{16}$$

with:

$$v_c^o \approx 2 \cdot b \cdot \frac{q \cdot E \cdot A}{Q} \tag{17}$$

for:

$$Q >> q \cdot (L_1 - l_1)$$

Limit slipping velocity, according to (5), is:

$$v_{S}(l_{1}) = \frac{k \cdot R}{\mu} \cdot \sqrt{\frac{v_{c}^{o}}{2 \cdot b \cdot (L_{1} - l_{1})}}$$
(18)

i.e.:

$$\nu_{S}(l_{1}) \approx \frac{k \cdot R}{\mu} \cdot \sqrt{\frac{q \cdot E \cdot A}{Q \cdot (L_{1} - l_{1})}}$$
(19)

If the elevator velocity exceeds v_c in the lifting phase, the dynamic load will increase (unstable movement). Under these conditions two cases are possible:

- If the rope slipping at the incoming point is prevented (pulleys with special clamps for rope gripping), the unprevented increase of the dynamic load appears in the lifting phase,
- At standard pulleys, the increased dynamic load also appears in the lifting phase, but only until the moment when the slipping begins in the rope incoming point, $v < v_S$. During the rest of the lifting process, the load decreases, i.e. the stable movement occurs.

The value of critical velocity depends on the rope damping characteristics (due to internal friction), slipping conditions in the rope incoming point (C) and the basic characteristics of the elevator. When elevators with high load capacity, high velocity and lifting height (fast passenger elevators and mine elevators) are concerned, the lifting velocity can exceed the critical velocity, making it necessary to check the stability of the lifting process already in the design phase.

4 CONCLUSION

For driving mechanisms used for vertical load lifting (elevators and cranes), it is necessary to provide adequate conditions for the correct dynamic analysis of their parts' behavior and especially for the rope as the basic element, already in the design phase. Due to a significant influence of rope free length changes, due to its slipping over the driving pulley and for the reason of its mechanical characteristics, it is impossible to apply a classical dynamic model of longitudinal oscillations for homogenous stick of a constant length, especially by the elevators with large velocities and large lifting heights (express and mine elevators). The introduced dynamic model for elevator driving mechanism provides the analysis of relevant parameters and a base for the assessment of lifting process stability through the critical velocity level, defined with the help of expressions (15) to (19).

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Impact of the Guard Grill on the Integral and Local Characteristics of an Axial Fan

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This paper presents the impact of the guard grill (GG) on the integral characteristic (IC) and local aerodynamic properties of the outflow of an axial fan (AF). The study of the GG's impact was conducted using measurements and numeric modelling at the integral and local levels. Measurements of IC were performed according to the ISO 5801 standard and the local velocity field was measured by using a five-hole probe. Numeric simulations of the machine's local characteristics were made in the same working conditions of the fan. The comparison of results was used in the analysis of the GG's impact on the fan's efficiency and in the analysis of the feasibility of the CFD modelling of flow conditions of the GG.

© 2008 Journal of Mechanical Engineering. All rights reserved. Keywords: axial flow fans, guard grids, integral characteristic, velocity fields

0 INTRODUCTION

In industrial practice the development of an AF is mostly performed on the assumption of the decisive impact of the rotor and the inlet of the fan on operating conditions, while the impact of the GG and the supporting construction is empirically evaluated or ignored. No research has yet been performed on the GG's actual impact and the supporting construction of the AF on the precise determining of the fan's operating characteristics and outflow kinematics.

Eck [2] and [3] presented the analytical procedures of designing axial and centrifugal fans, while Wallis [1] discuss the analytical procedures for designing AF together with inlet and outlet channels. The programming and analytical design of a fan based on a similar basis as described in the literature [1] to [3] was given by Zhou Dugao et al. where the results of the fan's optimisation were also presented [17]. A common feature of the presented literature [1] to [3] and [17] is that it fails to take into account the impact of the GG on operating characteristics. Several papers have also been written on experimental and numerical studies of flow conditions in the AF. The paper by Yang Li et al. [4], studying the impact of the direction of the blade curvature of similar fans on IC and by fivehole probe and hot wire anemometry the impact on the local vector velocity field, represents the basis for the experimental procedures used in the presented paper. Studies of local kinematics in the flow field of the AF are also discussed in papers [5] with hot wire anemometry [6], with laser Doppler anemometry [9] and [10], computer visualisation and [7] and [8] numeric modelling which reveals the extensive research in this field.

The significance of the GG's impact on the IC of the fan is also evident from the number of patents related to its design [12] to [14]. The question arises of the characteristics of the impact of the GG on the IC of the AF and on the flow kinematics in the flow tract of the fan. Further on, the paper focuses on the GG's impact through a numeric and experimental approach and presents the feasibility of evaluating the impact of the GG on the IC of the fan.

0.1 Description of the Fan with and without the Guard Grill

The analysis of the GG's impact on the fan's characteristics was performed on the AF with seven unprofiled blades curved forward. The basic geometrical characteristics of the fan were given by: diameter of the outlet $D_{\text{max}} = 510 \text{ mm}$ and the dimensionless hub to tip $R_{\text{pe}}/R_{\text{ob}} = 0.28$. The fan with the GG is presented in Fig. 1 with characteristic ratia adjusted to the maximum radius of the fan's rotor: distance of the GG from the outlet area to the outlet radius $s/R_{\text{max}} = 0.235$, with the start of the rectangular area of the GG to the axis of rotation at $s/R_{\text{max}} = 0.898$ and the end

of the rectangular area of the GG to the axis of rotation at $s/R_{max} = 0.592$. A fan with no GG is fixed on the outlet via the supporting construction with four legs as presented in Fig. 2. The deviation of four supporting consoles from the level of fixing the rotor is significantly bigger than the distance of the GG which permits an assessment of the hypothesis that the supporting consoles in Figure 2 do not affect the flow conditions at the outlet from the rotor.



Fig. 1. Fan with the guard grill



Fig. 2. Fixing of the fan to the outlet without guard grill

An experimental comparative study was performed on the presented versions of the AF with regard to the GG's impact on the IC and an experimental and numeric study with regard to the GG's impact on the velocity field at the fan's outlet in the same marginal working conditions.

0.2 Measuring the Integral Characteristics

Integral measurements of the fan with and without the GG were performed at the fan measurement station at the Hidria Institute Klima according to standard ISO 5801, Fig. 3.

The operating characteristics of both fans given as the dependence of pressure difference from the volume airflow on the fan were determined in the entire working range. The volume flow of fans q_V was measured on measurement nozzles with a sensor of differential pressure and static pressure in front of the nozzle. The Mensor 6100 sensor was used in both cases. The static pressure difference Δp at the fan was measured with the same sensor. The Vaisala HMT 330 measuring device for temperature and relative humidity was used to determine the ambient conditions and pertaining air density, while the barometric pressure was measured with the Vaisala PTB 220 measuring device. The connecting electrical power P_{el} and the rotating frequency of the fan's rotor were measured in addition to the aerodynamic parameters. Electrical power was measured with the Zimmer LGM450 digital power analyser and the rotating frequency n was measured with an HIK IJ measuring device. The efficiency of the fan is given by the Equation (1) which, together with the aerodynamic characteristic $\Delta p(q_v, P_{el})$, represents the IC of the fan:

$$\eta = \frac{\Delta p \, q_{\rm V}}{P_{\rm el}} \,, \tag{1}$$

The measurement uncertainty for the volume flow is $\pm 2\%$, for differential pressure $\pm 1.4\%$ and for electrical power 1% of the measurement range. As the measurement procedures and measuring equipment used did not change for either of the two tested fans, we can assume that the measurement uncertainty was the same in both cases.

0.3 Description of Measuring the Local Velocity

In order to study the GG's impact on the flow conditions, we later conducted measurements of the outflow from the fan. Measurements of local velocity were performed using a five-hole probe which is primarily suitable for determining the intensity and direction of a velocity field; however, it is unsuitable for a reliable assessment of the fluctuation characteristics of the velocity field which was not expected in the analysis.



Fig. 3. Test station for fans made occording to standard ISO 5801 and its section view [11] 1 – axuliry fan, 2 – flow straightener, 3- seattling mesh, 4-tested fan

Assumptions for the measurement were that the fluid flow is steady for each set working point and that it is symmetrical to the axis of the rotor's rotation.

To meet the condition of axial symmetry, the experiment was conducted in a coaxial channel allowing an unhindered entry and exit velocity field with no swirl flow structures. We provided the on-flow length of $3xR_{max}$, where R is the radius of the entry cylindrical channel, and the free exit of the airflow of the fan in the semi-area with a radius of $4xR_{max}$, in line with the requirements of the ISO 5801 standard. The outline of the station is presented in Figure 4.

Velocity values and the pertaining velocity components were calculated by using ratia of the measured pressure on the five-hole probe, the measured angle α and calibration algorithms of the probe. For this paper, the algorithms were provided by the manufacturer of the five-hole probe, United Sensor. Figure 5 presents the coordinate system selected for the experiment and enabling the appropriate presentation of the results.

The differences in static pressure on the individual pressure connections of the probe were acquired simultaneously. PVC tubes of a 5 mm diameter were used to guide the pressure signals to the pressure sensors of the Endres + Hauser PD 235 type. The sampling time interval was identical for all measured points at 60s. The sampling frequency by channel was 100 Hz with 6000 samples per channel. Measuring of each working point was repeated until a dissipation of measured values has fallen under 1.5% of an average measured pressure. The ambient conditions for the measurements determined by the temperature of the environment, barometric pressure and air humidity were measured before each measurement point.



Fig. 4. Outline of the station for measuring local characteristics with a five-hole probe





The measuring of the ambient conditions was undertaken in an identical manner and with identical equipment as measurements of the IC.

The overall measurement uncertainty in measuring the velocity by the five-hole probe [15] and [16] depended on the measurement uncertainty of the differential pressure and on average equalled $\pm 3\%$ of the measured value at velocities above 10 m/s [16]. The relative error was increasing with the reduction of velocity, as expected. Given the fact that areas with a smaller velocity were less important for the analysis and that the trend in changes of the velocity field is maintained, we can ignore the increased measurement uncertainty in that area.

The uncertainty of setting the five-hole probe is limited by the uncertainty of the positioning table and setting of the initial point ± 1.5 mm.

Measurements of local characteristics were made in three 'stable' working points of the fan on the outlet side of the fan (points A, B and C in the diagram in Figure 8). They were selected on the basis of a preliminary analysis of the integral measurements of the properties of both fans. As the working points at the selected number of revolutions of the fan's rotor were given in couples (q_V , Δp), the selection of the working point



Fig. 6. Area of outlet velocity measurement with measuring point position

for local measurements was already determined by the measurement of the pressure difference on the fan to simplify the experiment. Measurements of the local characteristics of the fan with and without the GG were made 7 mm behind the fan's orifice, Figure 6.

Ten measurements were made within the radii ratios r/R_{max} from 0.412 to 0.941 with the step $\Delta r/R_{\text{max}}$ 0.059.

0.4 Numeric Model Description

The numeric analysis included a definition of the geometric model, setting up the topological model, Figure 7 and a calculation of the flow field by using the SST (Shear Stress Transport) turbulent model of the software CFX 10.0 with identical marginal conditions as in the measurements.

SST turbulence model used Reynolds averaged Navier-Stokes equations for a steady conditions:

- mass conservation equation, Eq. 2,
- momentum conservation equation, Eq. 3
- along with equations of k-ω SST turbulence model, Eq. 4 and 5 are forming closed system of Eqs [8]:

$$\frac{\partial \left(\overline{u}_{j}\right)}{x_{j}} = 0 \tag{2}$$

$$\frac{\partial \left(\rho \overline{u}_{j} \overline{u}_{i}\right)}{\partial x_{j}} = -\frac{\partial \overline{p}}{\partial x_{i}} + \frac{\partial}{\partial x_{j}} \left[\mu \left(\frac{\partial \overline{u}_{i}}{\partial x_{j}} + \frac{\partial \overline{u}_{j}}{\partial x_{i}} \right) - \rho \overline{u_{i}' u_{j}'} \right]$$
(3)

$$\frac{\partial \left(\rho k \overline{u}_{j}\right)}{\partial x_{j}} = \frac{\partial}{\partial x_{j}} \left[\Gamma_{k} \frac{\partial k}{\partial x_{j}} \right] + \tilde{G}_{k} - Y_{k}$$

$$\tag{4}$$

$$\frac{\partial \left(\rho \omega \overline{u}_{i}\right)}{\partial x_{j}} = \frac{\partial}{\partial x_{j}} \left[\Gamma_{\omega} \frac{\partial \omega}{\partial x_{j}} \right] + G_{\omega} - Y_{\omega} + D_{\omega}$$
(5)

 G_k is the turbulence kinetic energy source, because of velocity gradients, G_{ω} is the source of ω . Γ_k and Γ_{ω} are the effective diffusivity of k and ω . Yk and Yk and Y ω represent dissipation of k and ω , because of turbulence. D ω is the crossdiffusion term.

The topological model is based on identical assumptions as in the experiment in local measurements of the flow field: the fluid flow is steady for each calculation point and the fluid flow is axially symmetrical with regard to the axis rotation. Air density was 1.2 kg/m^3 and the dynamic viscosity was 1.79×10^{-5} Pas.

We used those assumptions to establish a topological model with 1/7 segment of the geometric model [8] with a hexagonal grid for the fan with and without the GG, Figure 7, with the specifications presented in Table 1. The assumption in setting the topological grid on the fan with the GG was that the concentric protective wire covering roughly 35% of the available outflow area of the outlet has a much greater impact than the radial supporters of the fan, covering roughly 3% of the available outflow area of the outlet. An identical assumption was made for the fan without the GG, where the radial supporters were further distanced from the outflow outlet by the distance $s_{\rm no}/R_{\rm max} = 0.76$, which is characteristically more than in the case of a GG.

Marginal conditions of the numeric simulation were determined on the basis of the preliminary measurement of the fan's local velocities with and without the GG.

The convergence criterion was set with regard to the development of pressure and mass flow. The convergence criterion was met with a decrease in the average remains of each quantity in the entire calculation domain by 1×10^{-6} , which required at least 350 iteration steps.

1 MEASUREMENTS AND NUMERIC SIMULATION RESULTS

Below is a presentation of the results of measuring the integral and local characteristics of the fan and results of the numeric model for the fan with and without the GG. A comparison of the results of the experiment and the model is possible on those diagrams where integral and local characteristics were given separately. As regards IC, the analysis of the GG's impact on the fan's characteristics is primarily directed at assessing the GG's impact on the operating characteristic $\Delta p(q_V)$ and efficiency η while at the local level the analysis is directed at the GG's impact on the velocity distribution at the exit of the fan's airflow.

1.1 Integral Characteristic of the Fan

The diagrams in Figs. 8 and 9 present the measured IC of the fan. Values were set to constant a rotational speed of n = 1392 rpm and the air density of $\rho = 1.2$ kg/m³ so that the fans' characteristics with and without the GG are comparable. In identical set conditions the diagram also gives the modelled values (MV) for both version of the fans.



Fig. 7. Overall and detailed 3D calculation model of fan with the guard grill

Table 1.	Topological	data of the	calculation a	area of fan with	and without	the guard gril.
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	Fan with the guard grill	Fan without guard grill
Number of nodes	832512	527760



Fig. 8. Comparison of the integral characteristic differential pressure–flow of fan with and without the guard grill



Fig. 9. Comparison of the integral characteristic electrical power–flow of fan with and without the guard grill

The abscissa presents the volume airflow q_V , and the ordinates present the static pressure difference Δp and the efficiency of the fan which is set for the maximum efficiency of the machine.

The functional dependence of the static pressure difference from the flow $\Delta p(q_{\rm V})$ (the thin full curve) presented in Figure 8 and the dependence of efficiency from the flow $\eta(q_{\rm V})$ (the thin dotted curve) enable an assessment of the GG's impact. Figure 5 presents the fan's IC consumed power-flow $P_{\rm el}(q_{\rm V})$ (the dotted curve). Both figures present the characteristics of the fan with the GG marked with full squares and characteristics of the fan without the GG marked with empty squares. The modelled values of the IC of the fans $\Delta p(q_V)$ were given by three points for each version on the fan in Figure 8. The mark (-) represents the points for the version with the GG while the mark (+) represents points for the version without the GG.

In Fig. 8 we see that the fan with the GG produces a smaller flow at the same static pressure difference because the net represents a resistance to airflow. From the maximum flow with the difference of 4.2%, the difference of flows is decreasing with the flow reduction as the flow pressure loss on the protective grid is falling, except for the medium area where the difference increases. Further, a distinctive difference can be established when comparing the efficiencies where the difference is significantly bigger in higher flows. The increase in deviations in the area of the optimal working point and its movement in the direction of the increased airflow volume in the case of the fan without the GG is the result of the impact of the GG on the kinematics of the flow from the fan.

The results of numeric calculations of the IC with and without the GG in the three operating areas marked in the diagram are virtually the same. In both cases, the calculated points are located very near the measured characteristic of the fan without the GG. This leads to the conclusion that the GG's impact is inadequately evaluated in integral results of the numeric model, and that the model does not include all the flow characteristics on the coaxial supporters of the net.

Fig. 8 presents the dependence of the connecting electrical power on the airflow

volume on the fan. Flow $P_{el}(q_V)$ is related to the integral aerodynamic characteristic $\Delta p(q_V)$ and $\eta(q_{\rm V})$, Figure 9. The measurement results point to the division of the entire operating area in four parts. The first part between 2000 and 3600 m³/h, where the efficiency of both fans is very similar with a difference in power consumption which increases with an increased flow and is smaller in the fan with the net as it provides for smaller flows. The second part between 3600 and 6000 m³/h is where the difference between both efficiencies as well as the flow begin to increase. However, the difference between the power consumption levels starts to fall. The third part between 6000 and 7600 m³/h is where the difference between the efficiencies partly falls, yet both fans consume similar power. And the last, fourth part between 7600 m³/h and $q_{V,max}$ is where the difference between efficiency again increases but the difference between power consumption at an identical flow decreases.

This division was used as the basis for selecting three working points of the fan pertaining to individual parts. We expect to find the described characteristics to be the result of the flow kinematics and distinctly differing structures of the velocity field at the fan's outlet to be present at the selected points.

Figure 8 presents the selected points A, B and C in which measurements of the velocity field were performed as described in Section 2.2.

1.2 Local Characteristics of the Flow Field

Diagrams Figs. 10 to 15 present the results of the local characteristics. They are given in the form of a velocity distribution in the meridian plane in dependence on the selected working point (A, B and C; Fig. 8) and the position of the five-hole probe R in the radial direction from the hub to the external housing of the fan.

Presented are the measured values (the dotted line) and numerically modelled values of the velocity (the full line) of the fans with and without the GG for individual working points. Absolute velocity curves are marked as full circles, the axial components of the velocity vector with full triangles, the tangential components of the velocity vector with crosses

and the radial components of the velocity vector with empty circles, depending on the selected operating point.

Marks in the charts Figs. 10 to 12:

- Va CFD bm: Numerically calculated axial velocity of the fan without the GG;
- Va mer bm: Measured axial velocity of the fan without the GG;
- Vt CFD bm: Numerically calculated tangential velocity of the fan without the GG;
- Vt mer bm: Measured tangential velocity of the fan without the GG;
- Vr CFD bm: Numerically calculated radial velocity of the fan without the GG;
- Vr mer bm: Measured axial velocity of the fan without the GG;
- Vabs CFD bm: Numerically calculated absolute velocity of the fan without the GG and

- Vabs mer bm: Measured absolute velocity of the fan without the GG;

and for the charts in Figure 13 to Figure 15:

- Va CFD zm: Numerically calculated axial velocity of the fan with the GG;
- Va mer zm: Measured axial velocity of the fan with the GG;
- Vt CFD zm: Numerically calculated tangential velocity of the fan with the GG;
- Vt mer zm: Measured tangential velocity of the fan with the GG;
- Vr CFD zm: Numerically calculated radial velocity of the fan with the GG;
- Vr mer zm: Measured axial velocity of the fan with the GG;
- Vabs CFD zm: Numerically calculated absolute velocity of the fan with the GG and
- Vabs mer zm: Measured absolute velocity of the fan with the GG.



Fig.10. Local velocities of fan without the guard grill at operating point

Impact of the Guard Grill on the Integral and Local Characteristics of an Axial Fan

The comparison between the measured and calculated components of the velocity vector at working point A shows an excellent match between the results as regards the fans with and without the GG, Figs. 10 and 13. In both cases, the axial component of the velocity vector increases linearly with the radius – position of the measuring probe, equally in both versions (with and without the GG) its value is reduced at the circumference, which is the result of the impact of the boundary layer and impact of the gap between the blade and outlet.



Fig. 11. Local velocities of fan without the guard grill at operating point



Fig.12. Local velocities of fan without the guard grill at operating point C

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The differences between the measured absolute and modelled velocity in the area near the outlet of the fan are bigger in the case of the free outflow without the GG. These differences can be linked to deviations of measured values of the radial component of velocity for both versions of the fans and the pertaining modelled values. In the area around the housing there is a trend of an increase in the radial component of the velocity, which is more prominent in the measured values, leading to the conclusion that the flow is directed diagonally in the peripheral area.



Fig. 13. Local velocities of fan with the guard grill at operating point A



Fig. 14. Local velocities of fan with the guard grill at operating point B



Fig. 15. Local velocities of fan with the guard grill at operating point C

It is a familiar phenomenon resulting from the swirl of the flow and the related centrifugal forces [1] and [2]. At point A there is a minimal difference between the components of the tangential velocity for both versions of the fans and a relatively good match with the results of the numeric model. There is an increase in the tangential component of the velocity, which is line with the expectations [1] and [2].

The differences between both versions increase at working point B, as evident from Figs. 11 and 14. By comparing the absolute velocities we can find that the trend of increasing velocity accelerates with radius R in the entire measurement area. The distinct reduction in velocity on the fan's circumference does not appear in measured values. However, it is present on the curve presenting the modelled velocity values. At this working point and as opposed to working point A there is bigger difference between the measured and modelled axial component of velocity. The match between the measured and modelled values is also poorer in the tangential component of velocity, while in both versions of the fan there is a distinct increase in the tangential component of the velocity vector, which is in line with the literature [1] and [2] and a decrease in the axial component of the velocity vector resulting from the smaller flow. A comparison with the results for working point A shows a significant increase in the radial component of the velocity vector at the circumference which is the result of an increase in the tangential component of the velocity vector.

The trend of growing differences between the modelled and measured values increases with transition to working point C. The measurement uncertainty was greater at this point due to the smaller velocities and unregulated flow which should be taken into account in the analysis of results. For both fans, the numeric calculation showed the occurrence of a backflow in the form of a negative component of the axial velocity on the circumference which was confirmed by the measurements of the fan with the GG, Fig. 15. In comparison with the previous two working points, there is a strong increase in the tangential component of the velocity vector due to the bigger produced pressure, which is greater in the fan with the GG as the airflow has less interaction with the surroundings because of the additional GG. A significant increase in comparison with the previous measurements is also in the radial component of the velocity resulting from the increase in the tangential velocity.

The analysis of local velocities together with the integral measurements yielded the

following results for the individual operating points. In the fourth part, the GG primarily affects the size of the velocity field while no change in direction occurs. This is reflected in the IC with an almost parallel efficiency-flow curve and pressure-flow curve. The almost identical direction of the velocity field is presented in Figs. 10 and 13. As regards the fan with the GG, the trend in velocity is the same as in the fan without the GG, only it is significantly smaller. The GG affects both the size and direction of the velocity field at the third operating point, Figs. 11 and 14. The impact on the change in the direction of the velocity field is reflected in the increasing difference between the pressure-flow and efficiency-flow curves, Fig. 8, reaching the biggest difference exactly as the part ends. Around this point, the impact of the GG on the IC starts to decrease while the impact on the direction of the flow field is substantial, Figs. 12 and 15.

2 CONCLUSIONS

On the basis of the presented results, we may conclude as follows:

- (1) GG has a distinct impact on the aerodynamic properties of the AF. In addition to the aerodynamic resistance present due to the flow around the GG structure, impacts of the change in the velocity field should also be taken into account that form because of the effect of the geometry of the structure of the GG on the outflow from the fan's rotor.
- (2) By further development it is necessary to adjust the GG to the optimal on-flow velocity field from the fan without the GG to the meshed structure for decreasing the aerodynamic resistance and increasing the efficiency in the optimal working range.
- (3) Given the discrepancy between the model results and the measured values at the integral and local level, we can conclude that the GG has a distinct impact and should be included in the modelling of flow conditions on the fans. More attention should be paid to local phenomena in the flows around the GG.
- (4) Experimental measurements of the velocity field have a very good match with the numeric calculations in the operating

area of the AF. However, as the difference of the static pressure of the AF increases the irregularity of the flow rises and so the uncertainty in measuring the velocity field with a five-hole probe and in the numeric calculation grows.

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Influence of The Waste Oil Concentration in Water on The Efficiency of The Aeration Process in Refinery Wastewater Treatment

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Process, aeration system and aeration method for biological treatment of wastewater with activated sludge in bio-aeration tanks are chosen based on the flow parameters, composition of the wastewater and required characteristics of the purified water. Choosing an aeration system is a very complex question, as the capacity of oxygen introduced into the wastewater should correspond to the oxygen consumption in order to achieve the most efficient purification.

This paper presents the results of experimental aeration analysis of water contaminated with waste motor oil with variations of airflow introduced into the water. The processes of clean water aeration and water with different oil content (mass concentration from 5 to 10 mg/L). The purpose of these investigations performed on an experimental setup were analysis of the aeration process technical indicators in correlation with the thermodynamic and kinetic parameters, related to the presence of oil in the water in order to achieve more efficient refinery wastewater purification.

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Keywords: aeration system, refinery wastewater treatment, oxygen transport, relaxation processes

0 INTRODUCTION

Almost all production lines and auxiliary installations in refineries create certain amounts of wastewater containing oil during the production process. Refinery wastewater contains large amounts of oil that is removed in tanks using scrapers during pre-treatment processes. Flocculation-sedimentation and flocculation-flotation processes are used for complete removal of free hydrocarbons from water. After that, the water is biologically purified in tanks with activated sludge or in aerated lakes and lagoons.

Efficiency of aerobic biological processes in wastewater depends on the amount of dissolved oxygen introduced by aeration system. When aeration systems are designed, it is necessary to use aeration devices with an actual oxygen introduction capacity higher or equal to the actual need for oxygen in a certain phase of the treatment process of wastewater. Introduction of the necessary amount of air into bio-aeration tank and mixing of the whole suspension in the tank requires knowledge of technical characteristics of the aeration device (coefficient of oxygen transport, oxygen introduction capacity, oxygen transport efficiency and energy efficiency of oxygen transport). The coefficient of oxygen transport, k_La is the basic parameter for defining efficiency of the refinery wastewater aeration process.

Some of the technical indicators of aeration process installations stated in the literature and in the technical documentation of the producers for conditions of clean water aeration are supplied. Technical characteristics of an aeration process in the wastewater were determined in an experimental installation for clean water and water with 5 and 10 mg/L motor oil concentration and presented in this paper. Also the correlation of oxygen transport coefficient with thermodynamic and kinetic parameters related to the presence of oil in the water and air flow velocity were examined in order to achieve more efficient refinery wastewater purification.

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1 EXPERIMENTAL INSTALLATION AND DESCRIPTION OF THE INVESTIGATED PROCESS

The actual coefficient of oxygen transport was determined in an experimental installation by enriching water with oxygen. Dissolved oxygen was first chemically extracted from water and then waste oil with a defined mass concentration, was introduced in the water.

A polypropylene column with dimensions 700x700x2200 mm with accompanying connections and construction frame, was used for experimental work in batch conditions. The cross-section surface of the column was determined according to recommendations for an air distributor.

Experimental work was performed for batch working conditions and varying air flow of 2 and 10 m³/h. The water level in the column was 2m high and total volume was 980 liters. Water aeration with waste oil content of 5 and 10 mg/L was performed. Dissolved oxygen was previously removed using a chemical method of introducing sodium-sulfate in the presence of a cobaltchloride hex-hydrate catalyst. Investigation of clean water aeration not containing dissolved oxygen was also performed, in order to compare process parameters for air distribution in standard research conditions with the ones obtained in actual conditions for wastewater with corresponding characteristics.

Fig. 1 shows the experimental installation for investigating the influence of oil in water on the efficiency of the aeration process for different air flow values [1]. Table 1 contains a list of measured values and instruments on the described installation.

The column was initially filled with previously prepared water from which oxygen was extracted using a chemical method. After that, a determined amount of waste oil was added into the water.

Before experiment started, kinematical viscosity, density and surface tension of water in the experimental installation, had to be determined.



Figure 1. Diagram of experimental installation

1 -low pressure compressor; 2 -air inflow pipe valve; 3 -relief valve; 4 -air flow regulator; 5 -air flow measuring orifice plate; 6 -column with corresponding connections and framework; 7 -disk-shaped membrane air distributor; 8 -water supply; 9 -sampling connection

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experimental Complete system for defining process parameters of water aeration with certain characteristics, started by measuring the temperature of the surrounding air and water in the column. The over-pressure value before the distributor $(p_m)_d$ and the orifice plate $(p_m)_p$ i.e. the pressure difference before and after the orifice plate $(\Delta p_{\rm p})$, was measured, when the first air bubbles entered in the water. Air flow regulation was performed using a flow regulator and relief valve until a set value for the applied analysis system was attained. When the flow is stabilized, water samples were taken from the column in equal time intervals starts ($\Delta \tau = 60$ s) and the dissolved oxygen content was measured until the same value appeared three times. After one regime was analyzed, the compressor was switched off and the relief valve was opened completely. Water from the column was released into drains via a draining valve. The column was then filled with a fresh amount of water. Thus, the installation was ready for a new investigation regime, i.e. the described procedure was repeated.

2 DETERMINING TECHNICAL INDICATORS OF THE AERATION PROCESS

Oxygen transport coefficient $k_{L}a$ is a parameter used to determine the transport intensity of oxygen in water, i.e. the value at which the equilibrium state is reached.

Oxygen transport coefficient is calculated as the product of the oxygen transport coefficient in water $k_{\rm L}$ and the specific surface of the contact between air and water in the aeration process (*a*).

Based on the material balance of oxygen flow through the batch reactor with complete mixing the general equation of material balance is:

$$Qc_{\rm in} + V_{\rm G} c_{\rm ul} - V_{\rm L} R(\tau) = Qc + V_{\rm G} c_{\rm iz} + V_{\rm L} \frac{\mathrm{d}c}{\mathrm{d}\tau}, \qquad (1)$$

where:

$$V_{G}$$
 - air flow, m³/s

Q - water flow (for batch process conditions Q = 0), m³/s,

Table 1. Description of measuring equipment (list of measured values and instruments)

	Measured value	Denotation	Unit	Instrument	
1.	Air temperature	t _G	°C	Mercury thermometer	
2.	Air over-pressure	$(p_{\rm m})_{\rm p}$	mm Hg	U-pipe with mercury	
3.	Pressure loss through the orifice plate	$\Delta p_{ m p}$	mm Alc	Micro manometer with alcohol	
4.	Air over-pressure in front of the air distributor	$(p_{\rm m})_{\rm d}$	mm Hg	U-pipe with mercury	
5.	Water temperature in the column	$t_{ m L}$	°C	Mercury thermometer	
6.	Kinematical viscosity of water in the column	$ u_{ m L}$	m ² /s	Modified Oswald viscosimeter	
7.	Water density in the column	$ ho_{ m L}$	kg/m ³	Laboratory pycnometer	
8.	Water surface tension	$\sigma_{ m L}$	N/m	Microbalance	
9.	Starting mass concentration of dissolved oxygen in water	Co	mg/L	"HANNA INSTRUMENTS 9142"	
10.	Mass concentration of dissolved oxygen in water	$c(\tau)$	mg/L		

 $c_{\rm in}$ - mass concentration of oxygen in the influent, kg/m³,

c - mass concentration of oxygen in the effluent, kg/m^3 ,

 $c_{\rm ul}$ - mass concentration of oxygen in air at the input, kg/m³,

 c_{iz} - mass concentration of oxygen in air at the output, kg/m³,

 $R(\tau)$ - specific oxygen consumption during biological treatment, kg/(m³·s),

 $V_{\rm L}$ - water volume, m³.

The gas phase material balance gives:

$$\dot{V}_{\rm G}\left(c_{\rm ul}-c_{\rm iz}\right) = AK_{\rm L}\left[c^*\left(c_{\rm iz}\right)-c\right],\qquad(2)$$

where:

 $c^* = c^*(c_{iz})$ - equilibrium mass concentration of oxygen in dependence on the mass concentration

of oxygen in air at the output, kg/m³, $A = aV_{L}$ - total contact surface between air and water, m²,

a - specific surface of contact between air and water, m^2/m^3 .

The basic resistance to oxygen transport through the system occurs in a liquid hence the total transport coefficient of oxygen by water (K_L) is approximately equal to the oxygen transport coefficient through water (k_L).

Mass concentration of oxygen in air at the output, based on modified Henry's constant (*Hc*), is determined as follows:

$$c_{\rm iz} = Hcc^*(c_{\rm iz}). \tag{3}$$

Modified Henry's constant (*Hc*), is determined by equation:

$$Hc = Ha/(R_{\rm u}T_{\rm G}C_{\rm L}), \qquad (4)$$

where:

Ha- Henry's constant, Pa·kmol (O_2+L)/kmol O_2 , R_{μ} - the universal gas constant, J/(kmolK),

 $T_{\rm G}$ - absolute air temperature, K,

 $C_{\rm L}$ - molar concentration of oxygen in water, kmol/m³.

Equations (3) and (4) give the expression for calculating the output mass concentration of oxygen in air:

$$c_{iz} = \frac{Ha(V_{G} c_{ul} + V_{L} k_{L} ac)}{V_{G} Ha + V_{L} k_{L} a(R_{u} T_{G} C_{L})},$$
(5)

where:

 $k_{\rm L}a$ - coefficient of oxygen transport, 1/s.

Time dependence of dissolved oxygen concentration is presented by the following equation:

$$Y=mX,$$
 (6)

where:

$$Y = ln \frac{\frac{p_{ul}M_{O_2}}{R_u T_G} - \frac{Ha}{R_u T_{OK} C_L} c_o}{\frac{p_{ul}M_{O_2}}{R_u T_G} - \frac{Ha}{R_u T_{OK} C_L} c(\tau)},$$
(7)

m - slope coefficient of the equilibrium curve,

$$m = \frac{\frac{Ha}{R_{\rm u}T_{\rm OK}C_{\rm L}}k_{\rm L}a}{\frac{Ha}{R_{\rm u}T_{\rm OK}C_{\rm L}} + \frac{V_{\rm L}}{V_{\rm G}}k_{\rm L}a} - R(\tau),$$
(8)

$$X = \tau. \tag{9}$$

The value of Y is given for each value of $c(\tau)$ determined by experiments for defined time intervals.

The previous equations result in formulae for determining the coefficient of oxygen in wastewater as:

$$k_{\rm L}a = \frac{Ha(m+R(\tau))\dot{V}_{\rm G}}{\dot{V}_{\rm G} Ha - (m+R(\tau))V_{\rm L}R_{\rm u}T_{\rm ok}C_{\rm L}} \cdot$$
(10)

The coefficient of oxygen transport for standard conditions is determined in the following way [2]:

$$(k_{\rm L}a)_{\rm s} = \frac{(k_{\rm L}a)_{\rm t_{\rm L}}}{\theta^{\rm t_{\rm L}-20}},$$
 (11)

where:

 $(k_{\rm L}a)_{\rm t_r}$ - experimentally obtained oxygen

transport coefficient, 1/s,

 $t_{\rm L}$ - water temperature, ^oC ,

 $\theta = 1.024$ - temperature correction factor.

Based on the known value for the coefficient of oxygen transport $(k_L a)_s$ the following technical characteristics of aeration systems for aeration of wastewater are calculated

[3]: actual capacity of oxygen introduction OC', actual efficiency of oxygen transport E' and actual energy efficiency of oxygen transport E_e' .

Actual capacity of oxygen introduction is calculated as the product of the standard capacity of oxygen introduction and corresponding correction factors that convert standard investigation conditions to real ones:

$$OC' = \alpha OC \frac{\beta c_{\rm h}^* - c_{\rm o}}{c_{\rm s}^*} \theta^{(t_{\rm L} - 20)}, \qquad (12)$$

where:

OC - standard capacity of oxygen introduction into wastewater, kg/h,

$$OC = (k_{\rm L}a)_{\rm s} c_{\rm s}^* V_{\rm L}$$

 $c_{\rm s}^*$ - equilibrium mass concentration of dissolved oxygen in clean water, for standard conditions, kg/m³,

 $\alpha = 0.8 \div 0.94$ - relative transport degree of oxygen in wastewater,

 $\beta = 0.90 \div 0.97$ - relative saturation degree of wastewater with oxygen,

 c_h^* - equilibrium mass concentration of dissolved oxygen in clean water in actual conditions, corrected in relation to the height of the liquid column above the air distributor and molar participation of oxygen in air, mg/L,

$$c_{\rm h}^* = c_{\rm s}^* \left[1 + \frac{\rho_{\rm L}g}{p_{\rm n}} (H - h) \right]$$

 $p_{\rm n}$ - pressure corresponding to standard conditions, Pa,

H - total height of the water column, m, h - height from the bottom of the tank to the distributor, m,

 $\rho_{\rm L}$ - water density, kg/m³.

The actual capacity of oxygen introduction *OC*' into wastewater should correspond to oxygen consumption during biological treatment.

The actual efficiency of the transport system is expressed in percentages and represents the ratio between the real capacity of oxygen introduction and the total oxygen flow brought by the aeration system [4]:

$$E' = \frac{OC'}{\dot{G}_{O_2}} = \frac{OC'}{\left(\dot{V}_{G}\right)_n \rho_G y_{O_2}} \cdot 100,$$
(13)

where:

 G_{O_2} - mass flow of oxygen introduced into the water by the aeration system, kg/h,

$$\left(\dot{V_G} \right)_n$$
 - volume air flow
(for $p_n = 101.3$ kPa, $t_L = 20^\circ$ C), m³/h,

 $\rho_{\rm G}$ - air density for standard conditions ($t_{\rm G} = 20^{\circ}$ C), kg/m³,

 $y_{O_2} = 0.232$ - mass participation of oxygen in air, kg/kg.

The real energy efficiency of oxygen transport represents the ratio between the actual capacity of oxygen introduction and the applied power needed to drive the aeration device:

$$E'_{e} = \frac{OC'}{\sum_{i} P_{i}},$$
(14)

where:

 $\sum_{i} P_{i}$ - sum of power off all electrical motors

(aerator, pump, blower, etc.), kW.

Based on all the experiments performed on the described installation for varying air flow between 2 and 10 m³/h and waste motor oil content from 0 to 10 mg/L using the investigated model $(k_L a)_s$ values given in Table 2 were determined. The values obtained for $(k_L a)_s$ were used to determine the values of actual capacity of oxygen introduction *OC'*, real oxygen transport efficiency *E'* and actual energy efficiency of oxygen transport E_e' also given in Table 2.

Technical characteristics of the aeration system in dependence on oil presence in wastewater, are given as diagrams in Fig. 2.

An increase of polluting components' concentrations (oil) in the water, led to a reduction of oxygen transport coefficient values and affecting other technical characteristics of the investigated aeration system.

Waste oil concentration in $(V_{\rm G})_n$ $(k_{\rm L}a)_{\rm s}$ OC'Regime water E' E_{e} number m³/h 1/h % g/kWh mg/L g/h 2.361 2.887 11.754 1.465 35.078 5.843 2 3.052 12.425 0.617 15.024 -0.550 14.851 3 10.768 5.267 21.876 _ 5 28.442 4 2.361 2.248 9.528 1.118 5 5 5.862 2.378 9.519 0.463 11.510 5 4.527 18.999 0.625 12.898 6 10.735 7 2.379 1.993 8.204 24.489 10 1.002 8 5.899 2.160 8.837 0.426 10.686 10 9 10.724 2.577 10.730 0.271 7.284 10 6,0 30 5,5 5.0 25 4,5 4,0 20 **€** ^{3,5} oc', g/h (**k**^{-3,0} (**k**⁻³) 2,5 15 2,0 10 1.5 1,0 5 0,5 0,0 0 4 10 12 0 6 8 b) a) 2 0 4 6 8 10 12 (V_G)_n, m³/h (V_G)_n, m³/h 1,6 40 1,4 35 1,2 30

Table 2. Technical characteristics of the aeration system for corresponding experimental regimes



d)

25 g/kWh

20 Ë.

15 10

5

2

Content in water 5 mg/l

6

(V_G)_n, m³/h

8

10

12

OC'; c) actual oxygen transport efficiency E'; d) actual energy efficiency of oxygen transport E'_e

1,0

0,6

0,4

0,2 0,0

c)

0

2

4

6

(V_G)_n, m³/h

8

10

12

Content in water 0 mg/l

---- Content in water 10 mg/l

Е, % 0.8
3 DEPENDENCY OF TECHNICAL INDICATORS ON THERMODYNAMIC INDICATORS AND PRESENCE OF OIL IN THE WATER DURING AERATION PROCESS

The law of energy balance can explain reduction of oxygen transport coefficient in the water during aeration with increasing waste oil concentration:

$$\frac{dE_{tot}}{d\tau} = \frac{dE_{kin}}{d\tau} + \frac{dE_{pot}}{d\tau} \,. \tag{15}$$

Change velocity of the total oxygen molar energy, E_{tot}^{Θ} is determined with oxygen transport coefficient, $k_{\text{L}}a$ and influenced by the diffusion rate constant, k_{d} , the kinetic molar energy, $E_{\text{kin}}^{\Theta} = \frac{3RT}{2}$, as well as the adsorption rate constants, k_{a} in potential field through the free liquid surface, a [5]:

$$k_{\rm L}aE^{\Theta}_{\rm tot} = k_{\rm d}2E^{\Theta}_{\rm kin} + k_{\rm a}E^{\Theta}_{\rm pot}, \qquad (16)$$

or in the form:

$$k_{\rm L}a = k_{\rm d} + (k_{\rm a}/E^{\Theta}_{\rm tot}) \Delta_{\rm ads}G^{\Theta}$$
(17)

The potential molar energy, E_{pot}^{Θ} determined with oxygen bond strength, is equal with the oxygen molar adsorption Gibbs energy, $\Delta_{ads} G^{\Theta}$, and can be affected by the presence of the oil oxidation products and radical initiators (OH, H₂O₂, Fe²⁺, HCO₃⁻/CO₃⁻²⁻, PO₄⁻³⁻, humid acids, aryl-R etc) [6]. Standard oxygen adsorption Gibbs energy in this work is determined with the Gibbs-Helmholtz - van' t Hoff equation [7], applied for saturated water:

$$\mu(O_2)_{a q} = \mu(O_2)_g. \tag{18}$$

The chemical potential, $\mu(O_2)$, is determined with the standard value, $\mu^{\theta}(O_2)$ and with oxygen concentration ($c(O_2)$):

$$\mu(O_2) = \mu^{\theta}(O_2) + RT \ln c(O_2).$$
(19)

Determination of adsorption molar energy on the base the Equations (18) and (19) is performed using the obtained linear function (20):

$$\ln c(O_2)_{a q} = \ln c(O_2)_g - \Delta_{ads} G^{\Theta}(O_2)/RT, \quad (20)$$

and presented in the Table 3 as:

$$y = y_o - \operatorname{tg} \varphi \cdot x,$$

where:

$$\begin{split} y &= \ln c \ (\mathrm{O}_2)_{\mathrm{a}\,\mathrm{q}} \ , \\ y_\theta &= \ln c \ (\mathrm{O}_2)_{\mathrm{g}} \ , \\ \mathrm{tg} \ \varphi &= \Delta_{\mathrm{ads}} \ G^\Theta \left(\mathrm{O}_2\right) / R, \ x = 1 / T. \end{split}$$

The calculated oxygen adsorption Gibb-s energies for the investigated regimes on the basis of the obtained linear functions to equation (21) are presented in the Table 4:

$$\Delta_{\mathrm{ads}} G^{\Theta}(\mathcal{O}_2) = R \cdot g \, \varphi \,, \tag{21}$$

The calculated oxygen kinetic parameters (the adsorption rate constants and diffusion rate constants) in aeration process:

$$k_{\rm a} = 3RT {\rm tg} \,\varphi, \tag{22}$$

$$k_{\rm d} = y_{\rm o} \,, \tag{23}$$

are presented in the Fig. 3 and the Table 5.

Investigated regimes defined with: <i>c-h-q</i>	$y = y_0 - tg \varphi \cdot x$ were is to eqation (20) $y = \ln c(O_{2,aq}) \text{ and } x = 1/T$	Coefficient correlation, R^2
0-2-2, 0-2-6, 0-2-10	y = -3663.1x + 4.4554	0.9219
0-2-2, 0-2-6, 5-2-10, 5-1-6	y = 4367x - 23.624	0.9751
0-2-10, 5-2-10, 10 -2-10	y = 165691x - 585.58	0.7455
5-2-2 0-2-10, 5-1-2, 5-1-10	y = 8959.3x - 39.575	0.6137
5-2-6, 10-2-2, 10-2-6	y = 7726.9x - 35.569	0.9939
10-2-10, 10-1-2, 10-1-6, 10-1-10	y = 8492.3x - 38.096	0.9393

Table 3. Functions derived from the Equation (20) in conditions defined with c-h-q, c (oil concentration), mg/l, - h (height of water column), m, - q (air flow), m^3/h



Table 4. Thermodynamic and technical characteristics of examined regimes



Regime number	$(\overset{\cdot}{V_G})_n$ m ³ /h	$ \begin{array}{c} \Sigma \varDelta_{diss} G^{\Theta} \\ (\mathrm{O}_{2}_{,\mathrm{g/aq}}) \\ \mathrm{kJ/mol} \end{array} $	<i>k_La</i> 1/min	<i>k_d</i> 1/min	<i>k</i> _{<i>a</i>} 1/min	$1/\tau_m$ * 1/min	Compared measured with the calculated aeration rate constants
1	2.361	-5.9	0.0419	0.043	- 0.214	0.042	$\frac{1/\tau_m = k_L a = k_d}{(1/\tau_m)/k_a = 0.19}$
2	5.843	-5.9	0.0443		"	0.045	$1/\tau_m = k_L a = k_d$ (1/ τ_m)/k a = 0.21
3	10.768	1346	0.0779		"	0.048	$\frac{1/\tau_m / k_L a = 0.62}{1/\tau_m / k_d = 1.11}$ $1/\tau_m / k_a = 0.22$
4	2.361	-74.5	0.0339	0.032	- 0.143	0.04	$(1/\tau_m)/k_L a = 1.17$ $(1/\tau_m)/k_d = 1.25$ $(1/\tau_m)/k_a = 0.28$
5	5.862	-64.2	0.0339	"	"	0.043	$(1/\tau_m)/k_L a = 1.27$ $(1/\tau_m)/k_d = 1.34$ $(1/\tau_m)/k_a = 0.3$
6	0.735	-1441.2	0.0677	"	"	0.053	$(1/\tau_m)/k_L a = 0.78$ $(1/\tau_m)/k_d = 1.65$ $(1/\tau_m)/k_a = 0.37$
7	2.379	-64.2	0.0292	0.03	-0.043	0.042	$(1/\tau_m)/k_L a = 1.42$ $(1/\tau_m)/k_d = 1.39$ $(1/\tau_m)/k_a = 0.97$
8	5.899	-64.2	0.0315	"	"	0.032	$(1/\tau_m)/k_L a = 1.4$ $(1/\tau_m)/k_d = 1.4$ $(1/\tau_m)/k_a = 1.0$
9	10.724	-1447.6	0.0382			0.053	$(1/\tau_m)/k_L a = 1.38 (1/\tau_m)/k_d = 1.75 (1/\tau_m)/k_a = 1.22$

Table 5. Theoretical values of the relaxation processes rate constants in the aeration system for corresponding investigation regimes

* measured period of aeration time enabling that constant oxygen concentration was achieved

The oxygen transport coefficient as technical indicator of aeration kinetic in correlation with the rate constants, obtained on the base of the thermodynamic study of aeration kinetic as in correlation with directly measured saturation time is presented in the Table 5.

In clean water, the oxygen transport coefficient is equal with the diffusion rate constant and with parameter calculated on the base saturation time ($k_{\rm L}a = k_{\rm d} = 1/\tau_{\rm m}$), but adsorption rate constants were five time greater, $k_{\rm a}/(1/\tau_{\rm m}) = 4.8 \div 5.3$. In the presence of added oil, relative adsorption rate constants were from $3.6 \div 2.7$ for 5 mg/l oil concentration up to 1.34 $\div 0.82$ for 10 mg/L oil concentration. Equal adsorption rate constants with the saturation rate constants reduced aeration efficiency due to activated oxygen in the slow adsorption processes at 10 mg/L oil concentration.

4 CONCLUSION

Experimental research shows that 5 mg/L concentration of oil in the water, leads to reduction of the dissolved oxygen concentration from 10 to 15% compared to clean water, while 10 mg/L concentration of oil in the water reduces concentration of dissolved oxygen from 20 to 25% when compared to clean water. Aeration process technical indicators (oxygen introduction capacity. oxygen transport efficiency and energy efficiency of oxygen transport) for the clean water and water with 5 mg/L oil concentrations were significantly better with the maximum air flow, $10 \text{ m}^3/\text{h}$, than with the less flow of 2 m³/h and 6 m³/h. For an oil concentration of 10 mg/L, the air flow of 10 m³/h did not have a significant influence on the aeration process technical indicators.

Based on the research of the dependence of technical indicators on thermodynamic parameters and presence of oil in the water during aeration process it can be concluded that:

- 1. The oxygen transport coefficient is equal with the diffusion and saturation rate constants for clean water and contaminated water at the air flow of up to $6 \text{ m}^3/\text{h}$.
- 2. The oxygen transport coefficient for 10 m³/h air flow:
- is increased compared to diffusion rate constants 1.8 times and for water with 5 mg/L oil concentration 2.1 times for clean water
- is only 1.2, except for water with 10 mg/l oil concentration when the relative oxygen transport coefficient was compared to diffusion rate constants
- 3. The study of the Gibbs energy adsorption dependence from oil concentrations shows:
- oxygen adsorption is activated with parallel endothermic and exothermic oxygen processes at the air flows up to 2 m³/h and 6 m³/h in the clean water,
- oxygen adsorption is activated with spontaneous, exothermic electrons transitions with the most adsorption affinity at the air flow of 10 m³/h in the contaminated water.
- 4. Adsorption rate constants with increased oil concentrations become less negative and diffusion rate constants less positive values

of up to 10 mg/L oil concentration, when equal adsorption and saturation rate constants prevent the air flow influence on increasing the oxygen transport coefficient.

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Strength of Pressure Vessels with Ellipsoidal Heads

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A method for stress analysis in cylindrical pressure vessels with ellipsoidal heads, based on the axisymmetric shell theory, was proposed. The starting point were the approximate solutions of the differential equation system that were used to get mathematical expressions for determining internal forces, moments and displacements in the vessel walls.

Final expressions that can be applied were acquired by joining the membrane and moment theory and by setting and solving equations of boundary conditions. Diagrams that show distribution of internal forces and moments are in dimensionless form which enables their use regardless of dimensions and load. These expressions were used to develop a method for testing strength of pressure vessels with ellipsoidal heads in the design phase. Application of the method was shown on a selected numerical example, while a special computer programme was created for calculation purposes.

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Keywords: pressure vessels, ellipsoidal heads, internal forces, analytical solutions

0 INTRODUCTION

Because of the increasing use of pressure vessels in all types of industry (e.g. oil-refining industry, gas-processing industry, power plant industry, etc.), in transportation, and other processing facilities, and because of their many side effects on the environment, there is a need for accurate analysis of the shells in the design phase. The designer must know the real distribution of stresses, i.e. he must make calculations and analysis of stress in all critical parts of the vessel in order to make an adequate pressure vessel design. Stress is highest at the points of geometry parameter change, e.g. the transition from the shell of the vessel to the head, where higher wall bending usually occurs.

Simple empirical forms determined by norms are used for making calculations of stress in the walls of pressure vessels. This type of calculation does not give the real distribution of stress, the highest values or places where stress occurs. More accurate stress analysis can be done by using finite element method (FEM) or analytical method.

Analytical method for calculating strength of the thin-walled pressure vessels is based on the theory of axisymmetric shells. This theory has a system of the fourth order differential equation with internal forces, moments and displacements as unknown functions. The shell theory was established by famous researchers (Lourye 1947, Goldenveiser 1953, Novozhilov 1962 and others) who developed analytical methods for determining general solutions of the system. Accurate solutions of the system can be acquired in general form only for pressure vessels of simple geometry (e.g. cylindrical, conical and spherical). In order to apply those solutions in the process of designing, we must calculate numerous complex functions and infinite orders with complex variables, and we must also set and solve an equation of boundary conditions. These solutions can be used to make calculations on a computer, but the calculation procedure and the algorithm structure are very complex.

Far simpler analytical procedure for calculating pressure vessel strength can be developed on the basis of approximate solutions of axisymmetric shell theory equation system [1]. It is of vital importance because the approximate solutions contain an error of the same order of magnitude as do the correct solutions. Relative error of those solutions is of the order h/R (ratio of the wall thickness and the curvature radius), which is less than 1/20 at thin-walled pressure vessels, i.e. 5%, which can be neglected in engineering calculations.

Approximate solutions of the shell theory can be applied to pressure vessels with more complex shells.

This paper presents solutions for a cylindrical vessel with ellipsoidal heads. Final mathematical expressions for calculating internal

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forces, moments and displacements in the walls of head and cylindrical part were derived from approximate solutions. By using those expressions, and by connecting membrane and moment theory, a method for determining strength of pressure vessels with ellipsoidal heads, which is suitable for designing, was developed. A special computer programme was created for the application of this method. Computer calculation was done on a selected numerical example, and the analysis results were shown in a diagram.

1 MEMBRANE FORCES IN ELLIPSOIDAL HEAD

Stresses in the walls of pressure vessels occur due to different types of loads, depending on the purpose of the vessel and on different influences that a vessel is subjected to during exploitation. Internal pressure has the biggest influence on the amount of stress, so all other types of loads are considered to be less important.

Solutions of the shell theory equation show that internal forces which occur in the walls of the vessel can be, under certain conditions, determined by superimposing of two kinds of loads, membrane and moment. These conditions given in [2] are true for curvature radius derivations of vessel meridian R_1 and surface load q_n .

At membrane stress state, it is assumed that only normal forces occur in the wall, while shear forces and bending moments are not considered. In the design phase, pressure vessels should be formed so that the real stress is approximately the same as the membrane state in order to avoid bending of the walls and high stresses due to bending.

Cylindrical pressure vessels usually have semi-spherical, toroidal-spherical or ellipsoidal heads. Stress is best distributed in ellipsoidal heads. The middle plane of ellipsoidal head is shaped like a half of the ellipsoid of revolution which is the result of ellipse arc rotation around its minor axis. A major ellipsoid semi-axis is equal to cylindrical part radius (a = R), and a minor semi-axis is equal to the head height (b =H). A parameter determining ellipse form is $\gamma = a^2/b^2 - 1$, and $R_0 = a \cdot \sqrt{1+\gamma}$ is curvature radius of head vertex. Figure 1 shows the basic measures of cylindrical pressure vessels, closed by heads shaped like revolution ellipsoid half.



Fig.1. Pressure vessel with ellipsoidal heads

Internal forces in the head walls are determined according to the membrane theory [2] by using the following expressions:

$$N_{s}^{(0)} = \frac{p R}{2} \left(\frac{1+\gamma}{1+\gamma \sin^{2} \theta} \right)^{1/2},$$
 (1)

$$N_t^{(0)} = \left(1 - \gamma \sin^2 \theta\right) N_s^{(0)}.$$
 (2)

Membrane component of circular force has (in places of shell connection) tensile character on the cylinder, and compressive character on the head, which leads to difference in radial displacements according to [3]:

$$u_r^{(0)} - u_r^{(0)} = \frac{p R^2}{2 E h} (1+\gamma), \qquad (3)$$

where:

 $N_s^{(0)}$ - meridian normal force in the head wall,

 $N_t^{(0)}$ - circular normal force in the head wall, $u_r^{(0)}$ - radial displacement of the wall (at the head),

 $u_r^{(0)}$ - radial displacement of the wall (at the cylindrical part),

E - modulus of elasticity of the pressure vessel wall,

h - thickness of the wall,

- *p* pressure in the vessel,
- θ meridian angle,
- *R* radius of cylindrical part of the vessel.

Superscript⁽⁰⁾ denotes values pertaining to a membrane stress condition. Corresponding boundary condition values will be denoted by the superscript⁽¹⁾, and those without the superscript refer to total solutions.

As can be seen from the expression (1), meridian normal force of ellipsoidal head changes without interruption in the points of transition from the ellipsoid to the cylinder, which is very important for achieving the membrane state of stress. Both normal forces (meridian and circular) reach their maximum value in the head vertex. Value $\gamma = 3$ should be chosen for ellipse parameter, because maximum values of internal forces in the head are the same as values of circular force in the cylindrical part of the vessel.

There is a discontinuity of the membrane component of circular force in the places where head and the cylinder are connected. Since total displacements of the vessel walls should stay intact, additional forces and moments occur in the vessel walls, which lead to bending of the walls. Expressions for determining their values can be acquired from general solutions of shell theory differential equation system.

2 FORCES AND MOMENTS OF BOUNDARY EFFECT

Thin-walled axisymmetric shells can be divided according to the moment theory into short and long. The shells are considered long if:

$$\lambda = \int_{0}^{s_1} \beta \cdot ds \ge 3 , \quad \beta = \frac{\sqrt[4]{3(1-v^2)}}{\sqrt{R_2 \cdot h}}.$$
(4)

With long shells we can disregard the influence of the load of one end of the shell on the internal forces and displacements on the other end. Each end of the shell can be observed independently (i.e. without considering the conditions at the opposite end). Bending of the shell occurs due to radial forces F_0 and bending moments M_0 at the boundary of the shell, which is called boundary load. Boundary load of the long axisymmetric shell is shown in Figure 2.

The condition (4) is regularly achieved for the heads and cylindrical parts of thin-walled pressure vessels.



Fig.2. Boundary load of the long axisymmetric shell

With those conditions, general solutions of axisymmetric shell theory differential equation system, acquired by approximate solutions of the system by [4] can be reduced to:

$$u_{r}^{(1)} = -\frac{F_{0}}{2D\beta_{0}^{3}} \sin^{2}\theta_{0} \cdot e^{-x} \cos x -\frac{M_{0}}{2D\beta_{0}^{2}} \sin\theta_{0} \cdot e^{-x} (\cos x - \sin x),$$
(5)

$$\begin{aligned} \boldsymbol{\theta}^{(1)} &= -\frac{F_0}{2D\beta_0^2} \quad \sin\theta_0 \cdot e^{-x} \ \left(\cos x + \sin x\right) \\ &- \frac{M_0}{D\beta_0} \cdot e^{-x} \cos x \qquad , \end{aligned}$$
(6)

$$F_{r}^{(1)} = F_{0} \cdot e^{-x} \left(\cos x - \sin x\right) \\ - \frac{2M_{0} \ \beta_{0}}{\sin \theta_{0}} \cdot e^{-x} \sin x , \qquad (7)$$

$$M_{\rm s} = \frac{F_0}{\beta_0} \quad \sin \theta_0 \cdot e^{-x} \sin x$$
$$M_0 \cdot e^{-x} \left(\cos x + \sin x\right) , \tag{8}$$

$$M_{t} = v \cdot M_{s}, \quad N_{t}^{(1)} = \frac{Eh}{r} \cdot u_{r}^{(1)}, \\ N_{s}^{(1)} = F_{r}^{(1)} \cdot \cos\theta \qquad (9)$$

where:

j

$$x = \int_{0}^{s} \beta \cdot ds , \quad D = \frac{E h^{3}}{12 (1 - v^{2})}.$$
 (10)

In expressions (4) to (10) the following was used:

 $F_{\rm r}^{(1)}$ - radial force in the wall at boundary load,

 $N_{\rm s}^{(1)}$ - meridian normal force in the wall,

 $N_{\rm t}^{(1)}$ - circular normal force in the wall,

 $u_r^{(1)}$ - radial displacement of the wall,

 $\mathcal{G}^{(1)}$ - wall twist angle,

 $M_{\rm s}^{(1)}$ - meridian moment of bending,

 $M_{t}^{(1)}$ - circular moment of bending,

v - Poisson's ratio,

s - meridian arc length from the shell boundary to a certain point,

 s_1 - total shell length,

r - radius of the middle plane in a certain point of the wall,

R- circular radius of wall curvature.

All expressions with a subscript $_0$ are values for the shell boundary. For example, θ_0 and β_0 are values of θ and β in points x = 0, at the shell boundary.

In order to apply expressions (5) to (8) it is necessary to determine the values of the force F_0 and moment M_0 of boundary load. They are, by their nature, internal forces of moment theory at the joint of two different shells that superimpose forces of membrane theory, in order to achieve inner mechanical balance.



Fig.3. Forces and moments of boundary effect

Their values for the example under consideration (cylindrical part of the vessel with ellipsoidal heads) in Figure 3 will be determined by setting equations of boundary conditions.

Strain continuity at the joint (i.e. place where cylinder and ellipsoid are connected) depends on equal total displacements and twist angles of joined parts, which is expressed by the following equations:

$$u_r^{(0)'} + F_0 \cdot \delta_{11}' - M_0 \cdot \delta_{12}' = = u_r^{(0)''} - F_0 \cdot \delta_{11}'' - M_0 \cdot \delta_{12}'', \qquad (11)$$

$$-F_0 \cdot \delta_{21}' + M_0 \cdot \delta_{22}' = = -F_0 \cdot \delta_{21}'' - M_0 \cdot \delta_{22}'' .$$
(12)

Membrane components of the wall twist angle, both for ellipsoidal and cylindrical shell, equal null at the angle $\theta = \pi/2$, which can be shown by using membrane theory [5]. Coefficient values of generalized forces influence (F_0 and M_0) on generalized strain (u_r and ϑ), with long shells, according to [4] are:

$$\delta_{11}' = \delta_{11}'' = \frac{\sin^2 \theta_0}{2D\beta_0^3}, \ \delta_{12}' = \delta_{12}'' = \frac{\sin \theta_0}{2D\beta_0^2},$$
$$\delta_{22}' = \delta_{22}'' = \frac{1}{D\beta_0}, \tag{13}$$

where:

$$\beta_0 = \frac{\sqrt[4]{3(1-\nu^2)}}{\sqrt{Rh}} \,.$$

By solving equations system (11) and (12), and by considering values (3) we get:

$$M_0 = 0$$
, $F_0 = \frac{p}{8 \beta_0} (1 + \gamma)$. (14)

By substituting solution (14) into expressions (5) to (10), we get formulas for internal forces and displacements that occur due to boundary effect in the walls of ellipsoidal (') and cylindrical ('') part of the vessel:

$$u_{\rm r}^{(1)} = \frac{p R^2}{4 E h} (1+\gamma) e^{-x} \cos x , \qquad (15)$$

$$u_{\rm r}^{(1)} = - \frac{p R^2}{4 E h} (1+\gamma) e^{-x} \cos x , \qquad (16)$$

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$$M_{\rm s}' = - \frac{p}{8 \beta_0^2} (1+\gamma) e^{-x} \sin x, \qquad (17)$$

$$M_{\rm s}" = \frac{p}{8 \beta_0^2} (1+\gamma) e^{-x} \sin x , \qquad (18)$$

$$N_{t}^{(1)} = \frac{pR}{4} (1+\gamma) e^{-x} \cos x, \qquad (19)$$

$$N_t^{(1)} = - \frac{pR}{4} (1+\gamma) e^{-x} \cos x, \qquad (20)$$

$$F_{\rm r}^{(1)} = - \frac{p\sqrt{2}}{8\beta_0} (1+\gamma) e^{-x} \cos(x+\pi/4), \qquad (21)$$

$$F_{\rm r}^{(1)} = \frac{p\sqrt{2}}{8 \beta_0} (1+\gamma) e^{-x} \cos(x+\pi/4), \qquad (22)$$

where:
$$x = \int_{0}^{s} \beta \, ds$$
, $\beta = \frac{\sqrt[4]{3(1-v^2)}}{\sqrt{R_2 h}}$. (23)

Distribution of internal forces and moments along the meridian length of the pressure vessel wall, according to expressions (15) to (22), is shown in the diagrams (Figures 4, 5, and 6).

According to the diagram, circular force $N_t^{(1)}$ (Fig.4) and meridian bending moment M_s (Fig. 5) have the biggest influence on stress. Circular moment M_t , according to (9), has smaller values, and components of radial force $F_r^{(1)}$ are of lower order of magnitude so that they can be disregarded in stress calculation.

4 METHOD FOR TESTING STRENGTH

Total values of internal forces in the pressure vessel walls can be acquired by adding components of membrane and moment theory:

$$N_t = N_t^{(0)} + N_t^{(1)}$$
, $N_s = N_s^{(0)} + N_s^{(1)}$. (24)

If we analyse expressions (17) to (22), the following can be concluded. For designing purposes values of internal forces and moments can be observed as dimensionless magnitudes, e.g. $N_s/(p \cdot R)$ and $M_s/(p \cdot R \cdot h)$. Beside the position of the point on the vessel wall, they depend also on design parameters, i.e. on the head shape γ and ratio R/h.



Fig.4. Distribution of circular forces along the vessel wall determined by moment theory



Fig.5. Distribution of meridian moments along the vessel wall determined by moment theory



Fig.6. Distribution of radial forces along the vessel wall determined by moment theory

Therefore, by choosing values of these parameters we can test strength before determining final dimensions of the vessel.

In order to perform this procedure successfully, a special computer programme (in Fortran 77) [6], based on mathematical expressions, was created.

This programme can be used to calculate internal forces and moments by using expressions (1) and (2), (15) to (24) in a number of points distributed along the meridian of the vessel. It can be also used to determine values of main stresses on the inner and outer surface of the wall and to calculate equivalent stresses.

The programme is made in such a way that calculation is performed for the chosen ratio R/h and factor of head shape γ . It was carried out on the example R/h = 87 and $\gamma = 3$, and calculated values of internal forces and moments are shown in Table 1 only for a limited number of points. Distribution of total circular force, according to calculated values, is shown in Fig. 7.



Fig.7. Distribution of total circular forces of the cylindrical vessel with ellipsoidal heads

Equivalent stress according to HMH theory of strength [7] was chosen as a criterion for testing strength of pressure vessel. This programme can be used to calculate equivalent stresses at different points of the wall and to find critical points. Maximum value of equivalent stress was determined for the above mentioned example:

$$(\sigma_{\rm e})_{\rm max} = 1.42 \cdot \frac{p \cdot R}{h}$$
 (25)

which occurs in the points at the inner wall surface at meridian angle $\theta = 69.61^{\circ}$ (which is equal to polar angle $\varphi = 5.31^{\circ}$). On the outer surface of the wall, maximum stress occurs at the head vertex, and it is much lower in the mentioned example.

Stress determined by the expression (25) is the basis for the choice of pressure vessel material (concerning the necessary strength), after the shape and dimensions of the pressure vessel have been determined. These two steps in the design phase can be done iteratively, i.e. by varying design R/h and γ , as input data, we can

get values of maximum equivalent stresses for different dimension ratia.

5 CONCLUSION

Approximate solutions of axisymmetric shell theory used for expressions (15) to (22) are valid for steep shells, i.e. those with big θ angle. This condition is found in the vicinity of the place where the head and the cylinder are joined (i.e. where angle θ slightly differs from 90°). It is better to use these expressions instead of exact solutions because they consist of simple mathematical functions, so that there is no need to consider boundary conditions, since they are already included in final expressions.

These expressions clearly show the law of distribution of internal forces and moments in the walls by using diagrams, which was done for ellipsoidal heads with R/H=2 ratio, since they are commonly used in praxis. These diagrams can be used for analysis of stress and for choosing dimensions during the design phase of such heads. Computer programme enables quick acquisition of data and representation of such diagrams for other R/H ratia.

	Cylindrical vessel with ellipsoidal heads, $\gamma = 3$, $R/h = 87$				
$\hat{ heta}$	\hat{x}	$\frac{N_{\rm t}}{p \cdot R}$	$\frac{N_{\rm s}}{p \cdot R}$	$\frac{M_{\rm t}}{p \cdot R \cdot h}$	$\frac{M_{\rm s}}{p \cdot R \cdot h}$
0	11.59	1.00	1.00	0	0
$\pi/32$	9.950	0.9574	0.9859	0.000	0.000
$\pi/16$	8.417	0.8391	0.9474	0.000	0.000
$3\pi/32$	7.063	0.6682	0.8934	0.000	0.000
$\pi/8$	5.910	0.4699	0.8335	0.000	0.000
$5\pi/32$	4.947	0.2599	0.7746	0.000	0.0021
$3\pi/16$	4.145	0.0448	0.7206	0.0012	0.0040
$7 \pi/32$	3.472	-0.1689	0.6731	0.0009	0.0031
$\pi/4$	2.901	-0.3696	0.6325	-0.0012	-0.0040
9π/32	2.408	-0.5412	0.5984	-0.0055	-0.0182
$5\pi/16$	1.975	-0.6672	0.5704	-0.0116	-0.0386
$11\pi/32$	1.588	-0.7338	0.5477	-0.0186	-0.0618
$3\pi/8$	1.234	-0.7309	0.530	-0.0249	-0.0831
$13\pi/32$	0.9056	-0.6531	0.5166	-0.0289	-0.0963
$7\pi/16$	0.5946	-0.4996	0.5073	-0.0281	-0.0935
$15\pi/32$	0.2946	-0.2764	0.5018	-0.0196	-0.0654
$\pi/2$	0	0	0.5	0	0

Table 1. Values of internal forces and moments in the ellipsoidal head of cylindrical pressure vessel

The described method for calculating strength, performed by a computer programme, gives data on critical stress in the form (25), in a short period of time, for arbitrarily chosen geometry parameters of ellipsoidal head, which makes this method very suitable for designing purposes.

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Numerical and Experimental Analysis of a Shaft Bow Influence on a Rotor to Stator Contact Dynamics

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The shaft bow problem presents a real situation especially in case of slender rotors. This paper investigates the shaft bow influence on the rotor-stator contact dynamics. For this purpose the rotor is described as a simple Jeffcott model and the stator as an elastically suspended rigid ring. To test the numerical model, except a usual run down analysis, an emergency shut down after the sudden rotor unbalance increase is also analyzed. Numerical integration is carried out by the fourth-order Runge-Kutta method. Two different normal force models for the rotor-stator interaction are analyzed. For analyzed parameters, both linear and nonlinear (Hunt and Crossley), normal force models gave similar rotor and stator responses. To confirm some of the results and to tune the numerical model, the experimental investigation on the test rig was conducted.

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Keywords: rotor-stator contact, shaft bow, numerical simulations, rotor vibration, vibration measurements

0 INTRODUCTION

Contact or rub between rotor and stator is one of the most intensive research subjects in rotor dynamics. Most papers which are considering the rotor-stator contact phenomena can be classified in rigid rotor disc-rigid stator contact [1] to [3], bladed disc-stator contacts [4] and rotor-stator contact in retainer bearings when the rotor is supported by active magnetic bearings [5] and [6]. Choy and Padovan [1] developed a general analytical rub model using the following assumptions: simple Jeffcott rotor model, linear stiffness and damping characteristics, rigid casing supported by springs acting in the radial direction, mass inertia of the casing small enough to be neglected, simple Coulomb friction and onset of rub caused by unbalance. Bartha [2] performed an extensive numerical and experimental research of the backward whirl of rotors considering the rigid and elastically suspended stator. Von Grol and Ewins compared measurements and simulations for a windmilling imbalance in aero-engines which is very similar to the classical rotor to stator contact. They have revealed the rotor response rich in different subharmonics. Influence of torsion on the rotor-stator contact was introduced by Edwards et al. [7]. Attention was paid to the effects of torsion on a steady state

response of a system experiencing rotor to stator contact. The analyzed Jeffcott rotor model had only 3 degrees of freedom – d.o.f. because the stator was assumed to be rigid. Žigulić et al. [8] considered nonlinear dynamics of a rotor supported by two dry-friction bearings. The rotor was modeled with FEM-Finite element method and the obtained nonlinear system was integrated with Hilbert-Hughes-Taylor – HHT α method.

Karpenko et al. [9] have shown how the preloading of the snubber ring could stabilize the dynamic response of the rotor. Theoretical predictions of the two-degrees-of-freedom Jeffcott rotor model with the preloaded snubber ring subjected to out-of-balance excitation was studied in [10], while an experimental verification with special attention to the analysis of the shaft rotational speed (excitation frequency) as well as different eccentricities and preloading was given in [11].

This paper considers the contact between the rotor, with a bowed shaft, and the stator at the position of rotor disc (seal), after the appearance of sudden unbalance due to blade loss. Linear and nonlinear models of the normal contact force are compared for different shaft bow – mass unbalance combinations. In this study, not only lateral but also torsional d.o.f. are taken into account allowing the additional extension of the

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model with mechanical model of the induction motor. Results of numerical simulations and experiments are presented and compared using the rotor and stator responses in time domain, spectral maps and rotor orbits

1 MODELING OF THE ROTOR-STATOR SYSTEM

A rather simple mechanical model for the rotor-stator interaction description is used. The rotor is described by Jeffcott model with a shaft bow while the stator is modeled as an elastically suspended rigid ring. Except traditional consideration of the rotor lateral d.o.f. (x_r, y_r) , torsional d.o.f. (φ_r) is taken into account allowing an extension of the model with torsional d.o.f. of the induction motor (φ_m) as well as appropriate mechanical model. To keep consistency of the described consideration the stator had additional three d.o.f., two lateral (x_s, y_s) and one torsional $(\varphi_{\rm s})$. Fig. 1 shows the considered rotor-stator system. Next sections briefly discuss the differential equations of motion, the contact model and mechanical model of the induction motor.



Fig.1. Extended Jeffcott rotor model

1.1. Differential equations of motion

The equations of motion for the system described above can be derived using the Lagrange equations and have the following form:

$$\begin{split} m_{\rm r} \dot{x}_{\rm r} + c_{\rm r} \dot{x}_{\rm r} + k_{\rm r} x_{\rm r} &= m_{\rm r} e \dot{\phi}^2 \cos \phi + m_{\rm r} e \dot{\phi} \sin \phi + k_{\rm r} x_{\rm r0} - F_{\rm Cx} \\ m_{\rm r} \ddot{y}_{\rm r} + c_{\rm r} \dot{y}_{\rm r} + k_{\rm r} y_{\rm r} &= m_{\rm r} e \dot{\phi}^2 \sin \phi - m_{\rm r} e \ddot{\phi} \cos \phi + k_{\rm r} y_{\rm r0} - F_{\rm Cy} \\ J_{\rm r} \ddot{\phi}_{\rm r} + c_{\rm rt} \left(\dot{\phi}_{\rm r} - \dot{\phi}_{\rm m} \right) + k_{\rm rt} \left(\phi_{\rm r} - \phi_{\rm m} \right) = M_{\rm fr1} - M_{\rm C} \\ m_{\rm s} \ddot{x}_{\rm s} + c_{\rm s} \dot{x}_{\rm s} + k_{\rm s} x_{\rm s} = F_{\rm Cx} \\ m_{\rm s} \ddot{y}_{\rm s} + c_{\rm s} \dot{y}_{\rm s} + k_{\rm s} y_{\rm s} = F_{\rm Cy} \\ J_{\rm s} \ddot{\phi}_{\rm s} + c_{\rm st} \dot{\phi}_{\rm s} + k_{\rm st} \phi_{\rm s} = M_{\rm C} \\ J_{\rm m} \ddot{\phi}_{\rm m} + c_{\rm rt} \left(\dot{\phi}_{\rm m} - \dot{\phi}_{\rm r} \right) + k_{\rm rt} \left(\phi_{\rm m} - \phi_{\rm r} \right) = M_{\rm fr2} + M_{\rm m} \end{split}$$
(1)

The Equation (1) is taken from literature [12] and presents a further development of the models found in [2], [7] and [13], where m_r, m_s, c , $c_{\rm s}$, k, $k_{\rm s}$ are mass parameters, lateral damping coefficients and stiffnesses of the rotor and stator respectively, $c_{\rm rt}$, $c_{\rm st}$, $k_{\rm rt}$ and $k_{\rm st}$ are torsional damping coefficients and stifnesses of the rotor shaft and stator, e is the rotor mass eccentricity, $J_{\rm r}, J_{\rm s}$ and $J_{\rm m}$ represent mass moments of inertia of the rotor, stator and induction motor respectively. M_{fr1} and M_{fr2} represent the torques of total losses (friction in bearings and losses in motor fan). These torques were identified from the measurements conducted on the test rig. To simplify the already complicated model, rotor and stator are assumed as isotropic regarding lateral stiffness and damping coefficient, what is still common practice in literature [14].

Vector \mathbf{r}_{r0} from the Fig. 2 represents the static equilibrium position of the bowed shaft [13], where x_{r0} and y_{r0} are its X,Y components and are given by

$$x_{r0} = r_{0x} \cos \varphi_r - r_{0y} \sin \varphi_r$$

$$y_{r0} = r_{0x} \sin \varphi_r + r_{0x} \cos \varphi_r$$
(2)

Components r_{0x} and r_{0y} represent the static equilibrium position of the bowed shaft in the rotor fixed x,y,z coordinate system.



Fig.2. Jeffcott rotor with a bowed shaft

Contact forces F_{Cx} and F_{Cy} as well as the contact moment M_C are defined, according to Fig. 3, by the following expression:

$$F_{Cx} = F_{N} \cos \gamma - F_{T} \sin \gamma$$

$$F_{Cy} = F_{N} \sin \gamma + F_{T} \cos \gamma$$

$$M_{C} = F_{T} R$$
(3)

where the normal $F_{\rm N}$ and the tangential $F_{\rm T}$ contact forces have positive value when they are acting on the stator. Fig. 3 presents the geometry of the rotor - stator contact model with force definition.

1.2. Contact Models

Traditionally, for the rotor-stator impact modeling, two different methods have been namely the Newton's restitution applied, coefficient model and the Contact forceindentation model. The first model is based on the simplifying assumption of perfectly rigid bodies. Actual physical objects are compliant and hence the impact duration is strictly greater than zero. This more realistic view of impact phenomena led many researchers to consider the continuousdynamics models of collision where bodies deform during impact and the collision dynamics is treated as a continuous-time dynamic phenomenon. In its general form, the forceindentation relationship looks like [15]:

$$F_{\rm N} = F_{\rm k}\left(\delta\right) + F_{\rm c}\left(\delta,\dot{\delta}\right) + F_{\rm p}\left(\delta,\dot{\delta}\right) \tag{4}$$

where, $F_{\rm k}$ is the elastic (conservative) part of the normal contact force $F_{\rm N}$, $F_{\rm c}$ the viscous damping part and $F_{\rm p}$ the dissipative part due to plastic deformation. In the paper, plastic dissipations have been neglected. Fig. 4 shows the normal



Fig.3. Geometry of the rotor - stator contact model

contact force model presented as a spring-damper system.

In the development of the elastic relation $F_k(\delta)$, the Hertz theory has to be mentioned. A very commonly used expression is the forceindentation relation for sphere to sphere contact, according to Hertz:

$$F_{\rm k} = k_{\rm s} \delta^{3/2} \tag{5}$$

where, F_k is the normal force pressing the solids together, δ represents the approach of two spheres, i.e. total deformation of both surfaces while k_s is a constant depending on the sphere radii and elastic properties of the sphere materials.



Fig.4. Normal contact force model – a springdamper system

For the case of a cylindrical joint, Rivin [16] based his analysis on a half sine distribution of contact pressure between shaft and bearing. To calculate the pressure distribution, he considered two possible models for the force – deformation relation, linear and quadratic:

$$F_{\rm k} = k_1 \delta \tag{6}$$

$$F_{\rm k} = k_2 \delta^2 \tag{7}$$

Results of his experiments carried out with cast iron and hardened steel showed that the quadratic model offered a better correlation of the force deflection in a journal bearing when a high load per unit length was applied. For low loads, the linear model represented a better fit.

The simplest model of viscous dissipation $F_{\rm c}(\delta, \dot{\delta})$ which is known from the literature [1] is a linear damper,

$$F_{\rm c} = c\delta \tag{8}$$

where *c* is the viscous damping coefficient.

Non-linear expression for visco-elastic force, originally proposed by Hunt and Crossley [17], is given by

$$F_{\rm c} = \lambda \delta^n \dot{\delta} \tag{9}$$

where λ and *n* are model parameters. Fig. 5 shows responses of two analyzed models of the contact force regarding different combination of the elastic force models F_k and damping force models F_c .

The area bounded by hysteresis loop represents a loss of work during an impact of rotor and stator and can be expressed by,

$$W = \oint F_c \, \mathrm{d} \, x \tag{10}$$

The linear contact force model is described by the well known linear equation:

$$m\ddot{\delta} + c\dot{\delta} + k\delta = 0 \tag{11}$$

where *c* is the viscous damping coefficient, *k* is the stiffness, $\delta = |\mathbf{r}| - C_r$ is the local deformation of the rotor and stator in the contact point in normal direction, $\dot{\delta}$ is the relative velocity of the rotor and stator along normal direction during the contact, $|\mathbf{r}| = \sqrt{(x_r - x_s)^2 + (y_r - y_s)^2}$ is the magnitude of vector difference between rotor and stator displacement according to Fig. 3 and C_r is the radial clearance between rotor and stator.



Fig.5. Models of normal contact force

In order to eliminate the negative forces that appear in the linear model and with the improvement in the sense of avoiding the discontinuity, Hunt and Crossley formulated the equation:

$$m\ddot{\delta} + \left(\lambda\delta^{n}\right)\dot{\delta} + k\delta^{n} = 0 \tag{12}$$

Their theory gives a relationship between restitution coefficient and dissipation in contacts via assumption of linear dependency of the restitution coefficient ε and initial impact velocity v_0 i.e.

$$\varepsilon = 1 - \alpha v_0 \tag{13}$$

Thus, according to [17] Equation (12) can be written in the following form

$$m\ddot{\delta} + k\delta^n \left(\frac{3}{2}\alpha\dot{\delta} + 1\right) = 0 \tag{14}$$

where α is the constant dependent mainly on material and geometry of bodies in contact. Based on Equation (14) two separate parts of normal contact force can be identified,

$$F_{\rm N} = F_{\rm c} + F_{\rm k} = \frac{3}{2}\alpha k\delta^{\rm n}\dot{\delta} + k\delta^{\rm n}$$
(15)

According to [16], index n can vary between 1 and 2 for cylindrical joints, so the choice of n = 3/2, originally proposed by Hertz (see Equation (5)), seems fairly reasonable.

In this paper both normal force models (linear–Equations (6), (8) and nonlinear–Equation (15)) have been used and compared.

Tangential contact force has been represented by the Coulomb dry friction law. The definition of the contact model described above can be expressed as follows

$$F_{\rm N} = \begin{cases} 0 & \text{for } r < C_{\rm r} \\ F_{\rm k} + F_{\rm c} & \text{for } r > C_{\rm r} \end{cases}$$
$$F_{\rm T} = \mu_{\rm c} F_{\rm N} \operatorname{sgn}(v_{\rm sl})$$
(16)

where μ_c is the coefficient of friction and v_{sl} is the sliding velocity between rotor and stator surfaces in contact. The angle γ from equation (3) represents the normal force angle or direction of the smallest gap when rotor and stator are not in contact i.e.

$$\tan \gamma = \frac{y_{\rm r} - y_{\rm s}}{x_{\rm r} - x_{\rm s}} \tag{17}$$

To determine $\dot{\delta}$ and v_{sl} further vector calculation should be taken. Fig. 6 shows the velocity definition for the rotor and stator in contact. If the normal direction is known,

$$n_0 = \frac{r}{|r|}, \quad r = [r_x, r_y] = [x_r - x_s, y_r - y_s]$$
 (18)

as well as the velocities of rotor and stator centers $v_r = [v_{xr}, v_{yr}]$ and $v_s = [v_{xs}, v_{ys}]$, one can easily obtain the relative normal rotor-to-stator velocity,

$$\delta = v_{\rm rn} - v_{\rm sn} \tag{19}$$

where $v_{\rm m} = v_{\rm r} \cdot \boldsymbol{n}_0$ and $v_{\rm sn} = v_{\rm s} \cdot \boldsymbol{n}_0$ represent the rotor and stator velocities in normal direction.



Fig.6. Velocity definition for rotor and stator in contact

The sliding velocity v_{sl} can be further defined as,

$$v_{\rm sl} = v_{\rm rt} - v_{\rm st} \tag{20}$$

where $v_{\rm rt} = \mathbf{v}_{\mathbf{r}} \mathbf{t}_0 + \dot{\varphi}_{\rm r} R$

and $v_{st} = \mathbf{v}_s \mathbf{t}_0 + \dot{\phi}_s (R + C_r)$ represent the rotor and stator velocities on its periphery in tangential direction. *R* is the radius of the rotor while $\dot{\phi}_r$, $\dot{\phi}_s$ are angular velocities of the rotor and stator.

1.3. Mechanical Model of the Induction Motor

The driving moment of the induction electric motor has been modeled according to the Kloss expression [18] and [19]

$$M_{\rm m} = M_{\rm max} \frac{2 + 2s_{\rm mn}}{\frac{s\varphi}{s_{\rm mn}} + \frac{s_{\rm mn}}{s\varphi} + 2s_{\rm mn}}$$
(21)

where M_{max} and s_{mn} are the maximum motor torque or breakdown torque and slip for maximum motor torque, s is the actual slip of motors rotor with respect to the synchronous electromagnetic field and $\varphi = f_{\text{sf}}/f_{\text{s}}$ is the ratio of actual frequency given from frequency inverter and nominal frequency given by electrical network (f_{s} =50 Hz).

If we replace M_{max} in Equation (21) with $M_{\text{max b}} = -M_{\text{max}} (1+s_{\text{mn}})/(1-s_{\text{mn}})$, the Kloss expression for regenerative braking is obtained. The regenerative braking starts when actual speed of the motor's rotor is greater than the synchronous speed. The motor is then behaving as an electromagnetic brake. Fig. 7 shows the speed torque characteristics of an induction motor.



Fig.7. Speed torque characteristics of an induction motor

2 NUMERICAL SIMULATIONS

The numerical integration has been carried out by the fourth-order Runge-Kutta method. Time step of $\Delta t = 2 \cdot 10^{-5}$ s, has been applied in all simulations where the contact between rotor and stator was reasonably expected. Otherwise $\Delta t = 1 \cdot 10^{-4}$ s has been used. In some earlier simulations a greater time step was used but it turned out that it couldn't describe some specific rotor behaviors. Time step during one simulation has been fixed, so the subsequent frequency analysis could be done.

The parameters used in all calculations are as follows: rotor mass $m_r = 4.25$ kg, stator mass $m_s = 3.838$ kg, mass moments of inertia $J_r =$ $4.532 \cdot 10^{-3}$ kgm², $J_m = 2.115 \cdot 10^{-3}$ kgm², $J_s =$ $2.11 \cdot 10^{-2}$ kgm², rotor disc radius R = 0.06 m, stiffnesses $k_r = 131540$ N/m, $k_{rt} = 1080$ Nm/rad, $k_{\rm s} = 1.237 \cdot 10^6$ N/m, $k_{\rm st} = 7917$ Nm/rad, damping coefficients $c_{\rm r} = 22.43$ Ns/m, $c_{\rm rt} = 0.0374$ Nms/rad, $c_{\rm s} = 11.96$ Ns/m, $c_{\rm st} = 0.0258$ Nms/rad, radial clearance $C_{\rm r} = 0.4$ mm, motor breakdown torque $M_{\rm max} = 10.25$ Nm and motor speed of max. breakdown torque (at power network frequency $f_{\rm s}$ = 50 Hz) $n_{\rm m} = 2400$ rpm.

The identified dependence between angular deceleration and angular velocity of the rotor, obtained from free run down tests (induction motor switched off), had the following form

$$\varepsilon = -4 \cdot 10^{-6} \omega^3 + 0.002 \omega^2 - 0.5376 \omega - 10.513$$

rad/s² (22)

2.1. Numerical Simulation of Motor Controlled Run-Down (RD) Analysis

The numerical RD analysis has been performed to identify the basic dynamic characteristics of the rotor-stator system and to correlate them with the experimental results. Knowing, from the modal testing (rotor response on the impulse force excitation at standstill) that the first rotor natural bending frequency is equal to $f_{r1} = 28.0$ Hz and according to the Jeffcott (or Laval) theory it is the only natural frequency, RD analysis has been focused on speed range n = 40to 0 Hz. The second rotor natural bending frequency obtained by the modal testing has been $f_{r2} = 148.5$ Hz hence it has been far away enough from the first natural frequency and its influence has been negligible.

The basic goal of this part of numerical analysis was to discover the appropriate relation of the mass unbalance eccentricity vector e and the static equilibrium position of the bowed shaft \mathbf{r}_{r0} according to the measured maximum lateral displacement amplitude of the rotor disc and the qualitative shape of the rotor lateral response in time domain, Fig. 8. Final parameters have been e = 0.046 mm, $r_{\rm r0}$ = 0.05 mm and phase lag of the e in relation to \mathbf{r}_{r0} i.e. $\alpha_0 = \pi$ rad (see Fig. 2). RD analysis has been controlled by a linear frequency (speed law) of the synchronous ramp electromagnetic field of induction motor, $f_{\rm sf} = 40$ to 0 Hz in the time period of 34 s, while its rotor behaved according to the presented mechanical model. At the beginning of the numerical simulation, 2 s have been taken additionally for disappearance of any transients caused by initial conditions.



Fig.8. Simulated lateral rotor response in horizont direction for RD analysis and speed law

Fig. 9 shows the spectral map of the lateral rotor response in the horizontal direction where as expected only the first order can be seen. On the same figure at the beginning of the simulation, the rotor critical speed is excited but only because of the numerical transient due to initial conditions.



Fig.9. Spectral map of the lateral rotor response in horizontal direction x_r

2.2. Numerical Simulations of Sudden Rotor Unbalance Increase (SRUI) Analysis

The basic assumption has been that the rotor is running at a nominal speed above the first rotor critical speed and that the rotor is well balanced. To achieve the initial condition which is similar to the steadily running speed and to avoid any transients, the sudden unbalance was introduced 1 second after the start of simulation. After sudden appearance of the additional unbalance m_{add} it has been assumed that the

permanent monitoring system of the turbo machine has recorded unallowable vibrations and as a precaution measure it starts an immediate shut down procedure. This was accomplished by switching off the induction motor. Then the whole rotor starts to decelerate due to given torques of losses, expressed by the identified rotor angular deceleration mentioned above. The rotor is then approaching to its critical speed and because of additional unbalance m_{uadd} it contacts the stator. When the additional unbalance is big enough, the rotor contacts the stator immediately after appearance of the additional unbalance i.e. before switching off the motor. The time period for response of the shut down procedure after sudden unbalance appearance has been set to 1 s. Two different contact normal force models have been compared i.e. linear and nonlinear. Linear model has contact stiffness $k_{\rm C} = 7.7 \cdot 10^7$ N/m and contact damping $c_{\rm C} = 3.2 \cdot 10^3$ Ns/m and nonlinear model has contact toughness $k_{\rm C} = 1.8 \cdot 10^8$ Pa/m^{1/2} and parameter $\alpha = 3$ s/m. The output inverter frequency before unbalance appearance was set to $n_{\rm sf} = 40$ Hz. It turns out that the frictional losses in test rig are quite big so the rotor decelerates, with small unbalance and without contact, in some 5 to 6 seconds from 40 Hz to standstill.

In Figs. 10 to 12 the rotor response without contact to stator is presented, without any additional unbalance, $m_{add} = 0$ (or presented via rotor mass eccentricity $e_{add} = 0$). If someone compares Fig. 9 and Fig. 12, except differences in the shape of first order (linear in Fig. 9 and nonlinear in Fig. 12), on Fig. 12 exists additional horizontal line at 28 Hz. The reason for this mainly lies in fast rotor speed change while passing though its critical speed.



Fig.10. Rotor center orbit during free RD (red dashed line represents a clearance circle)



Fig.11. Simulated lateral rotor response in horizontal direction for free RD analysis with speed law $e_{add} = 0$



Fig.12. Spectral map of rotor displacements in horizontal direction, free RD, $e_{add} = 0$



Fig.13. Rotor center orbit during RD, rotorstator contact appeared, $e_{add} = 0.124$ mm



Fig.14. Lateral rotor and stator responses in horizontal direction, contact appeared, $e_{add} = 1.24 \cdot 10^{-4}$ m, $r_{r0} = 0.05$ mm, $\alpha_0 = \pi$ rad



Fig.15. Rotor speed laws for simulations without and with contact



Fig.16. Spectral map of rotor displacement in horizontal direction, contact appeared



Fig.17. Spectral map of rotor relative torsional angles $\varphi_m - \varphi_r$, contact appeared



Fig.18. Spectral map of stator lateral velocities in horizontal direction, contact appeared



Fig.19. Spectral map of stator torsional angles, contact appeared

Figures 13 to 19 show rotor and stator response after the appearance of the additional mass unbalance (presented in rotor mass eccentricity) $e_{add} = 0.124$ mm, so the total rotor mass eccentricity is equal to $e = e_{in} + e_{add} = 1.17$ mm. The magnitude of the vector \mathbf{r}_{r0} was $r_{r0} = 0.05$ mm and the phase lag was $a_0 = \pi$ rad. In this simulation the linear model of normal contact force has been applied and coefficient of friction in contact was $\mu_C = 0.18$, according to Bartha [2].

In Fig. 11 we had the supercritical selfbalancing effect because the mass eccentricity ewas smaller than the magnitude of the shaft bow vector r_{r0} . On the contrary in Fig. 14 the self balancing effect is placed subcritically at the $t_1 =$ 3.8 s of simulation because after sudden unbalance increase, the total mass unbalance eccentricity e became greater than r_{r0} .

Fig. 15 shows the difference between speed laws for simulations with and without the rotor-stator contact appearance. As expected a rotor deceleration for simulation where the rotorstator contact appeared, was more intensive especially in a time period with established rotorstator contact (see contact indication in Fig. 15). After introduction of SRUI, the rotor made several impacts to the stator and then a permanent contact was established. In Figs. 16 to 19 there is a narrow time period around $t_1 = 1$ s of simulation where multiple harmonics can be seen. This is due to intermittent rotor-stator contact. It is interesting to see after the rotor-stator separation $(t_1 = 2.75 \text{ s})$ that the rotor and stator continued to vibrate with their own flexural natural frequencies, $f_r = 28$ Hz (Fig. 16), $f_s = 90$ Hz (Fig. 18) and torsional frequencies $f_{\rm rt} = 137.9$ Hz (Fig. 17), $f_{st} = 97.7$ Hz (Fig. 19).

In Fig. 20, the influence of the phase lag of the mass unbalance eccentricity vector **e** in relation to the static equilibrium position of the bowed shaft \mathbf{r}_{r0} i.e. α_0 has been analyzed. Presented contact forces are related to $\alpha_0 = 180^\circ$, 135° and 90° . In the same time normal contact forces of the linear and nonlinear models have been compared, so the diagrams in the left column correspond to the linear and in the right column to the nonlinear model of the normal contact force.

The main difference between the linear and nonlinear model of the contact normal force, presented in Fig. 20, can be seen in the region immediately after appearance of SRUI. Nonlinear model generally had a greater force response for the first impact, while in permanent contact both linear and nonlinear model had almost the same values. Details about the differences between analyzed normal force models and influence of the phase lag α_0 are presented in Table 1.

The normal force for the permanent rotorstator contact (with sliding) F_{Cperm} has showed expected increasing trend with decreasing phase lag α_0 , because α_0 smaller than 180° means greater

Table 1. Results of contact normal force analysis regarding phase lag of the mass eccentricity vector e in relation to shaft bow radius vector r_{r0}

α_0 ,	$F_{C \mathrm{max}}/F_{C \mathrm{perm}}$, N		
	Linear	Nonlinear	
180°	327.7 / 96.8	611.9 / 97.1	
135°	279.8 / 99.3	462.2 / 99.5	
90°	203.7 / 104.8	277.1 / 105.0	

unbalance excitation. The reason for contradictory trend of decreasing of the F_{Cmax} (max. value of the normal force for the first impact) with decreasing α_0 (increasing unbalance excitation) lies in the fact that α_0 of the e_{in} and e_{add} is the same. This caused an increasing tendency of the radius of initial rotor orbit and therefore the decrease of the maximum normal velocity at the first rotor-stator impact.

3 EXPERIMENTAL ANALYSIS

The analysis has been performed on the test rig which can be seen on Fig. 21. The test rig has been specifically designed for the rotor-stator contact measurements. It consists of the following main components: robust foundation, mounting plate, roller bearings with their supports, rotor, stator, flexible coupling, induction motor with speed controller, measuring system based on Bruel and Kjaer non-contacting displacement sensors and National Instruments data acquisition card PCI NI 4472 with adequate software (LabVIEW, Matlab). Measurements showed that the first critical speed of the rotor was 28 Hz, while the first natural frequency of the stator was 90 Hz. Radial clearance between the rotor and stator was 0.4 mm. Before the experiments the rotor had been balanced.



Fig. 20. Normal forces in contact between rotor and stator; left column – linear model, right column nonlinear model; a) $\alpha_0 = 180^\circ$, b) $\alpha_0 = 135^\circ$, c) $\alpha_0 = 90^\circ$

Also before every single experiment the contact surfaces of the rotor and stator were lubricated with WD40 spray [2] to decrease the coefficient of friction to the value of approximately $\mu = 0.18$ and to minimize the possibility of destruction of sensors and other structural parts of the test rig.

3.1 Experimental, Motor Controlled, Run-Down (RD) Analysis

To verify the numerical results presented in chapter 2.1, an experimental RD analysis has been performed.



Fig.21. Test rig for rotor-stator contact investigation

The speed change has been controlled via frequency inverter Micromaster 440 from Siemens. The speed rate has been the same as in the numerical simulation only the starting speed has been 70 Hz. Figs. 22 and 23 show measured lateral responses of the rotor in horizontal x direction with speed law and spectral map of the rotor lateral response in the horizontal direction. This figures can be compared to Figs. 8 and 9 respectively, for a speed range n = 40 to 0 Hz. In Fig. 23 except for the first harmonic, other higher harmonics can be also seen due to various imperfections of the experimental model like radial and angular misalignment etc. Fig. 24 shows a spectral map of stator displacements in horizontal direction. Although there wasn't direct contact between the rotor and stator, the stator was excited as it can be seen on Fig. 24. This is due to the fact that rotor bearing supports and stator supports are rigidly connected on the mounting plate, so this situation allows an indirect stator excitation through the mounting plate.

Measured flexural natural frequencies were, $f_r = 28$ Hz (Fig. 23), $f_s = 90$ Hz (Fig. 24) and stator torsional frequency $f_{st} = 102.5$ Hz (Fig. 24).

To excite effectively both stator natural frequencies (with second harmonic), the 70 Hz as a starting speed of this RD analysis was chosen.



Fig.22. Measured lateral rotor x_r response in horizontal direction and speed law



Fig.23. Measured spectral map of rotor displacement in horizontal direction



Fig.24. Measured spectral map of stator displacement in horizontal direction

3.2 Experimental Sudden Rotor Unbalance Increase (SRUI) Analysis

Because this test rig couldn't simulate the SRUI scenario during the rotor operation, additional mass has been added instead of being removed but with rotor at standstill. The rotor has been quickly accelerated to rotational speed of 40 Hz to reduce the rubbing while passing through rotor critical speed. The power supply to the motor was then switched off and a whole rotor started a free deceleration. Two measurements have been presented: free rotor RD with good rotor balance condition as in previous analysis (without appearance of the rotor-stator contact), Figs. 25 and 26 and free RD of the rotor with additional mass unbalance (presented in the rotor mass eccentricity) $e_{add} = 0.124$ mm, Figs. 27 to 30.

The first measurement (Figs. 25 and 26) can be compared directly with the numerical simulation presented in the Figs. 11 and 12. Looking just on Figs. 11 and 25 a great correspondence can be noticed. Comparing spectral maps of simulated and measured lateral rotor displacements, i.e. Figs. 12 and 26, some differences regarding presence of higher harmonics in measured response can be observed. The reason for their presence is explained in the previous chapter and their influence is reflected on extended life of rotor bending natural frequency present in the response till the rotor standstill.

The second measurement (Figs. 27 to 30) can be compared with the numerical simulation presented in Figs. 13 to 19 with remark that in simulation the motor is switched off in $t_1 = 2$ s (Fig. 14) while in measurement in $t_1 = 1.65$ s. Furthermore, rotor and stator were, in measurement, in permanent contact from the very beginning, while in numerical SRUI simulation they got into contact after additional mass unbalance activation i.e. in $t_1 = 1.0$ s of simulation.

Taking into account the aforementioned reasons for differences between simulation and measurement, a great correspondence between Figs. 14 and 27 can be noticed again. In measured spectral map of the rotor lateral response, Fig. 28, vibrations have strong first harmonic component i.e. $1 \times$ rotating frequency, while there are higher harmonics too. They were present also in the

rotor lateral response without appearance of the rotor-stator contact (Fig. 26), because of the above mentioned imperfections of the rotor experimental model. Both lateral and torsional natural frequencies of the stator are present in each of its depicted responses, i.e. in the lateral response (Fig. 29) and in the torsional response (Fig. 30), after the final rotor-stator separation. This happened mainly because of the vicinity of stator natural frequencies and due to indirect measurement of torsional stator rotations via 2 parallel non-contacting displacement sensors in horizontal (lateral) direction. Although both stator frequencies can be seen in two separated responses, lateral frequency is more emphasized in lateral response while torsional frequency is more emphasized in its torsional response. Thus in the case of torsionally elastic systems susceptible to the rotor-stator contact it is important to combine the measurements with results of the numerical simulations to make a decision of the character of a particular natural frequency.

4 CONCLUSIONS

The presented numerical model of the rotor with a bowed shaft - stator system has proved to be capable for analyzing different contact situations. Linear and nonlinear normal contact force models were compared in the same simulated situations and they showed similar responses for permanent rotor-stator contact, while nonlinear force model generally had a greater response for the first rotor-stator impact. This didn't have a big influence on a subsequent rotor and stator response and their stability, for the analyzed parameters. Both normal force models, for permanent rotor-stator contact (with sliding) F_{Cperm} , showed increasing tendency with decreasing phase lag α_0 , because if α_0 is smaller than 180° the unbalance excitation will be greater.

Experimental analysis showed more realistic rotor and stator responses with more different harmonics present than in the results of numerical simulations. This was mainly due to tested model imperfections and measuring principle. In the real practice measured signals are even more unclear and mixed with different noise, so it is important to combine the measurements with adequate numerical simulations to recognize the needed information.



Fig.25. Measured lateral rotor x_r response in horizontal direction with speed law, contact-free situation, free RD, $e_{add} = 0$ mm



Fig.26. Spectral map of rotor displacements in horizontal direction, free RD, $e_{add} = 0$ mm



Fig.27. Measured lateral rotor x_r and stator x_s responses in horizontal direction, free RD, contact appeared, $e_{add} = 0.124$ mm



Fig.28. Spectral map of rotor lateral displacements in horizontal direction, free RD, contact appeared, $e_{add} = 0.124$ mm



Fig.29. Measured spectral map of stator lateral displacement in horizontal direction, free RD, $e_{add} = 0.124 \text{ mm}$



Fig.30. Measured spectral map of stator torsional rotations, free RD, $e_{add} = 0.124$ mm

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MINLP Optimization of a Single-Storey Industrial Steel Building

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The paper presents the topology and standard sizes optimization of a single-storey industrial steel building, made from standard hot rolled I sections. The structure consists of main portal frames, connected with purlins. The structural optimization is performed by the Mixed-Integer Non-linear programming approach (MINLP). The MINLP performs a discrete topology and standard dimension optimization simultaneously with continuous parameters. Since the discrete/continuous optimization problem of the industrial building is non-convex and highly non-linear, the Modified Outer-Approximation/Equality-Relaxation (OA/ER) algorithm has been used for the optimization. Alongside the optimum structure mass, the optimum topology with the optimum number of portal frames and purlins as well as all standard cross-section sizes have been obtained. The paper includes the theoretical basis and a practical example with the results of the optimization.

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Keywords: industrial buildings, topology optimization, sizing optimization, non-linear programming, MINLP

0 INTRODUCTION

Single-storey frame structures are extensively used for industrial, leisure and commercial buildings. In order to obtain efficient frame designs, researchers have introduced various optimization techniques, suitable either for continuous or discrete optimization. O'Brien and Dixon [1] have proposed a linear programming approach for the optimum design of pitched roof frames. Guerlement et al. [2] have introduced a practical method for single-storey steel structures, based on a discrete minimum weight design and Eurocode 3 [3] design constraints. Recently, Saka [4] has considered an optimum design of pitched roof steel frames with haunched rafters by using a genetic algorithm. One of the latest researches reported in this field is the work of Hernández et al. [5], where the authors have considered a minimum weight design of the steel portal frames with software developed for the structural optimization. It should be noted that all the mentioned authors deal with the discrete sizes optimization only at fixed structural topologies.

This paper discusses the simultaneous topology, standard sizes and continuous parameter optimization of an unbraced singlestorey industrial steel building. The optimization of the portal frames and purlins was performed by the Mixed-Integer Non-linear Programming approach (MINLP). The MINLP is a combined discrete and continuous optimization technique. In this way, the MINLP performs the discrete topology (i.e. the number of frames and purlins) and the standard dimension (i.e. the standard cross-section sizes of the columns, beams and purlins) optimization simultaneously with the continuous optimization of the parameters (e.g. the structure mass, internal forces, deflections, etc.).

The MINLP discrete/continuous optimization problems of frame structures are in most cases comprehensive, non-convex and highly non-linear. The optimization is proposed to be performed through three steps. The first one includes the generation of a mechanical superstructure of different topology and standard dimension alternatives, the second one involves the development of an MINLP model formulation and the last one consists of a solution for the defined MINLP optimization problem.

The objective of the optimization is to minimize the mass of the single-storey industrial building. The mass objective function is subjected to the set of equality and inequality constraints

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known from the structural analysis and dimensioning. The dimensioning of steel members is performed in accordance with Eurocode 3.

The Modified Outer-Approximation /Equality-Relaxation algorithm is used to perform the optimization, see Kravanja and Grossmann [6], Kravanja et al. [7] and [8]. The two-phase MINLP optimization is proposed. It starts with the topology optimization, while the standard dimensions temporarily relaxed are into continuous parameters. When the optimum topology is found, the standard dimensions of the cross-sections are reestablished and the simultaneous discrete topology and standard dimension optimization of the beams, columns and purlins is then continued until the optimum solution is found.

1 SINGLE-STOREY INDUSTRIAL BUILDING

The paper presents the topology and standard sizes optimization of unbraced rigid single-storey industrial building steel structures, Fig. 1. The columns, beams and purlins are proposed to be built up of standard hot rolled steel I sections.

The considered portal frame structures are optimized under the combined effects of the self-

weight of the frame members, a uniformly distributed surface variable load (snow and wind), a concentrated horizontal variable load (wind) and an initial frame imperfection. The purlins are designed to transfer the permanent load (the selfweight of the purlins and the weight of the roof) and the variable load (snow and wind). The internal forces are calculated by the elastic firstorder method. The dimensioning of the steel members is performed in accordance with Eurocode 3 for the conditions of both the ultimate limit state (ULS) and the serviceability limit state (SLS).

When the ULS is considered, the elements are checked for the axial, shear and bending moment resistance, for the interaction between the bending moment and the axial force, the interaction between the axial compression/buckling and the buckling resistance moment.

The total deflection δ_{max} subject to the overall load and the deflections δ_2 subjected to the variable imposed load are calculated to be smaller than the limited maximum values: span/200 and span/250, respectively. The frame horizontal deflections are also checked for the recommended limits: the relative horizontal deflection of the portal frame should be smaller then the frame height/150.



Fig. 1. Single-storey industrial steel building

2 MINLP MODEL FORMULATION FOR MECHANICAL SUPERSTRUCTURES

It is assumed that a non-convex and nonlinear discrete/continuous optimization problem can be formulated as a general MINLP problem (MINLP-G) in the form:

min
$$z = \mathbf{c}^{\mathrm{T}}\mathbf{y} + \mathbf{f}(\mathbf{x})$$

s.t. $h(\mathbf{x}) = \theta$
 $g(\mathbf{x}) \le \theta$ (MINLP-G)
 $B\mathbf{y} + C\mathbf{x} \le b$
 $\mathbf{x} \in X = \{\mathbf{x} \in \mathbb{R}^{\mathrm{n}} : \mathbf{x}^{\mathrm{LO}} \le \mathbf{x} \le \mathbf{x}^{\mathrm{UP}}\}$
 $\mathbf{y} \in Y = \{0, 1\}^{\mathrm{m}}$

where x is a vector of continuous variables specified in the compact set X and y is a vector of discrete, mostly binary 0-1 variables. Functions f(x), h(x) and g(x) are non-linear functions involved in the objective function z, the equality and inequality constraints, respectively. All functions f(x), h(x) and g(x) must be continuous and differentiable. All functions f(x), h(x) and g(x) must be continuous and differentiable. Finally, $By+Cx \le b$ represents a subset of mixed linear equality/inequality constraints.

The above general MINLP-G model formulation has been adapted for the optimization of mechanical superstructures. The resulting MINLP formulation for mechanical superstructures (MINLP-MS) that is more specific, particularly in variables and constraints, can be used also for the modelling the steel industrial buildings. It is given in the following form:

min
$$z = c^{T} y + f(x)$$

s.t. $h(x) = 0$
 $g(x) \le 0$
 $A(x) \le a$
 $Ey \le e$ (MINLP-MS)
 $Dy^{e} + R(x) \le r$
 $Ky^{e} + L(d^{en}) \le k$
 $Py + S(d^{st}) \le s$
 $x \in X = \{x \in R^{n}: x^{LO} \le x \le x^{UP}\}$
 $y \in Y = \{0,1\}^{m}$

The MINLP model formulation for mechanical superstructures is proposed to be described as follows:

- Included are continuous variables $x=\{d, p\}$ and discrete binary variables $y=\{y^e, y^{st}\}$. Continuous variables are partitioned into design variables $d=\{d^{en}, d^{st}\}$ and into performance (non-design) variables p, where subvectors d^{en} and d^{st} stand for continuous and standard dimensions, respectively. Subvectors of the binary variables y^e and y^{st} denote the potential existence of structural elements inside the superstructure (the topology determination) and the potential selection of standard dimension alternatives, respectively.
- The mass or economical objective function z involves fixed mass or cost charges in the linear term $c^{T}y$ and dimension dependant mass or costs in the term f(x).
- Parameter non-linear and linear constraints h(x)=0, $g(x) \le 0$ and $A(x) \le a$ represent a rigorous system of the design, loading, resistance, stress, deflection, etc. constraints known from the structural analysis.
- Integer linear constraints $Ey \le e$ are proposed to describe the relations between binary variables.
- Mixed linear constraints $Dy^{e}+R(x) \le r$ restore interconnection relations between currently selected or existing structural elements (corresponding $y^{e}=1$) and cancel relations for currently disappearing or nonexisting elements (corresponding $y^{e}=0$).
- Mixed linear constraints $Ky^{e+}L(d^{en}) \le k$ are proposed to define the continuous design variables for each existing structural element. The space is defined only when the corresponding structure element exists $(y^{e}=1)$, otherwise it is empty.
- Mixed linear constraints $Py+S(d^{st}) \le s$ define standard design variables d^{st} . Each standard dimension d^{st} is determined as a scalar product between its vector of *i*, *i* \in *I*, discrete standard dimension constants $q=\{q_1, q_2, q_3, ..., q_i\}$ and its vector of subjected binary variables $y^{st}=\{y^{st}_1, y^{st}_2, y^{st}_3, ..., y^{st}_i\}$, see Eq. (1). Only one discrete value can be selected for each standard dimension since the sum of the binary variables must be equal to 1 Eq. (2):

$$d^{\rm st} = \sum_{i \in I} q_i y_i^{\rm st} \tag{1}$$

$$\sum_{i \in I} y_i^{st} = 1 \tag{2}$$

3 OPTIMIZATION MODEL FRAMEOPT

The MINLP optimization model FRAMEOPT (FRAME OPTimization) for the optimization of the single storey industrial steel buildings has been developed with relating to the above MINLP model formulation for mechanical structures.

The following assumptions and simplifications have been defined in the model FRAMEOPT and considered in the optimization:

- Considered was a single load case only, where the partial safety factors and combination of actions were defined according to Eurocodes. The optimization of the structure was performed under the combined effects of:
 - the self-weight of the structure (the line uniform load of columns, beams and purlins) and the weight of the roof (the vertical surface load) plus
 - snow and vertical wind (the uniformly distributed vertical surface variable load) plus
 - horizontal wind (the horizontal force at the top of the columns).
- Equal steel portal frames and equal purlins were proposed to compose the structure.
- Steel members were proposed to be made from standard hot rolled European wide flange I sections (HEA sections).
- The global portal frame geometry including the span, height and the beam inclination was proposed to be fix through the optimization.
- Vertical and horizontal bracing systems as well as wall sheeting rails were not included in this calculation/optimization.
- The internal forces and deflections were calculated by the elastic first-order method.
- The portal frames were classified as nonsway steel portal frames. The ratio between the design value of the total vertical load $N_{\rm Sd}$ and the elastic critical value for failure in a sway mode $N_{\rm cr}$ was constrained: $N_{\rm Sd}/N_{\rm cr} \leq 0.1$.
- The portal frame was calculated as a laterally supported frame. Hereby, the steel members

were checked only for the in-plane instability. Columns were designed for the compression/buckling resistance plus the lateral torsional buckling. Beams were checked for the in-plane bending moment resistance.

 Buckling lengths of columns were calculated as the in-plane buckling lengths for the nonsway mode.

As an interface for mathematical modelling and data inputs/outputs GAMS (General Algebraic Modeling System), i.e. a high level language, was used [9]. The proposed optimization model includes the structure's mass objective function, parameter structural nonlinear and linear constraints, integer and mixed integer logical constraints, sets, input data (constants) and variables.

3.1. Mass objective function

The mass objective function of the industrial building structure is defined by Eq. (3). The mass of the structure *MASS* comprises the masses of columns, beams and purlins. A_C , A_B and A_P represent the cross-section areas of the column, beam and purlins, respectively. *h* denotes the height of the column, L_B is the length of the frame beam and L_L is the length of the industrial building (and purlins). *NOFRAME* represents the number of portal frames and *NOPURLIN* denotes the number of purlins. Each portal frame is constructed from two columns and two beams, see Fig. 2.

$$MASS = 2 \cdot (A_{\rm C} \cdot h \cdot \rho) \cdot NOFRAME + 2 \cdot (A_{\rm B} \cdot L_{\rm B} \cdot \rho) \cdot NOFRAME + (A_{\rm p} \cdot L_{\rm L} \cdot \rho) \cdot NOPURLIN$$
(3)

3.2. Parameter structural non-linear and linear constraints

The first constraints of the model represent the constraints (4) to (7) which determine the relations between the continuous cross-sectional dimensions and the cross-sectional height of the column $h_{\rm C}$. These equations accelerate the convergence of the optimization when standard dimensions are re-established. They define the section breadth $b_{\rm C}$, the flange thickness $t_{\rm f,C}$, the



Fig. 2. Portal frame and cross-sections of elements

web thickness $t_{w,C}$ and the cross-section area A_C (see Fig. 2) for the column. The second moments of the area about the y-y and z–z axis, Iy, C and Iz,C, the torsional constant It, C and the warping

constant I ω , C for the frame column are given by Eqs. (8) to (11). Similar cross-sectional constraints are defined for the frame beam, Eqs. (12) to (16), and for the purlins, see Eqs. (17) to (24).

$$b_{\rm C} = -8.7681 \cdot 10^{-12} \cdot h_{\rm C}^7 + 3.5913 \cdot 10^{-9} \cdot h_{\rm C}^6 - 5.9883 \cdot 10^{-7} \cdot h_{\rm C}^5 + 5.1897 \cdot 10^{-6} \cdot h_{\rm C}^4 - 2.4578 \cdot 10^{-3} \cdot h_{\rm C}^3 + 6.007 \cdot 10^{-2} \cdot h_{\rm C}^2 - 5.8757 \cdot h_{\rm C} + 29.294$$
(4)

$$t_{\rm f,C} = 1.5801 \cdot 10^{-8} \cdot h_{\rm C}^4 + 3.4958 \cdot 10^{-6} \cdot h_{\rm C}^3 + 2.3488 \cdot 10^{-4} \cdot h_{\rm C}^2 - 1.9322 \cdot 10^{-3} \cdot h_{\rm C} + 0.76681$$
(5)

$$t_{\rm w,C} = -1.0598 \cdot 10^{-5} \cdot h_{\rm C}^2 + 2.4652 \cdot 10^{-3} \cdot h_{\rm C} + 0.23804 \tag{6}$$

$$A_{\rm c} = 2 \cdot b_{\rm c} \cdot t_{\rm f,c} + \left(h_{\rm c} - 2 \cdot t_{\rm f,c}\right) \cdot t_{\rm w,c} \tag{7}$$

$$I_{y,C} = \frac{2 \cdot b_{C} \cdot t_{f,C}^{3}}{12} + \frac{t_{w,C} \cdot \left(h_{C} - 2 \cdot t_{f,C}\right)^{3}}{12} + 2 \cdot b_{C} \cdot t_{f,C} \cdot \left(\frac{h_{C}}{2} - \frac{t_{f,C}}{2}\right)^{2}$$
(8)

$$I_{z,C} = \frac{2 \cdot t_{f,C} \cdot b_{C}^{3}}{12} + \frac{(h_{C} - 2 \cdot t_{f,C}) \cdot t_{w,C}^{3}}{12}$$
(9)

$$I_{t,C} = \frac{1}{3} \cdot \left(2 \cdot b_{C} \cdot t_{f,C}^{3} \right) + \frac{1}{3} \cdot \left(h_{C} - 2 \cdot t_{f,C} \right) \cdot t_{w,C}^{3}$$
(10)

$$I_{\omega,C} = \frac{I_{z,C}}{4} \cdot \left(h_{C} - 2 \cdot t_{f,C}\right)^{2}$$

$$\tag{11}$$

$$b_{\rm B} = -8.7681 \cdot 10^{-12} \cdot h_{\rm B}^{7} + 3.5913 \cdot 10^{-9} \cdot h_{\rm B}^{6} - 5.9883 \cdot 10^{-7} \cdot h_{\rm B}^{\cdot 5} + 5.1897 \cdot 10^{-6} \cdot h_{\rm B}^{4} - -2.4578 \cdot 10^{-3} \cdot h_{\rm B}^{3} + 6.007 \cdot 10^{-2} \cdot h_{\rm B}^{2} - 5.8757 \cdot h_{\rm B} + 29.294$$
(12)

$$t_{\rm f,B} = 1.5801 \cdot 10^{-8} \cdot h_{\rm B}^4 + 3.4958 \cdot 10^{-6} \cdot h_{\rm B}^3 + 2.3488 \cdot 10^{-4} \cdot h_{\rm B}^2 - 1.9322 \cdot 10^{-3} \cdot h_{\rm B} + 0.76681$$
(13)

$$t_{\rm w,B} = -1.0598 \cdot 10^{-5} \cdot h_{\rm B}^2 + 2.4652 \cdot 10^{-3} \cdot h_{\rm B} + 0.23804 \tag{14}$$

$$A_{\rm B} = 2 \cdot b_{\rm B} \cdot t_{\rm f,B} + (h_{\rm B} - 2 \cdot t_{\rm f,B}) \cdot t_{\rm w,B}$$
(15)

$$I_{y,B} = \frac{2 \cdot b_{B} \cdot t_{f,B}^{3}}{12} + \frac{t_{w,B} \cdot \left(h_{B} - 2 \cdot t_{f,B}\right)^{3}}{12} + 2 \cdot b_{B} \cdot t_{f,B} \cdot \left(\frac{h_{B}}{2} - \frac{t_{f,B}}{2}\right)^{2}$$
(16)

$$b_{\rm p} = -8.7681 \cdot 10^{-12} \cdot h_{\rm p}^7 + 3.5913 \cdot 10^{-9} \cdot h_{\rm p}^6 - 5.9883 \cdot 10^{-7} \cdot h_{\rm p}^{.5} + 5.1897 \cdot 10^{-6} \cdot h_{\rm p}^4 - -2.4578 \cdot 10^{-3} \cdot h_{\rm p}^3 + 6.007 \cdot 10^{-2} \cdot h_{\rm p}^2 - 5.8757 \cdot h_{\rm p} + 29.294$$
(17)

$$t_{\rm f,p} = 1.5801 \cdot 10^{-8} \cdot h_{\rm p}^4 + 3.4958 \cdot 10^{-6} \cdot h_{\rm p}^3 + 2.3488 \cdot 10^{-4} \cdot h_{\rm p}^2 - 1.9322 \cdot 10^{-3} \cdot h_{\rm p} + 0.76681$$
⁽¹⁸⁾

$$t_{\rm w,P} = -1.0598 \cdot 10^{-5} \cdot h_{\rm P}^2 + 2.4652 \cdot 10^{-3} \cdot h_{\rm P} + 0.23804 \tag{19}$$

$$A_{\rm P} = 2 \cdot b_{\rm P} \cdot t_{\rm f,P} + \left(b_{\rm P} - 2 \cdot t_{\rm f,P}\right) \cdot t_{\rm w,P} \tag{20}$$

$$I_{y,P} = \frac{2 \cdot b_{C} \cdot t_{f,P}^{3}}{12} + \frac{t_{w,P} \cdot (h_{P} - 2 \cdot t_{f,P})^{3}}{12} + 2 \cdot b_{P} \cdot t_{f,P} \cdot \left(\frac{h_{P}}{2} - \frac{t_{f,P}}{2}\right)^{2}$$
(21)

$$I_{z,P} = \frac{2 \cdot t_{f,P} \cdot b_{P}^{3}}{12} + \frac{(b_{P} - 2 \cdot t_{f,P}) \cdot t_{w,P}^{3}}{12}$$
(22)

$$I_{t,P} = \frac{1}{3} \cdot \left(2 \cdot b_{P} \cdot t_{t,P}^{3} \right) + \frac{1}{3} \cdot \left(h_{P} - 2 \cdot t_{t,P} \right) \cdot t_{w,P}^{3}$$
(23)

$$I_{\omega,\mathrm{P}} = \frac{I_{z,\mathrm{P}}}{4} \cdot \left(h_{\mathrm{P}} - 2 \cdot t_{\mathrm{f},\mathrm{P}}\right)^2 \tag{24}$$

The length of the frame beam $L_{\rm B}$ is calculated according to Eq. (25) and the angle of the inclination of the beam α is defined by Eq. (26). *L* represents the frame span and *f* denotes the overheight of the frame beam:

$$L_{\rm B} = \sqrt{\left(L/2\right)^2 + f^2}$$
(25)

$$\alpha = \arctan\left(f/(L/2)\right) \tag{26}$$

The uniformly distributed vertical surface variable load q_z , the uniformly distributed horizontal surface variable load q_y , the self-weight per unit length of the portal frame g, the concentrated design horizontal variable wind load P (for the ULS) and wind load P_w (for the SLS) are defined by Eqs. (27) to (31):

$$q_{z} = \left(s \cdot \cos^{2}(\alpha) + w_{v}\right) \cdot e_{f}$$
⁽²⁷⁾

$$q_{\rm y} = s \cdot \cos(\alpha) \cdot \sin(\alpha) \cdot e_{\rm f} \tag{28}$$

$$g = A_{\rm B} \cdot \rho + (A_{\rm P} \cdot \rho \cdot e_{\rm f})/e_{\rm p}$$
⁽²⁹⁾

$$P = \gamma_{\rm q} \cdot w_{\rm h} \cdot e_{\rm f} \cdot h/2 \tag{30}$$

$$P_{\rm w} = w_{\rm h} \cdot e_{\rm f} \cdot h/2 \tag{31}$$

Where s, w_v and w_h represent snow, the vertical and horizontal wind per m² (the variable imposed load); e_f stands for the intermediate distance between the portal frames, ρ is the density of steel, γ_q is the partial safety factor for the variable load and *h* represents the height of the columns. The number of the portal frames *NOFRAME*, the number of purlins *NOPURLIN* and the maximal intermediate distance between the purlins e_p are determined by Eqs. (32) to (38), where L_L represents the length of the industrial building, *MINNO*^{frame} and *MAXNO*^{frame} denote the minimal and maximal number of defined portal frames, and *MINNO*^{purlin} and *MAXNO*^{purlin} stand for the minimal and maximal number of purlins.

$$NOFRAME = L_{\rm L}/e_{\rm f} + 1 \tag{32}$$

$$NOFRAME \ge MINNO^{\text{frame}}$$
 (33)

$$NOFRAME \leq MAXNO^{\text{frame}}$$
 (34)

$$NOPURLIN = 2 \cdot \left(L_{\rm B}/e + 1 \right) \tag{35}$$

$$NOPURLIN \ge MINNO^{\text{purlin}}$$
 (36)

 $NOPURLIN \leq MAXNO^{\text{purlin}}$ (37)

$$e_{\rm p} \le 250 \ [cm] \tag{38}$$

Eqs. (39) to (45) represent the constraints which determine the portal frames to be non-sway frames. The column stiffness coefficient $K_{\rm C}$, the effective beam stiffness coefficient $K_{\rm B}$, the distribution factor of the column for the sway frame $\eta_1^{\rm s}$, and the plane buckling length of the column for a sway frame mode $\beta_{\rm sway}$ are calculated by Eqs. (39) to (42). The value of the distribution factor η_2 is taken to be 1 because of the pinned connection of the columns.

$$K_{\rm c} = \frac{I_{\rm c}}{h} \tag{39}$$

$$K_{\rm B} = \frac{I_{\rm B}}{s} \tag{40}$$

$$\eta_{\rm I}^{\rm s} = \frac{K_{\rm C}}{K_{\rm C} + 1.5 \cdot K_{\rm B}} \tag{41}$$

$$\beta_{\text{sway}} = \sqrt{\frac{1 - 0.2 \cdot (\eta_1^{\text{s}} + \eta_2) - 0.12 \cdot \eta_1^{\text{s}} \cdot \eta_2}{1 - 0.8 \cdot (\eta_1^{\text{s}} + \eta_2) + 0.6 \cdot \eta_{1\,1}^{\text{s}} \cdot \eta_2}}$$
(42)

Eq. (43) represents the elastic critical load ratio (N_{sd}/N_{cr}) which defines the steel portal frame to be a non-sway frame. The distribution factor of the column for the non-sway frame η_1^{NS} and the plane buckling length of the column for the nonsway frame mode $\beta_{non-sway}$ are given by Eqs. (44) to (45):

$$\left[\frac{P \cdot h}{L} + \frac{\left(\gamma_{q} \cdot q_{z} + \gamma_{g} \cdot g\right) \cdot L}{2}\right] / \left[\frac{\pi^{2} \cdot E \cdot I_{y,C}}{\left(\beta_{sway} \cdot h\right)^{2}}\right] \le 0.1 \quad (43)$$

$$\eta_{\rm l}^{\rm NS} = \frac{K_{\rm C}}{K_{\rm C} + 0.5 \cdot K_{\rm B}} \tag{44}$$

$$\beta_{\text{non-sway}} = 0.5 + 0.14 \cdot (\eta_{1}^{\text{NS}} + \eta_{2}) + 0.055 \cdot (\eta_{1}^{\text{NS}} + \eta_{2})^{2}$$
(45)

The ULS constraints for the frame columns are defined by Eqs. (46)-(52). Eq. (46) represents the condition for the design bending moment resistance of the column ($M_{sd} < M_{cl,Rd}$), where the substituted expressions *A*, *B* and *C* are given by Eqs. (46 a,b,c). f_y is the yield strength of the structural steel, γ_g is the partial safety factor for the permanent load, γ_q is the partial safety factor for the variable load and γ_{M0} is the resistance partial safety factor. The design shear

resistance $(V_{sd} < V_{pl,Rd})$ and the design axial resistance $(N_{sd} < N_{pl,Rd})$ of the columns are checked by Eqs. (47) to (48).

The reduction factor resulting from the flexural buckling κ , the elastic critical moment for lateral torsional buckling M_{CR} and the reduction factor resulting from lateral torsional buckling κ_{LT} are determined by Eqs. (49) to (51). The substituted expression D in the constraint (49) is defined by Eq. (49 a). C_1 and C_2 are the equivalent uniform moment factors, E is the elastic modulus of steel, G is the shear modulus of steel, k and k_w are effective length factors, π is the Ludolf's number, λ_1 is slenderness, and α_b and $\alpha_{\rm LT}$ are the imperfection factors. The requirement for the interaction between axial compression/buckling and bending moment lateral-torsional buckling is handled by the constraint in Eq. (52).

Eqs. (53) to (56) represent the ULS constraints for beams of the portal frames. The design bending moment resistance of the beam $(M_{\rm sd} < M_{\rm el,Rd})$, the design shear resistance $(V_{\rm sd} < V_{\rm pl,Rd})$ and the design axial resistance $(N_{\rm sd} < N_{\rm pl,Rd})$ are determined by Eqs. (53) to (55). The interaction between axial compression and bending moment is checked by Eq. (56).

Purlins run continuously over the portal frames. The design bending moment resistance about the y-y axis $(M_{y,sd} < M_{el,Rd})$, and the design bending moment resistance about the z-z axis $(M_{z,sd} < M_{el,Rd})$ of the purlins are calculated by Eqs. (57) to (58). The requirement for the interaction between both the mentioned bending moments is handled by the constraint in Eq. (59) The design shear resistance $(V_{sd} < V_{pl,Rd})$ of the purlins are checked by Eq. (60).

The SLS constraints for the portal frames and the purlins are defined by Eqs. (61) to (65). The horizontal deflection of the portal frame Δ and its maximal value are defined by Eqs. (61) to (62). The substituted expressions U and V in constraint (61) are determined by Eqs. (61 a,b).

The vertical deflection of the portal frame δ_F is defined by Eq. (63). This deflection must be smaller than the recommended upper value: the frame span L/250, see Eq. (64). The vertical deflection of the purlins is also checked, see Eq. (65).

$$\frac{\left(\gamma_{q} \cdot q_{z} + \gamma_{g} \cdot g\right) \cdot L^{2} \cdot (3 + 5 \cdot A)}{16 \cdot (B + A \cdot C)} + \frac{P \cdot h \cdot (B + C)}{2 \cdot (B + A \cdot C)} \leq \frac{2 \cdot I_{y,C} \cdot f_{y}}{h_{C} \cdot \gamma_{M0}}$$
(46)

$$A = 1 + f/h; \qquad B = 2 \cdot \left(\frac{I_{y,B}}{I_{y,C}} \cdot \frac{h}{L_{B}} + 1\right) + A; \qquad C = 1 + 2 \cdot A \qquad (46 \text{ a,b,c})$$

$$\left(\frac{\left(\gamma_{q}\cdot q_{z}+\gamma_{g}\cdot g\right)\cdot L^{2}\cdot\left(3+5\cdot A\right)}{16\cdot\left(B+A\cdot C\right)}+\frac{P\cdot h\cdot\left(B+C\right)}{2\cdot\left(B+A\cdot C\right)}\right)\right)/h\leq\frac{\left(1.04\cdot h_{c}\cdot t_{w,C}\right)\cdot f_{y}}{\sqrt{3}\cdot\gamma_{M0}}$$
(47)

$$\frac{P \cdot h}{L} + \frac{\left(\gamma_{q} \cdot q + \gamma_{g} \cdot g\right) \cdot L}{2} \le \frac{A_{c} \cdot f_{y}}{\gamma_{M0}}$$

$$\tag{48}$$

$$\kappa = \frac{1}{0.5 \cdot \left[1 + \alpha_{\rm b} \cdot (D - 0.2) + D^2\right] + \sqrt{\left[0.5 \cdot \left(1 + \alpha_{\rm b} \cdot (D - 0.2) + D^2\right)\right]^2 - D^2}}$$
(49)

$$D = \frac{\beta_{\text{non-sway}} \cdot h}{\sqrt{I_{y,C}/A_{C}} \cdot \lambda_{1}}$$
(49 a)

$$M_{\rm CR} = C_1 \cdot \frac{\pi^2 \cdot E \cdot I_{z,\rm C}}{\left(K \cdot h\right)^2} \cdot \left\{ \sqrt{\frac{I_{\omega,\rm C}}{I_{z,\rm C}} + \frac{\left(K \cdot h\right)^2 \cdot G \cdot I_{u,\rm C}}{\pi^2 \cdot E \cdot I_{z,\rm C}}} + \left(C_2 \cdot \frac{h}{2}\right)^2 - \left(C_2 \cdot \frac{h}{2}\right) \right\}$$
(50)

$$\kappa_{\rm LT} = \frac{1}{0.5 \cdot \left(1 + \alpha_{\rm LT} \cdot \left(\sqrt{\frac{2 \cdot I_{\rm y,C} \cdot f_{\rm y}}{M_{\rm CR} \cdot h_{\rm C}}} - 0.2\right) + \frac{2 \cdot I_{\rm y,C} \cdot f_{\rm y}}{M_{\rm CR} \cdot h_{\rm C}}\right) + \sqrt{\left(0.5 \cdot \left(1 + \alpha_{\rm LT} \cdot \left(\sqrt{\frac{2 \cdot I_{\rm y,C} \cdot f_{\rm y}}{M_{\rm CR} \cdot h_{\rm C}}} - 0.2\right) + \frac{2 \cdot I_{\rm y,C} \cdot f_{\rm y}}{M_{\rm CR} \cdot h_{\rm C}}\right)\right)^2 - \frac{2 \cdot I_{\rm y,C} \cdot f_{\rm y}}{M_{\rm CR} \cdot h_{\rm C}}}$$
(51)

$$\frac{\frac{P \cdot h}{L} + \frac{\left(\gamma_{q} \cdot q_{z} + \gamma_{g} \cdot g\right) \cdot L}{2}}{\kappa \cdot A_{C} \cdot f_{y}/\gamma_{M0}} + \frac{\frac{\left(\gamma_{q} \cdot q_{z} + \gamma_{g} \cdot g\right) \cdot L^{2} \cdot (3 + 5 \cdot A)}{16 \cdot (B + A \cdot C)} + \frac{P \cdot h \cdot (B + C)}{2 \cdot (B + A \cdot C)}}{\kappa_{LT} \cdot 2 \cdot I_{y,C} \cdot f_{y}/(h_{C} \cdot \gamma_{M0})} \le 1.0$$
(52)

$$\frac{\left(\gamma_{q} \cdot q_{z} + \gamma_{g} \cdot g\right) \cdot L^{2} \cdot \left(3 + 5 \cdot A\right)}{16 \cdot \left(B + A \cdot C\right)} + \frac{P \cdot h \cdot \left(B + C\right)}{2 \cdot \left(B + A \cdot C\right)} \leq \frac{2 \cdot I_{y,B} \cdot f_{y}}{h_{B} \cdot \gamma_{M0}}$$
(53)

$$\left(\frac{P \cdot h}{L} + \frac{\left(\gamma_{q} \cdot q_{z} + \gamma_{g} \cdot g\right) \cdot L}{2} \cdot \cos(\alpha) - \left(\frac{\left(\gamma_{q} \cdot q_{z} + \gamma_{g} \cdot g\right) \cdot L^{2} \cdot (3 + 5 \cdot A)}{16 \cdot (B + A \cdot C)} + \frac{P \cdot h \cdot (B + C)}{2 \cdot (B + A \cdot C)}\right) / h \cdot \sin(\alpha) \right) \leq \frac{\left(1.04 \cdot h_{B} \cdot t_{w,B}\right) \cdot f_{y}}{\sqrt{3} \cdot \gamma_{M0}}$$
(54)

$$\frac{\left(\frac{\left(\gamma_{\mathsf{q}}\cdot q_{\mathsf{z}}+\gamma_{\mathsf{g}}\cdot g\right)\cdot L^{2}\cdot \left(3+5\cdot A\right)}{16\cdot \left(B+A\cdot C\right)}+\frac{P\cdot h\cdot \left(B+C\right)}{2\cdot \left(B+A\cdot C\right)}\right)}{2\cdot \left(B+A\cdot C\right)}\right)}{h}+\frac{P\cdot h}{L}+\frac{\left(\gamma_{\mathsf{q}}\cdot q_{\mathsf{z}}+\gamma_{\mathsf{g}}\cdot g\right)\cdot L}{2}\cdot \sin(\alpha) \leq \frac{A_{\mathsf{B}}\cdot f_{\mathsf{y}}}{\gamma_{\mathsf{M0}}}$$

$$\frac{\left(\frac{\left(\gamma_{q}\cdot q_{z}+\gamma_{g}\cdot g\right)\cdot L^{2}\cdot(3+5\cdot A)}{16\cdot(B+A\cdot C)}+\frac{P\cdot h\cdot(B+C)}{2\cdot(B+A\cdot C)}\right)}{A_{B}\cdot f_{y}/\gamma_{M0}}+\frac{\left(\gamma_{q}\cdot q_{z}+\gamma_{g}\cdot g\right)\cdot L^{2}\cdot(3+5\cdot A)}{16\cdot(B+A\cdot C)}+\frac{P\cdot h\cdot(B+C)}{2\cdot(B+A\cdot C)}{2\cdot(B+A\cdot C)}\leq 1.0$$
(56)

(55)

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$$0.1057 \cdot \gamma_{g} \cdot \left(A_{P} \cdot \rho + m_{r}\right) \cdot e_{f}^{2} + 0.1057 \cdot \gamma_{q} \cdot \left(s \cdot \cos(\alpha) \cdot e_{p} + w_{v} \cdot e_{p}\right) \cdot e_{f}^{2} \leq \frac{2 \cdot I_{y,P} \cdot f_{y}}{h_{P} \cdot \gamma_{M0}}$$

$$(57)$$

$$0.1057 \cdot \gamma_{q} \cdot (s+m_{r}) \cdot \sin(\alpha) \cdot e_{p} \cdot e_{f}^{2} \leq \frac{2 \cdot I_{z,p} \cdot f_{y}}{h_{p} \cdot \gamma_{M0}}$$

$$(58)$$

$$\frac{0.1057 \cdot \gamma_{g} \cdot (A_{p} \cdot \rho + m_{r}) \cdot e_{f}^{2} + 0.1057 \cdot \gamma_{q} \cdot (s \cdot \cos(\alpha) \cdot e_{p} + w_{v} \cdot e_{p}) \cdot e_{f}^{2}}{2 \cdot I_{y,p} \cdot f_{y} / h_{p} \cdot \gamma_{M0}} + \frac{0.1057 \cdot \gamma_{q} \cdot (s + m_{r}) \cdot \sin(\alpha) \cdot e_{p} \cdot e_{f}^{2}}{2 \cdot I_{y,p} \cdot f_{y} / h_{p} \cdot \gamma_{M0}} \leq 1.0$$
(59)

$$2 \cdot I_{z,P} \cdot f_{y} / h_{P} \cdot \gamma_{M0}$$

$$0.567 \dots (4.5) = +0.567 \dots (5.57) \dots$$

$$0.567 \cdot \gamma_{\rm g} \cdot (A_{\rm p} \cdot \rho) \cdot e_{\rm f} + 0.567 \cdot \gamma_{\rm q} \cdot \left(s \cdot \cos(\alpha) \cdot e_{\rm p} + w_{\rm v} \cdot e_{\rm p}\right) \cdot e_{\rm f} \le \frac{1.04 \cdot \eta_{\rm p} \cdot \eta_{\rm w, P} \cdot \gamma_{\rm y}}{\sqrt{3} \cdot \gamma_{\rm M0}} \tag{60}$$

$$\Delta = \frac{1}{E \cdot I_{y,C}} \cdot \left(\frac{1}{3} \cdot L \cdot \left((V - U) \cdot h\right)\right) + \frac{1}{E \cdot I_{y,B}} \cdot \left(\frac{1}{6} \cdot L \cdot \left(V - 3 \cdot U\right) \cdot h - \frac{1}{3} \cdot L \cdot \left(\frac{(q + g) \cdot L^2}{8} \cdot h\right)\right)$$
(61)

$$U = \frac{(q+g) \cdot L^2 \cdot (3+5 \cdot A)}{16 \cdot (B+A \cdot C)} \qquad \qquad V = \frac{P \cdot h \cdot (B+C)}{2 \cdot (B+A \cdot C)}$$
(61 a,b)

 $\varDelta \leq h \, / \, 150$

 $\delta_{\rm F} = L/250$

$$=\frac{(q+g)}{24\cdot E\cdot I_{\rm VB}}\cdot \left(\frac{L}{2} + \frac{2\cdot V}{(q+g)\cdot L}\right)\cdot \left(\frac{\left(\frac{L}{2} + \frac{2\cdot V}{(q+g)\cdot L}\right)^3 - \left(2\cdot L + \frac{8\cdot V}{(q+g)\cdot L}\right)\cdot \left(\frac{L}{2} + \frac{2\cdot V}{(q+g)\cdot L}\right)^2 + \frac{12\cdot (V-U)}{(1+g)\cdot L}\right)$$
(63)

$$\delta_{\rm F} = \frac{(q+g)}{24 \cdot E \cdot I_{\rm y,B}} \cdot \left[\frac{L}{2} + \frac{2 \cdot V}{(q+g) \cdot L} \right] \cdot \left[\left(\frac{12 \cdot (V-U)}{(q+g)} \cdot \left(\frac{L}{2} + \frac{2 \cdot V}{(q+g) \cdot L} \right) + L^3 + \frac{(12 \cdot U - 4 \cdot V) \cdot L}{(q+g)} \right] \right]$$
(65)

$$\left[\left(s \cdot \cos(\alpha) \cdot e_{p} + \left(w_{v} \cdot e_{p} + A_{p} \cdot \rho \cdot e_{p} + m_{r} \cdot e_{p}\right) \cdot \cos(\alpha)\right) \cdot \left(269/42000\right) \cdot e_{f}^{4}/\left(E \cdot I_{y,P}\right)\right] \cdot \cos(\alpha) + \left[\left(s \cdot e_{p} + A_{p} \cdot \rho \cdot e_{p} + m_{r} \cdot e_{p}\right) \cdot \sin(\alpha) \cdot \left(269/42000\right) \cdot e_{f}^{4}/\left(E \cdot I_{z,P}\right)\right] \cdot \sin(\alpha) \le \frac{e_{f}}{250}$$
(65)

3.3. Integer and mixed integer logical constraints

The logical constraint in Eq. (66) defines the number of portal frames, where y_n denotes the binary variable which is subjected to each portal frame. Eq. (67) defines only one possible vector of binary variables for each frame topology. Eq. (68) calculates the even number of purlins, where the binary variables y_m are subjected to the purlins. Eq. (69) defines only one possible vector of the binary variables for each purlin topology.

$$NOFRAME = \sum_{n} y_{n}$$
(66)

$$y_{\rm n} \le y_{\rm n-l} \tag{67}$$

$$NOPURLIN = 2 \cdot \sum_{m} y_{m}$$
(68)

$$y_{\rm m} \le y_{\rm m-1} \tag{69}$$

Eqs. (70) to (79) calculate the standard cross-sections for the columns with their discrete dimensions and characteristics. The latter are determined as scalar products between their vectors of *i*, $i \in I$, the discrete standard constants $(q_i^{A_c}, q_i^{A_c}, \dots)$ and their vector of subjected binary variables y_i , see Eqs. (70) to (78). Only one discrete value is then selected for each standard section since the sum of the binary variables must be equal to 1, see Eq. (79).

$$A_{\rm C} = \sum_{\rm i} q_{\rm i}^{\rm A_{\rm C}} \cdot y_{\rm i} \qquad \qquad {\rm i} \in I \qquad (70)$$

(62)

(64)

$$h_{\rm C} = \sum_{i} q_{i}^{\rm h_{\rm C}} \cdot y_{i} \qquad i \in I \qquad (71)$$
$$b_{\rm C} = \sum_{i} q_{i}^{\rm h_{\rm C}} \cdot y_{i} \qquad i \in I \qquad (72)$$

$$t_{w,c} = \sum_{i}^{i} q_{i}^{t_{w,c}} \cdot y_{i} \qquad i \in I \qquad (73)$$

$$I_{f,C} = \sum_{i} q_{i}^{I,C} \cdot y_{i} \qquad i \in I \qquad (74)$$
$$I_{y,C} = \sum_{i} q_{i}^{I,C} \cdot y_{i} \qquad i \in I \qquad (75)$$

$$I_{z,c} = \sum_{i}^{i} q_{i}^{L_{z,c}} \cdot y_{i} \qquad i \in I$$
(76)

$$I_{t,C} = \sum_{i} q_{i}^{I_{t,C}} \cdot y_{i} \qquad i \in I \qquad (77)$$
$$I_{t,C} = \sum_{i} q_{i}^{I_{u,C}} \cdot y_{i} \qquad i \in I \qquad (78)$$

$$\sum_{i} y_{i} = 1 \tag{79}$$

Similarly, Eqs. (80) to (86) determine the discrete values of the cross-sectional characteristics for the frame beams and Eqs. (87) to (96) for purlins.

$$A_{\rm B} = \sum_{j} q_{j}^{\rm A_{\rm B}} \cdot y_{j} \qquad \qquad j \in J$$
(80)

$$h_{\rm B} = \sum_{j} q_{j}^{\rm h_{\rm B}} \cdot y_{j} \qquad \qquad j \in J \qquad (81)$$

$$b_{\rm B} = \sum_{j} q_{j}^{b_{\rm B}} \cdot y_{j} \qquad \qquad j \in J \qquad (82)$$
$$t_{\rm w,B} = \sum_{j} q_{i}^{t_{\rm w,B}} \cdot y_{i} \qquad \qquad j \in J \qquad (83)$$

$$t_{\mathrm{f,B}} = \sum_{j} q_{j}^{t_{\mathrm{f,B}}} \cdot y_{j} \qquad j \in J \qquad (84)$$

$$I_{y,B} = \sum_{j}^{J} q_{j}^{I_{y,B}} \cdot y_{j} \qquad j \in J$$
(85)

$$\sum_{j} y_{j} = 1 \tag{86}$$

$$A_{p} = \sum_{k} q_{k}^{A_{p}} \cdot y_{k} \qquad k \in K \qquad (87)$$
$$h_{p} = \sum q_{k}^{A_{p}} \cdot y_{k} \qquad k \in K \qquad (88)$$

$$h_{-} = \sum_{k} a^{b_{p}} \cdot v, \qquad k \in K$$

$$(89)$$

$$f_{w,P} = \sum_{k} q_{k}^{t_{w,P}} \cdot y_{k} \qquad k \in K$$
(90)

 $\mathbf{k} \in K$

$$I_{z,P} = \sum_{k} q_{k}^{I_{z,P}} \cdot y_{k} \qquad k \in K$$
(93)

$$I_{t,P} = \sum_{k} q_{k}^{I_{t,P}} \cdot y_{k} \qquad k \in K \qquad (94)$$

$$I_{\omega,P} = \sum_{k}^{k} q_{k}^{I_{\omega,P}} \cdot y_{k} \qquad \qquad k \in K$$
(95)

$$\sum_{k} y_{k} = 1$$
(96)

3.4. Sets, input data (constants) and variables

The following sets, input data (constants: scalars and parameters) as well as continuous and binary variables are involved in the optimization model FRAMEOPT:

Sets:

- *i* set for the standard dimension alternatives for columns, $i \in I$
- *j* set for the standard dimension alternatives for beams, $j \in J$
- k set for the standard dimension alternatives for purlins, $k \in K$
- *m* set for the number of purlins, $m \in M$
- *n* set for the number of portal frames (columns and beams), $n \in N$

Scalars (constants, input data):

- f denotes the overheight of the frame beam [cm]
- f_y yield the strength of the structural steel [kN/cm²]
- *h* height of the column [cm]
- *k* effective length factor [-]
- $k_{\rm w}$ effective length factor [-]
- $m_{\rm r}$ mass of the roof plates [kg/cm²]
- s snow (variable imposed load) $[kN/cm^2]$
- w_v vertical wind (variable imposed load) [kN/cm²]
- w_h horizontal wind (variable imposed load) [kN/cm²]

 C_1, C_2 equivalent uniform moment factors [-]

E elastic modulus of steel [kN/cm²]

- G shear modulus of steel [kN/cm²]
- *L* frame span [m]
 - $L_{\rm L}$ length of the industrial building [m]
- *MINNO*^{frame} minimum number of defined portal frames [-]
- (92) *MAXNO*^{frame} maximum number of defined portal frames [-]
 - *MINNO*^{purlin} minimum number of purlins [-]

 $I_{\mathbf{y},\mathbf{P}} = \sum_{\mathbf{y}} q_{\mathbf{k}}^{\mathbf{I}_{\mathbf{y},\mathbf{P}}} \cdot y_{\mathbf{k}}$
MAXNO ^{purlin}	maximum	number	of purlins	[-]	
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- $\alpha_{\rm b}$ imperfection factor [-]
- $\alpha_{\rm LT}$ imperfection factor [-]
- γ_{q} partial safety factor for the variable load [-]
- γ_g partial safety factor for the permanent load [-]
- γ_{M0} resistance partial safety factor [-]
- η_2 distribution factor [1]
- λ_1 slenderness [-]
- π Ludolf's number [-]
- ρ density of steel [kg/m³]

Parameters (constants, input data):

- $q_{i}^{A_{c}}$ vector of *i*, $i \in I$, discrete standard constants for cross-section area of the column
- $q_{j}^{A_{B}}$ vector of $j, j \in J$, discrete standard constants for cross-section area of the beam
- $\boldsymbol{q}_{k}^{A_{p}}$ vector of $k, k \in K$, discrete standard constants for cross-section area of the purlin
- $q_{i}^{b_{c}}$ vector of *i*, $i \in I$, discrete standard constants for overall breadth of the column
- $q_{j}^{b_{B}}$ vector of $j, j \in J$, discrete standard constants for overall breadth of the beam
- $q_{k}^{b_{p}}$ vector of $k, k \in K$, discrete standard constants for overall breadth of the purlin
- $\boldsymbol{q}_{i}^{t_{tc}}$ vector of *i*, $i \in I$, discrete standard constants for flange thickness of the column
- $q_{j}^{t_{f,B}}$ vector of $j, j \in J$, discrete standard constants for flange thickness of the beam
- $q_{k}^{t_{t,p}}$ vector of $k, k \in K$, discrete standard constants for flange thickness of the purlin
- $q_i^{t_{w,c}}$ vector of *i*, $i \in I$, discrete standard constants for web thickness of the column
- $q_{j}^{t_{w,B}}$ vector of $j, j \in J$, discrete standard constants for web thickness of the beam
- $q_{k}^{t_{w,P}}$ vector of $k, k \in K$, discrete standard constants for web thickness of the purlin
- $q_i^{I_{tc}}$ vector of *i*, $i \in I$, discrete standard constants for torsional constant of the column
- $q_{k}^{l_{i,p}}$ vector of $k, k \in K$, discrete standard constants for torsional constant of the purlin

- $\boldsymbol{q}_{i}^{\mathrm{I}_{\mathrm{y},\mathrm{c}}}$ vector of $i, i \in I$, discrete standard constants for second moment of area about
- the y y axis of the column $q_{i,j,B}^{I_{y,B}}$ vector of $j, j \in J$, discrete standard
 - constants for second moment of area about the y y axis of the beam
- $q_k^{1_{y,p}}$ vector of $k, k \in K$, discrete standard constants for second moment of area about the y y axis of the purlin
- $q_i^{1,c}$ vector of *i*, $i \in I$, discrete standard constants for moment of area about the z -z axis of the column
- $q_k^{1_{z,p}}$ vector of $k, k \in K$, discrete standard constants for moment of area about the z -z axis of the purlin
- $q_i^{l_{u,C}}$ vector of $i, i \in I$, discrete standard constants for warping constant of the column
- $q_k^{I_{u,p}}$ vector of $k, k \in K$, discrete standard constants for warping constant of the purlin

Continuous variables:

- $b_{\rm B}$ overall breadth of the beam [cm]
- $b_{\rm C}$ overall breadth of the column [cm]
- $b_{\rm P}$ overall breadth of the purlin [cm]
- *e*_f intermediate distance between the portal frames [cm]
- *e*_p intermediate distance between the purlins [cm]
- *g* self-weight of the portal frame [kN/cm]
- $h_{\rm B}$ cross-sectional height of the beam [cm]
- $h_{\rm C}$ cross-sectional height of the column [cm]
- $h_{\rm P}$ cross-sectional height of the purlin [cm]
- *q*_z uniformly distributed horizontal surface variable load [kN/cm]
- *q*_y uniformly distributed vertical surface variable load [kN/cm]
- $t_{\rm f,B}$ flange thickness of the beam [cm]
- $t_{\rm f,C}$ flange thickness of the column [cm]
- $t_{\rm f,P}$ flange thickness of the purlin [cm]
- $t_{\rm w,B}$ web thickness of the beam [cm]
- $t_{\rm w,C}$ web thickness of the column [cm]
- $t_{\rm w,P}$ web thickness of the purlin [cm]
- $A_{\rm B}$ cross section of the beam [cm²]
- $A_{\rm C}$ cross section of the column [cm²]
- $A_{\rm P}$ cross section of the purlin [cm²]

- $I_{t,C}$ torsional constant of the column [cm⁴]
- $I_{t,P}$ torsional constant of the purlin [cm³]
- $I_{y,B}$ second moment of area about the y y axis of the beam [cm⁴]
- $I_{y,C}$ second moment of area about the y y axis of the column [cm⁴]
- $I_{y,P}$ second moment of area about the y-y axis of the purlin [cm⁴]
- $I_{z,C}$ second moment of area about the z z axis of the column [cm⁴]
- $I_{z,P}$ second moment of area about the z-z axis of the purlin [cm⁴]
- $I_{\omega,C}$ warping constant of the column [cm⁶]
- $I_{\omega,P}$ warping constant of the purlin [cm⁶]
- $K_{\rm C}$ stiffness coefficient of the column [m³]
- $K_{\rm B}$ stiffness coefficient of the purlin [m³]
- $L_{\rm B}$ length of the beam [cm]
- $M_{\rm CR}$ elastic critical moment for lateral torsional buckling [kNcm]
- M_{eLRd} design elastic moment resistance [kNcm]

NOFRAME number of the portal frames [-]

NOPURLIN number of the purlins [-]

 $N_{\rm pl,Rd}$ design plastic axial resistance[kN]

- $N_{\rm sd}$ design axial force [kN]
- *M*_{sd} design bending moment [kNcm]
- P concentrated horizontal variable load multiplied by the partial safety factor [kN]
- $P_{\rm w}$ concentrated horizontal variable load [kN]
- $V_{\rm pl,Rd}$ design plastic shear resistance [kN]
- $V_{\rm sd}$ design shear force [kN]

 α angle of the inclination of the beam [rad]

- $\beta_{\text{non-sway}}$ plane buckling length of the column for a non-sway frame [-]
- β_{sway} plane buckling length of the column for a sway frame [-]
- $\delta_{\rm F}$ vertical deflection of the portal frame [cm]
- Δ horizontal deflection of the portal frame [m]
- η_1^{NS} distribution factors of the column for the non-sway frame [-]
- $\eta_1^{\rm s}$ distribution factors of the column for the sway frame [-]
- κ reduction factor due to the flexural buckling [-]
- κ_{LT} reduction factor for lateral-torsional buckling [-]

Binary variables:

- y_i binary variable assigned to the *i*-th, $i \in I$, standard dimension alternative of the columns
- y_j binary variable assigned to the *j*-th, $j \in J$, standard dimension alternative of the beams
- y_k binary variable assigned to the *k*-th, $k \in K$, standard dimension alternative of the purlins
- $y_{\rm m}$ binary variable assigned to the *m*-th, $m \in M$, topology alternative of the purlins
- y_n binary variable assigned to the *n*-th, $n \in N$, topology alternative of the frames

Substituted expressions:

- *A,B,C,D* functions which are substituded in Eqs. 46 a,b,c and 49 a
- *U,V* functions which are substituded in Eqs. 61 a,b

4 OPTIMIZATION

Outer-Approximation The Modified /Equality-Relaxation algorithm (OA/ER) (Kravanja and Grossmann [6]) was used to perform the optimization. The OA/ER algorithm consists of solving an alternative sequence of Non-linear Programming optimization subproblems (NLP) and Mixed-Integer Linear Programming master problems (MILP), Fig. 3. The former corresponds to the optimization of parameters for a building structure with a fixed topology and standard dimensions and yields an upper bound to the objective to be minimized. The latter involve a global linear approximation to the superstructure of alternatives in which a new topology and standard sizes are identified. When the problem is convex the search is terminated when the predicted lower bound exceeds the upper bound, otherwise it is terminated when the NLP solution can be improved no more. The OA/ER algorithm guarantees the global optimality of solutions for convex and quasi-convex optimization problems.

The OA/ER algorithm as well as all other algorithms do not generally guarantee that the solution found is the global optimum. This is due to the presence of non-convex functions in the models that may cut off the global optimum. In order to reduce undesirable effects of



nonconvexities, the following nonstructured and structured convexifications are applied for the

Fig. 3. Steps of the OA/ER algorithm

decomposition and the deactivation of the objective function linearization, the use of the penalty function, the use of the upper bound on the objective function to be minimized as well as the global convexity test and the validation of the outer approximations. By the use of the mentioned modifications, the likelihood of obtaining better results by the OA/ER algorithm, is significantly increased. A more extended information about these modifications may be found elsewhere, see Kravanja and Grossmann [6], and Kravanja et al. [7].

The optimum solution of a complex nonconvex and non-linear MINLP problem with a high number of discrete decisions is in general very difficult to obtain. The optimization is thus proposed to be performed sequentially in two different phases to accelerate the convergence of the OA/ER algorithm. The optimization is proposed to start with the discrete topology optimization of the building, while the standard MILP master problem of the OA/ER algorithm: the deactivation of linearizations, the dimensions are temporarily relaxed into continuous parameters. Topology and continuous parameter optimization is soluble (a smaller combinatorial problem) and accumulates a good global linear approximation of the superstructure (a good starting point for the next phase overall optimization). When the optimum topology is found, the standard sizes of the cross-sections are re-established and the simultaneous discrete optimization of the topology and standard dimensions of the beams, columns and purlins is then continued until the optimum solution is found.

The two-phase strategy requires that the binary variables should be defined in one uniform set. In the first phase, only the binary variables which are subjected to topology alternatives become active. Binary variables of standard dimension alternatives are temporarily excluded (set on value zero) until the beginning of the second phase, in which they participate in the simultaneous overall optimization. The same holds for standard dimension logical constraints. In the first phase they are excluded, while the second phase includes them into the optimization.

The data and variables are initializated only once in the beginning of the optimization. An advantage of this strategy is also in the fact that binary variables for topology and standard dimensions need not be initialized: after the first NLP, the first phase always starts in the subspace of the topological binary variables only, while the second phase starts with the MILP master subproblem which then predicts a full set of binary variables for the successive NLP. Under the convexity condition, the two-phase strategy guarantees a global optimality of the solution.

The optimization model may contain up to thousand binary 0-1 variables of the alternatives. Most of them are subjected to standard dimensions. Since this number of 0-1 variables is too high for a normal solution of the MINLP, a reduction procedure was developed, which automatically reduces the binary variables of alternatives into a reasonable number. The optimization at the second phase includes only those 0-1 variables which determine the topology and standard dimension alternatives close to the values, obtained at the first MINLP optimization phase.

5 NUMERICAL EXAMPLE

The paper presents an example of the topology and standard dimension optimization of a single-storey industrial building. The building is 25 meters wide (*L*), 75 meters long (L_L) and 6 meters height (*H*), see Fig. 4. The structure consists of equal non-sway steel portal frames which are mutually connected with purlins. The overheight of the frame beam (*f*) is 0.50 m.

The portal frame is subjected to selfweight of the structure and the roof g, and to the variable loads of snow and wind. The mass of the roof is $m_r = 0.20 \text{ kg/m}^2$. The variable imposed loads: $s = 2.00 \text{ kN/m}^2$ (snow), $w_v = 0.125 \text{ kN/m}^2$ (vertical wind) and $w_h = 0.50 \text{ kN/m}^2$ (total horizontal wind) are defined in the model input data. Both, the horizontal concentrated load at the top of the columns and the vertical uniformly distributed line load on the beams and purlins are calculated automatically through the optimization considering the calculated intermediate distance between the portal frames and purlins.

The material used is steel S 355. The yield strength of the steel (f_y) is 35.5 kN/cm², the density of steel (ρ) is 7.850·10⁻³ kg/cm³, the elastic modulus of steel (E) is 210 GPa and the shear modulus (G) is 80.76 GPa. The partial safety factor for the permanent load (γ_{q}) and for the combined (snow plus wind) variable load (γ_q) are both 1.35. The resistance partial safety factor (γ_{M0}) is 1.1. The imperfection factor $(\alpha_{\rm b})$ is 0.34, the imperfection factor (α_{LT}) is 0.21, the distribution factor (η_2) is 1, slenderness for the steel S 355 (λ_1) is 76.4, the effective length factors (k and k_w) are 1.0, the equivalent uniform moment factors for beams (C_1) and (C_2) are 1.879 and 0, respectively. While the defined minimum and maximum numbers of portal frames (MINNO^{frame} and MAXNO^{frame}) are 1 and 30, the minimal and maximal numbers of purlins (MINNO^{purlin} and MAXNO^{purlin}) are 1 and 20.



Fig. 4. Global geometry of the single-storey industrial building

The lower and upper bounds as well as the activity levels (starting points) of the independent continuous variables are shown in Table 1. The bounds and starting points of other dependant continuous variables are defined by using equations from the optimization model regarding the independent variables.

Table 1. Bounds and activity levels of theindependent variables

Variable	Lower	Activity	Upper
x	bound	level	bound
	x^{LO}	x^{L}	x^{UP}
$h_{ m C}$	50 cm	80 cm	99 cm
$h_{ m B}$	30 cm	60 cm	70 cm
$h_{ m P}$	10 cm	20 cm	30 cm
NOFRAMES	1	20	30
NOPURLINS	1	20	20

An industrial building superstructure was generated in which all possible structures were embedded by 30 portal frame alternatives, 10 various purlin alternatives and a variation of different standard sizes. The superstructure comprised 24 different standard hot rolled European wide flange I sections, i.e. HEA sections (from HEA 100 to HEA 1000) for each column, beam and purlin seperately. Vectors \boldsymbol{q} of 24 discrete values for different standard sections were defined. For example, the vectors for the section's heights $\boldsymbol{q}_i^{h_c}$, $\boldsymbol{q}_j^{h_g}$, $\boldsymbol{q}_k^{h_p}$ and the crosssection areas $\boldsymbol{q}_i^{A_c}$, $\boldsymbol{q}_j^{A_g}$, $\boldsymbol{q}_k^{A_p}$ are for the columns, beams and purlins defined as follows:

 $\boldsymbol{q}_{i}^{h_{C}} = \boldsymbol{q}_{i}^{h_{B}} = \boldsymbol{q}_{k}^{h_{P}} = \{9.6, 11.4, 13.3, 15.2, 17.1, \}$

19.0, 23.0, 25.0, 27.0, 29.0, 31.0, 33.0, 35.0, 39.0, 44.0, 49.0, 54.0, 59.0, 64.0, 69.0, 79.0, 89.0, 99.0}

$$\boldsymbol{q}_{i}^{A_{C}} = \boldsymbol{q}_{j}^{A_{B}} = \boldsymbol{q}_{k}^{A_{P}} = \{21.2, 25.3, 31.4, 38.8, 45.3,$$

53.8, 64.3, 76.8, 86.8, 97.3, 113.0, 124.0, 133.0, 143.0, 159.0, 178.0, 198.0, 212.0, 226.0, 242.0, 260.0, 286.0, 321.0, 347.0}

Regarding construction alternatives, the superstructure consists of a n possible number of

portal frames, $n \in N$, $N = \{1,2,3,...,30\}$, and 10 various even (2m) numbers of purlins, $m \in M$, $M = \{1,2,3,...,10\}$, which give $30 \cdot 10 = 300$ different topology alternatives. Since *i*, *j* and *k* different standard sections are also defined for columns, beams and purlins separately, $i \in I$, $j \in J$, $k \in K$, $I = J = K = \{1,2,3,4,5,6,7,...,24\}$, there exist $n \cdot m \cdot i \cdot j \cdot k = 30 \cdot 10 \cdot 24 \cdot 24 \cdot 24 = 4147200$ different discrete construction alternatives alltogether.

The optimization was performed by the proposed MINLP optimization approach. The task of the optimization was to find the minimal structure mass, the optimum topology (the optimum number of portal frames and purlins) and the optimum standard sizes.

The optimization was carried out by a user-friendly version of the MINLP computer package MIPSYN, the successor of PROSYN [6] and TOP [7], [8] and [10]. The Modified OA/ER algorithm and the two-phase optimization were applied, where GAMS/CONOPT2 (Generalized reduced-gradient method) [11] was used to solve the NLP subproblems and GAMS/Cplex 7.0 (Branch and Bound) [12] was used to solve the MILP master problems.

The two-phase MINLP optimization was applied. After the first performed continuous NLP (the initialization), the first phase started with the discrete topology optimization at the relaxed standard dimensions, see also the convergence of the Modified OA/ER algorithm in Table 2. At this level, only the binary variables y_n and y_m for topology optimization, parameter structural nonlinear and linear constraints, Eqs.(4) to (65), and the logical constraints for topology optimization, Eqs. (66) to (69), were included. When the optimum topology was reached (110.161 tons at the 2nd MINLP iteration, all the following solutions were poorer), the optimization proceeded with a simultaneous discrete topology and standard dimension optimization at the second level. At this phase, the binary variables y_i , y_i and y_k of standard sizes for columns, beams and purlins, as well as the logical constraints for standard dimensions, Eqs. (70) to (96), were added into the optimization. The final optimum solution of 122.144 tons was obtained at the 6th main MINLP iteration (all the following solutions were not as good).

MINLP	MINLP	Result	Торс	ology	Cross-sections [cm ²]				
Iteration	Subphaze	Mass [tons]	Frames	Purlins	Column	Beam	Purlin		
Phase 1: topology optimization									
1	Initialization								
1.	1.NLP	107.254	11.955	12.008	283.501	212.036	39.892		
2	1.MILP	107.763			321.721	176.232	30.527		
<i>Z</i> .	2.NLP	110.161	13	14	276.759	206.027	32.137		
2	2.MILP	114.351			326.447	174.106	30.509		
5.	3.NLP	111.339	14	14	270.667	200.658	28.442		
	Ph	ase 2: topology	and standa	ard dimens	sion optimization	on			
4	3.MILP	125.260			321.00	198.00	38.80		
4.	4.NLP*	125.231	14	14	HEA 900	HEA 550	HEA 160		
5	4.MILP	115.708			321.00	212.00	38.80		
5.	5.NLP*	115.209	12	14	HEA 900	HEA 550	HEA 160		
6	5.MILP	122.144			321.00	212.00	38.80		
0.	6.NLP	122.144	13	14	HEA 900	HEA 550	HEA 160		
7.	6.MILP	126.713			321.00	212.00	38.80		
	7.NLP	126.713	13	16	HEA 900	HEA 550	HEA 160		

 Table 2. Convergence of the Modified OA/ER algorithm

* Locally infeasible

The optimum result represents the mentioned minimal structure mass of 122.144 tons, the obtained optimum topology of 13 portal frames and 14 purlins, see Fig. 5, and the calculated optimum standard sizes of the columns (HEA 900), beams (HEA 550) and purlins (HEA 160), see Fig. 6.

At the second phase, where all the calculated dimensions were standard ones, a feasible optimum result was very difficult to be obtained. The optimization model contained a high number of 4147200 different discrete construction alternatives.



Fig. 5: Optimum design of the single-storey industrial building



Fig. 6: Optimum design of the portal steel frame

The prescreening procedure of alternatives was thus applied, which automatically reduced the binary variables of alternatives into a reasonable number. The optimization at the second phase included only those 0-1 variables which determined the topology and standard dimension alternatives close to the (continuous) values, obtained at the first phase. For topology, column, beam and purlin only 3 binary variables were used (1 variables under and 2 over the continuous value). In this way, only 15 binary variables were used in the second phase instead of all 112 binary variables. The number of 4147200 discrete construction alternatives was significantly reduced to $n \cdot m \cdot i \cdot j \cdot k = 3 \cdot 3 \cdot 3 \cdot 3 \cdot 3 = 243$ alternatives, which considerably improved the efficiency of the search.

6 CONCLUSIONS

The paper presents the simultaneous topology and standard sizes optimization of a single-storey industrial steel building. The optimization is proposed to be performed by the Mixed-Integer Non-linear Programming (MINLP) approach. The Modified OA/ER algorithm and the two-phase MINLP optimization strategy were applied. The proposed two-phase optimization starts with the topology optimization of the frames and purlins, while the standard dimensions are temporarily relaxed into

continuous parameters. When the optimum topology is found, the standard dimensions of the cross-sections are re-established and the simultaneous topology and discrete standard dimension optimization of beams, columns and purlins is then continued until the optimum solution is found. Without performing the twophase MINLP strategy and the prescreening procedure of alternatives no feasible optimum result was obtained. The proposed MINLP was found to be a successful optimization technique for solving this type of structures.

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Geometric Accuracy by 2-D Printing Model

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In order to improve the accuracy of three-dimensional printing (3DP), we carried out a set of experiments on two-dimensional (2-D) printed model. The combinations of three main processing factors were considered for every experiment: layer thickness, building orientation and infiltrant type. The results of experiments provided the appropriate anisotropic scale factor and proved the adequacy of anisotropic scaling. The analysis of results indicates that smaller layer thickness, infiltration with epoxy resin and sample orientation towards building direction of X axis provide better sample accuracy. © 2008 Journal of Mechanical Engineering. All rights reserved.

Keywords: rapid prototyping, three-dimensional printing, geometric accuraccy

0 INTRODUCTION

The models produced by threedimensional printing (3DP) are not so accurate when compared to other Rapid Prototyping (RP) technologies. Several authors emphasized the problem of 3DP accuracy in their researches. Dimitrov et al. pointed out in their review of 3DP that the accuracy is a significant weakness when compared to other RP technologies [1]. Pilipović et al. noticed that 3DP test samples exceed tolerance limits for tensile testing and are less accurate than samples made with a hybrid Polyjet technique [2].

The researchers are constantly trying to improve 3DP accuracy which is evident in the number of published papers [3] to [8]. The effort for improvement is also reflected on a 3DP market - the manufacturers not only develop and deliver the new models of printers, but also new improved materials, device firmware, software and improved alternative spare parts for the existing models.

The primary objective of our paper was to determine the appropriate scale factors for the considered 3DP system in the selected combination of materials and processing parameters. The possible use of the appropriate scale factors for improving printer accuracy is also pointed out in [6] and it gave us the additional motivation to perform the experiments. The secondary objective was to determine the influence of 3DP processing factors on the accuracy and to determine the most accurate combination. Therefore, we carried out the experiments to improve the accuracy of the considered 3DP system.

1 3DP EQUIPMENT AND MATERIALS

The 3D printer, used for these experiments, was the model Z310, a product of Z Corporation. It is a low-cost monochrome 3D printer suitable for RP education or for small and medium sized companies. The printer firmware version was 10.158 and test samples were prepared in a printer software ZPrint version 7.5.23 [9] and [10].

The considered 3D printer combines a layered approach from RP technologies and a conventional ink-jet printing. It prints a binder fluid through the conventional ink-jet print head into a powder, one layer onto another, from the lowest model's cross-section to the highest (Fig. 1). After printing, the printed models are dried in a building box, then removed from the powder bed, de-powdered by compressed air, dried in the oven and infiltrated for maximum strength.

There are several base materials, i.e. powder types, available for the above mentioned 3D printer. For our experiment we used a plasterbased powder zp130 with an appropriate binder zb56. The powder zp130 is recommended for the accuracy and for delicate models. It is a mixture of plaster, vinyl polymer and sulfate salt [11].

All test samples were dried two times as recommended in [10] and [12]: first in the printer's building box for one hour and then, after de-powdering, in the oven for at least two hours at 55°C.

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Fig. 1. 3D printer

After drying, the samples were infiltrated with the appropriate infiltrant, regarding an appropriate combination of the experiment. The applied infiltrants were: cyanoacrylate based Loctite 406; epoxy resin based Loctite Hysol 9483 and normal wax Cera Alba.

2 DESIGN OF THE EXPERIMENT

In our experiments we used models for tensile tests defined in standard ISO 527:1993. Tensile test model is selected because it has sufficient different dimensions and the models can be used for additional tests of mechanical properties if needed. The measured dimensions on test samples were: total length (L), width at the end (W), neck width (W1) and height (H). The position of dimensions on the test sample and nominal values of dimensions are presented in Fig. 2. Since the tensile test model has the small hight regarding other dimensions, we considered it as a 2-D model.

Two sets of test samples were planned. For the first set we considered non-scaled samples, so called the base set. The samples of the second set will be scaled with appropriate scale factors obtained from the measuring of the base set.

In both sets, the test samples were printed in a combination of considered processing factors. We considered combinations of the following processing factors: layer thickness, building orientation and infiltrant type.



Fig. 2. Measured dimensions of the test sample

Layer thickness can be selected from two possible values -0,1 mm and 0,0875 mm. The first thickness is default for the printer and it is therefore marked with number 1 in the beginning of the label for a particular combination of factors.

The model can be oriented in any possible direction inside the printer building box. We considered two directions: sample oriented with largest dimension L towards the building axis X and towards axis Y. Thereby, the samples were aligned at the bottom plane of the building box. In the experiment label of a particular combination of factors, sample orientation is marked on the second place of the label with letter X or Y respectively. The orientation towards the building axis Z was omitted since it considerably prolonged the printing time.

The type of the infiltrant is marked with letter at the end of the label for a particular combination of factors. The letters used for a particular infiltrant type are: W for wax, E for epoxy resin and C for cyanoacrylate.

For every particular experiment i.e. particular combination of factors in each set, 6 test samples were printed and measured (Fig. 3).



Fig. 3. 3D printing of test samples

We measured dimensions of test samples with digital caliper Lux Profi model 572587, with the measurement range 0-150 mm and accuracy 0.01 mm.

3 RESULTS

The measurement results of the base set samples are presented in Table 1. Common statistical values are calculated and presented in table rows, grouped by the measured dimension and combinations. Calculated values are: arithmetic mean (\overline{x}); standard deviation (*S*); relative standard deviation (*RSD*); average error (Δx); relative error ($\Delta x/X$) and relative deviation (\overline{x}/X).

The relative standard deviation adequately expresses the precision of a particular experiment combination regarding the measured dimensions. It is the absolute value of the coefficient of variation, usually expressed as a percentage and calculated by:

$$RSD = \frac{S}{x} \cdot 100 \tag{1}$$

If average standard deviations for every measured dimension in the base set of samples are presented together in a single chart, it can be noticed that the standard deviation is higher for larger dimensions (Fig. 4).



Fig. 4. Average stand. deviations in the base set

However, if relative standard deviations for each measured dimension are considered, it can be noticed that the relative standard deviation is higher for smaller dimensions (Fig. 5). Furthermore, average relative errors are higher for larger dimensions (Fig. 6).



Fig. 5. Average relative standard deviations in the base set



Fig. 6. Average relative errors in the base set

Relative errors and relative standard deviations show significant differences between measured dimensions. Such significant differences justify the anisotropic scaling of every particular dimension instead of the common overall scale factor.

Layer thickness			0.1	mm					0.087	75 mm		
Infiltrant	w	ax	Epoxy	y resin	Cyanoa	crylate	W	ax	Epoxy	y resin	Cyano	acrylate
Orientation	X	Y	X	Y	X	Y	X	Y	X	Y	X	Y
Experiment label	1XW	1YW	1XE	1YE	1XC	1YC	2XW	2YW	2XE	2YE	2XC	2YC
Dimension L (Length), m	m										
$\frac{-}{x}$	151.03	150.81	151.09	150.90	151.09	150.89	151.37	151.25	151.36	151.22	151.51	151.27
S	0.10	0.17	0.07	0.13	0.09	0.12	0.05	0.18	0.13	0.06	0.09	0.12
RSD, %	0.07	0.11	0.04	0.09	0.06	0.08	0.03	0.12	0.08	0.04	0.06	0.08
Δx	1.03	0.81	1.09	0.90	1.09	0.89	1.37	1.25	1.36	1.22	1.51	1.27
$\Delta x/X$, %	0.69	0.54	0.72	0.60	0.73	0.59	0.91	0.83	0.90	0.81	1.01	0.85
\overline{x}/X	1.0069	1.0054	1.0072	1.0060	1.0073	1.0059	1.0091	1.0083	1.0090	1.0081	1.0101	1.0085
Dimension W	(Width), m	m										
$\frac{-}{x}$	20.57	20.57	20.60	20.73	20.59	20.62	20.99	20.96	20.99	20.94	21.13	21.15
S	0.06	0.10	0.08	0.10	0.13	0.08	0.06	0.13	0.09	0.06	0.11	0.09
RSD, %	0.31	0.46	0.38	0.47	0.65	0.40	0.27	0.62	0.44	0.28	0.54	0.43
Δx	0.57	0.57	0.59	0.73	0.59	0.62	0.98	0.96	0.98	0.93	1.13	1.15
$\Delta x/X$, %	0.38	0.38	0.40	0.49	0.39	0.41	0.66	0.64	0.66	0.62	0.76	0.77
\overline{x}/X	1.0283	1.0283	1.0298	1.0366	1.0294	1.0308	1.0493	1.0479	1.0493	1.0468	1.0567	1.0575
Dimension H (Height), m	m										
$\frac{-}{x}$	4.22	4.19	4.28	4.25	4.38	4.40	4.21	4.15	4.18	4.18	4.43	4.48
S	0.06	0.04	0.05	0.05	0.08	0.07	0.08	0.03	0.03	0.01	0.09	0.08
RSD, %	1.44	0.85	1.05	1.06	1.75	1.49	1.87	0.68	0.78	0.28	1.94	1.76
$\Delta \mathbf{x}$	0.22	0.19	0.27	0.25	0.38	0.40	0.21	0.15	0.18	0.18	0.43	0.48
$\Delta x/X$, %	0.15	0.12	0.18	0.17	0.25	0.26	0.14	0.10	0.12	0.12	0.28	0.32
\overline{x}/X	1.0550	1.0463	1.0688	1.0633	1.0954	1.0992	1.0521	1.0367	1.0458	1.0446	1.1067	1.1192
Dimension W1	(Neck wid	th), mm										
$\frac{-}{x}$	10.53	10.48	10.61	10.61	10.60	10.61	11.08	11.04	11.02	10.86	11.27	11.27
S	0.04	0.03	0.06	0.11	0.10	0.07	0.11	0.17	0.05	0.04	0.16	0.12
RSD, %	0.41	0.31	0.59	1.04	0.91	0.65	1.01	1.58	0.42	0.40	1.38	1.03
Δx	0.53	0.48	0.61	0.61	0.60	0.61	1.08	1.04	1.02	0.86	1.27	1.27
$\Delta x/X$, %	0.35	0.32	0.41	0.40	0.40	0.41	0.72	0.69	0.68	0.57	0.84	0.85
\overline{x}/X	1.0532	1.0475	1.0608	1.0607	1.0603	1.0613	1.1075	1.1035	1.1015	1.0857	1.1267	1.1270

Table 1. Summary of base set measurements for L, W, H and W1 dimensions

For the particular combination in scaled set of samples, anisotropic scale factor (AS) of particular dimension is a multiplicative inverse of the matching relative deviation from the same combination in the base set:

$$AS_D = \left(\frac{\overline{x}}{\overline{X}}\right)^{-1} \tag{2},$$

where index D denotes anisotropic scaling direction and could be either X, Y or Z.

For example, the anisotropic scale factor in building direction X of the experimental combination 1XW is:

$$AS_{X}(1XW) = \left(\frac{151.03}{150}\right)^{-1} = 1.0069^{-1} = 0.99318$$
(3).

But for the combination 1YW, in which the largest dimension of the sample (L) is oriented in building direction Y, the width (W) of the sample is oriented along building direction X, therefore:

$$AS_{X}(IYW) = \left(\frac{20.57}{20}\right)^{-1} = 1.028\dot{3}^{-1} = 0.97245 \quad (4).$$

Table 2. Anisotropic scale factors

Such an exchange of length and width values in the calculation of anisotropic scale factor must be done for every experiment combination with samples oriented along the building direction Y.

The complete list of applied anisotropic scale factors for each experiment combination of scaled set is given in Table 2.

Label	1XW	1YW	1XE	1YE	1XC	1YC	2XW	2YW	2XE	2YE	2XC	2YC
AS_X	0.9932	0.9725	0.9928	0.9647	0.9928	0.9701	0.9910	0.9543	0.9910	0.9553	0.9900	0.9450
AS_Y	0.9725	0.9946	0.9711	0.9940	0.9714	0.9941	0.9531	0.9918	0.9531	0.9919	0.9464	0.9923
AS_Z	0.9479	0.9558	0.9357	0.9404	0.9129	0.9098	0.9505	0.9646	0.9562	0.9573	0.9036	0.8882

Appropriate anisotropic scale factors were entered in ZPrint software regarding printed

samples and combination. Table 3. shows measurement results of scaled set samples.

Table 3. Summary of scaled set measurements

Experiment Label	1XW	1YW	1XE	1YE	1XC	1YC	2XW	2YW	2XE	2YE	2XC	2YC
Dimension L (Length), mm												
\overline{X}	150.09	150.11	149.92	149.96	150.19	150.26	150.05	150.26	150.13	150.22	150.23	150.20
S	0.12	0.08	0.05	0.04	0.07	0.20	0.11	0.07	0.10	0.03	0.19	0.13
RSD, %	0.08	0.05	0.04	0.03	0.05	0.14	0.08	0.04	0.07	0.02	0.13	0.09
Δx	0.09	0.11	-0.08	-0.04	0.19	0.26	0.04	0.25	0.13	0.22	0.23	0.20
\overline{x}/X	1.0006	1.0007	0.9994	0.9997	1.0013	1.0017	1.0003	1.0017	1.0009	1.0015	1.0016	1.0013
Dimension W	(Width), 1	mm										
\overline{X}	20.17	20.14	19.96	19.90	20.27	20.32	20.16	20.10	20.10	20.19	20.15	20.00
S	0.03	0.07	0.02	0.05	0.11	0.09	0.13	0.05	0.08	0.08	0.07	0.23
RSD, %	0.14	0.36	0.12	0.26	0.54	0.42	0.63	0.26	0.40	0.40	0.35	1.13
Δx	0.17	0.14	-0.04	-0.11	0.27	0.32	0.16	0.10	0.10	0.19	0.15	0.00
\overline{x}/X	1.0083	1.0068	0.9980	0.9948	1.0136	1.0158	1.0082	1.0052	1.0049	1.0093	1.0074	0.9998
Dimension H	(Height), r	nm										
\overline{X}	3.98	4.09	4.03	4.02	4.09	4.11	4.02	4.10	4.02	4.13	4.04	4.10
S	0.02	0.03	0.03	0.03	0.03	0.03	0.02	0.01	0.03	0.02	0.08	0.05
RSD, %	0.59	0.74	0.78	0.75	0.77	0.68	0.53	0.22	0.69	0.55	2.03	1.10
Δx	-0.02	0.09	0.03	0.01	0.09	0.11	0.02	0.10	0.02	0.11	0.04	0.10
\overline{x}/X	0.9958	1.0217	1.0075	1.0038	1.0225	1.0283	1.0046	1.0250	1.0054	1.0317	1.0096	1.0242
Dimension W	1 (Neck wi	idth), mm										
\overline{X}	10.51	10.39	10.37	10.26	10.64	10.47	10.75	10.68	10.64	10.42	10.88	10.89
S	0.14	0.04	0.06	0.05	0.11	0.11	0.10	0.12	0.06	0.02	0.06	0.10
RSD, %	1.31	0.36	0.59	0.44	1.03	1.07	0.93	1.09	0.55	0.19	0.57	0.95
Δx	0.51	0.39	0.37	0.26	0.64	0.47	0.75	0.68	0.64	0.42	0.88	0.89
\overline{x}/X	1.0513	1.0392	1.0370	1.0262	1.0640	1.0468	1.0753	1.0683	1.0635	1.0417	1.0877	1.0887

4 ANALYSIS AND DISCUSSION

In order to determine whether the measured results are significantly affected by surface roughness, standard deviations of measured dimensions are compared with the average roughness of samples. The level of roughness for 3D printed samples is N10, and respectively average roughness is 12.5 μ m [2]. Standard deviations for both base and scaled set are several times higher than average roughness, therefore surface roughness did not have any significant influence on the measured results.

The overall effect of anisotropic scaling could be assessed if the results of both base and scaled set are compared (Fig. 7). A comparison of average errors between the sets reveals a considerable reduction of errors for all scaled dimensions in the scaled set of samples. A smaller error reduction is obtained for the neck width (W1) which is not scaled by own scale factor but indirectly by the anisotropic scale factor of width (W).



Fig. 7. Average errors between sets

A comparison of average relative standard deviations between the sets shows that there is no significant difference in deviations before and after scaling, except for the height of samples (Fig. 8).

The analysis of relative deviations for each experiment factor in the scaled set, reveals active relations within the considered factor. For dimensions: largest dimension – length (L) - has the lowest average relative deviation (1.000889), while the smallest – height (H) - has the highest relative deviation (1.015000). When comparing infiltrants: the samples infiltrated with the epoxy resin have the lowest average relative deviation (1.004735), while the samples infiltrated with the cyanoacrylate have the highest values (1.010592).

The samples oriented along the building direction X have smaller deviations (1.004993) than those oriented along Y direction (1.009605). Finally, the samples printed with layer thickness 0,1 mm do not show any significant difference in relative deviation compared to those printed with layer thickness 0.0875 mm: 1.006682 to 1.007915.



Considering particular combination and dimension, combinations 1YE and 2XW had the least average relative deviation of specimen's length (1.0003). For width, combinations 2YC (0.9998) had minimum and closest to it was 1XE (0.998), while for height it was combination 1YE (1.0038).



Fig. 9. Sorted average relative errors of scaled set combinations

A comparison of average relative errors for combinations of the scaled set reveals that the combination with the lowest relative error is 1YE (Fig. 9). The combination 1YC had the highest relative error. It is interesting to notice that the most accurate combination and as the second most inaccurate - 2YE share the same infiltrant type and building orientation. Also confusing could be the fact that the most accurate combination is printed with higher layer thickness as the second most inaccurate combination. Most likely the reason for this is that the accuracy of combination 2YE is significantly influenced by the operator's skills either during the cleaning of samples or while applying the epoxy resin. Relative errors of other combinations are equally distributed regarding the particular factor.

In order to understand completely the effect of factors on relative deviation and printing accuracy, we performed a factorial analysis of variance (ANOVA). Since there is a significant difference in relative deviations between the measured dimensions, a dimension is considered as an additional factor for ANOVA calculations. It is considered with 3 levels – length, width and height. Other factors and their levels are considered according to the experiment design. ANOVA of relative deviations is presented in Table 4. The abbreviations used for factors are: D - dimension: LT - laver thickness: 0 orientation; I - infiltrant.

The ANOVA calculation reveals that the dimension of printed sample was the most influential factor on the relative deviation. Beside the dimension, sample orientation and infiltrant type also showed a significant influence on the relative deviation. However, the analysis did not prove any influence of layer thickness variations.

Despite the lack of influence when considered alone, the layer thickness variations had a significant influence in the interaction with the infiltrant type. The interaction of dimension and sample orientation is also strongly emphasized according to the analysis of variance. Other interactions of factors have either smaller or insignificant influence when compared to the critical value of Fischer's distribution with the probability of 0.95 and respective degrees of freedom (1; 180).

To avoid an oversight of the possible significant interaction of particular variants within a particular experiment combination, we also performed a differential analysis of the influence of particular factor variants within the other factor.

Factors	Sum of squares	Degrees of freedom	Mean square	Variance	Probability	F(0.95)
D	0.0073	2	0.0037	103	0.0000	3.89
LT	0.0001	1	0.0001	2	0.1306	3.89
0	0.0011	1	0.0011	32	0.0000	3.89
Ι	0.0013	2	0.0006	18	0.0000	3.89
D*LT	0.0001	2	0.0001	2	0.1335	3.89
D*O	0.0029	2	0.0014	40	0.0000	3.89
LT*O	0.0002	1	0.0002	4	0.0373	3.89
D*I	0.0008	4	0.0002	6	0.0003	3.89
LT*I	0.0020	2	0.0010	29	0.0000	3.89
O*I	0.0002	2	0.0001	3	0.0609	3.89
D*LT*O	0.0004	2	0.0002	6	0.0040	3.89
D*LT*I	0.0009	4	0.0002	6	0.0002	3.89
D*O*I	0.0005	4	0.0001	3	0.0145	3.89
LT*O*I	0.0006	2	0.0003	8	0.0004	3.89
D*LT*O*I	0.0006	4	0.0002	4	0.0027	3.89
Error	0.0064	180	0.0000			

Table 4. Anova of relative deviations in scaled set

For that purpose, a thorough reading of results for each combination and reading of appropriate multifactor charts were performed (Fig. 10 and Fig. 11).

Differential analysis reveals one particular effect in the interaction of dimensions and building orientation. Although the interaction of those factors appears significant for the experiment, it is mainly caused by big variances of height. Other dimensions did not show any significant variances regarding the building orientation.



Fig. 10. Average relative deviations for layer thickness 1



Fig.11. Average relative deviations for layer thickness 2

Multifactor charts show that relative deviations of samples printed with smaller layer thickness 2, are grouped in narrower clusters than relative deviations of larger layer thickness 1. Therefore, smaller layer thickness 2 provided a better overall precision of samples.

5 CONCLUSION

The obtained experiment results clearly show that the scaling of 3D printed samples provides a better dimensional precision of samples. Furthermore, the significant differences in relative errors and relative standard deviations between measured dimensions prove the adequacy of anisotropic scaling for samples with significant differences between the main dimensions. A limitation of the anisotropic scaling exists for dimensions that are not scaled directly with their own scale factor, as it was presented in this paper for dimension of neck width (W1).

Although a comparison of average relative errors reveals the combinations with minor errors, from the factorial analysis of variance it can be concluded that an appropriate combination of 3DP processing factors should be selected regarding the significant dimension of the printed sample. For the accuracy of the largest dimension, i.e. length, the sample should be printed with default layer thickness 0,1 mm, oriented toward building direction X and infiltrated with the epoxy resin. Since there is a significant interaction of factors, i.e. layer thickness and the infiltrant type, from differential analysis it could be concluded that a combination of the default layer thickness and infiltration with the wax is also the appropriate selection for both considered building orientations. The accuracy of the medium dimension (width) could be improved if the printed sample is infiltrated with the epoxy resin, while other processing factors did not show any significant influence. The accuracy of the smallest dimension (height) could be improved if the printed sample is infiltrated with wax or epoxy resin and the sample is oriented towards building direction X.

From the overall analysis the following could be concluded for the particular processing factor:

- The layer thickness: smaller layer thickness provides better overall precision of samples.
- The infiltrants: the most accurate are samples infiltrated with epoxy resin and

the least accurate are samples infiltrated with cyanoacrylate.

• The building orientation: the samples oriented towards direction X are more accurate than samples oriented towards direction Y.

Further researches could also be performed to improve the accuracy of the considered 3DP system like the experimental measurement of test sample behaviour and a finite-element based simulation using a non-linear visco-elastic constitutive model similar to the one published in [13].

Although it is possible that some of the presented conclusions are valid for similar machines or even rapid prototyping techniques, all conclusions should be considered only for the selected 3D printer and selected materials. New materials and new equipment for 3D printing are developed constantly and could demand new analyses.

6 ACKNOWLEDGMENT

The authors wish to thank Professor Pero Raos and system engineer Miroslav Mazurek for supporting the work on this paper; and also Vedran Galeta for the assistance in experiments; and Professor Željka Rosandić for proofreading the manuscript. The work presented in this paper was financially supported by the Ministry of Science, Education and Sports, Republic of Croatia through the scientific project No. 152-1521473-3111.

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Vsebina

Strojniški vestnik - Journal of Mechanical Engineering letnik 54, (2008), številka 10 Ljubljana, oktober 2008 **ISSN 0039-2480**

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Preprost zančni sistem za aktivno dušenje hrupa v kanalu

Jan Černetič^{*} - Jurij Prezelj - Mirko Čudina Univerza v Ljubljani, Fakulteta za strojništvo, Slovenija

Sistem za aktivno dušenje hrupa ponavadi vsebuje elektronske filtre. Vendar za učinkovito delovanje ti filtri niso vedno nujno potrebni. V tem primeru odpade potreba po elektronskem krmilniku in ustrezni programski opremi, kar naredi sistem veliko bolj preprost in cenejši. V tem prispevku je predstavljen zančni sistem za aktivno dušenje hrupa v eksperimentalnem prezračevalnem kanalu, brez uporabe elektronskih filtrov. Narejena je bila simulacija, s katero se je ugotavljalo učinkovitost preprostega analognega zančnega sistema za aktivno dušenje hrupa. Za oceno najvišje možne ravni dušenja je bila uporabljena prenosna funkcija celotnega dejanskega sistema. Opravljeni so bili preizkusi, kjer je bila izmerjena raven dušenja oz. učinkovitost danega sistema, le-ta pa je bila nato primerjana s teoretičnimi vrednostmi iz simulacije. Izkazalo se je, da je v določenem frekvenčnem območju takšen sistem dovolj učinkovit zi uporabo v nekaterih prezračevalnih kanalih.

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Ključne besede: aktivno dušenje hrupa, povratne zanke, ADH simuliranje, fazni zamik



Slika 2. Shema zančne strukture ADH

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Dinamična analiza mehanizmov za dviganje bremen

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Prispevek obravnava problem dinamičnega obnašanja mehanizmov za dviganje bremen (kot so dvigala). Pri večjih dvižnih višinah se uporabljajo naprave z veliko hitrostjo, ki skrajšajo čas krogov in povečajo zmogljivost. V tem primeru ni mogoče uporabiti standardnega postopka dinamične analize. V prispevku je predlagan postopek za določevanje ustreznega dinamičnega modela in pripadajočih enačb. Postopek omogoča analizo vplivnih veličin kot so spreminjanje proste dolžine vrvi, spodrsavanje elastične cevi na bobnu ali škripcu in dušenje zaradi notranjega trenja v vrvi. © 2008 Strojniški vestnik. Vse pravice pridržane.

Ključne besede: dvigala, škripci, dinamični modeli, dolžina vrvi, stabilnost gibanja



Slika 1. Pogonski mehanizem dvigala

Vpliv zaščitne mreže na integralne ter lokalne značilke aksialnega ventilatorja

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V prispevku je predstavljen vpliv zaščitne mreže (ZM) na integralno karakteristiko (IK) in lokalne aerodinamske lastnosti izstopnega toka aksialnega ventilatorja (AV). Študija vpliva ZM je bila izvedena z meritvami in numeričnim modeliranjem na integralnem in lokalnem nivoju. Meritve IK so bile izvedene po standardu ISO 5801, meritve lokalnega hitrostnega polja pa so bile izvedene s petluknjično sondo. Ob enakih delovnih pogojih ventilatorja so bili izvedeni tudi numerični izračuni lokalnih značilk stroja. S primerjavo rezultatov pa je bila izvedena analiza vplivnosti mrežice na učinkovitost ventilatorja in analiza smotrnosti CFD-modeliranja tokovnih razmer na mrežici.

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Ključne besede: zaščitni ventilatorji, zaščitne mreže, integralne značilke, hitrostna polja



Slika 1. Izvedba ventilatorja z zaščitno mrežo

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Vpliv koncentracije odpadnega olja v vodi na učinkovitost postopka prezračevanja pri čiščenju rafinerijskih odpadnih vod

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Postopek, sistem in postopek prezračevanja za biološko čiščenje odpadnih vod z aktivnim blatom v bio-prezračevalnih posodah se izbirajo na osnovi parametrov pretoka, sestave odpadnih vod in zahtevanih lastnosti prečiščene vode. Izbira sistema za prezračevanje je zelo zahteven problem, saj se mora za najboljši učinek čiščenja količina kisika, dodanega odpadnim vodam, ujemati s porabo kisika.

V prispevku so predstavljeni rezultati eksperimentalne analize prezračevanja vode, onesnažene z odpadnim motornim oljem, pri različnih pretokih dodanega zraka. Obravnavani so bili postopki prezračevanja čiste vode in vode z različnimi vsebnostmi olja (masna koncentracija od 5 do 10 mg/l). Namen eksperimentalne analize je bil analizirati tehnične kazalnike postopka prezračevanja v odvisnosti s termodinamičnimi in kinetičnimi parametri, povezanimi s prisotnostjo olja v vodi, za doseganje bolj učinkovitega čiščenja rafinerijskih odpadnih vod.

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Ključne besede: prezračevanje, rafinerijske odpadne vode, prenos kisika, relaksacijski postopki



Slika 1. Shema postavitve eksperimenta

1 – nizkotlačni kompresor; 2 – zaporni ventil za dovod zraka; 3 – razbremenilni ventil; 4 – regulator pretoka zraka; 5 – zaslonka za merjenje pretoka zraka, 6 – kolona s priključki in ogrodjem, 7 – zračni dovodnik z diskasto membrano; 8 – dovod vode, 9 – vzorčevalni priključek

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Trdnost tlačnih posod z elipsoidnim pokrovom

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Predlagana je metoda za analizo napetosti v valjastih tlačnih posodah z elipsoidnim pokrovom na osnovi teorije osno-simetričnih lupin. Izhodišče so bile aproksimativne rešitve sistema diferencialnih enačb, iz katerih so bili pridobljeni matematični izrazi za določanje notranjih sil, momentov in pomikov v stenah posod.

Končni uporabni izrazi so bili pridobljeni z združitvijo teorije membran in momentov ter z nastavitvijo in reševanjem enačb robnih pogojev. Diagrami, ki prikazujejo porazdelitev notranjih sil in momentov, so v brezrazsežni obliki, ki omogoča njihovo uporabo ne glede na izmere in obremenitve. Ti izrazi so bili uporabljeni za razvoj metod za preizkušanje trdnosti tlačnih posod z elipsoidnim pokrovom v fazi konstrukcije. Prikazana je uporaba metode na izbranem numeričnem primeru, za namene izračunavanja pa je bil razvit poseben računalniški program.

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Ključne besede: tlačne posode, elipsoidni pokrovi, notranje sile, analitične rešitve



Slika 4. Porazdelitev radialnih sil v steni posode, določena s teorijo momentov

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Numerična in eksperimentalna analiza vpliva povesa gredi na dinamiko stika rotorja in statorja

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Poves gredi je stvaren problem, ki se pojavlja zlasti pri vitkih rotorjih. Prispevek obravnava vpliv povesa gredi na dinamiko stika rotorja in statorja. Rotor je v ta namen opisan kot enostaven Jeffcottov model, stator pa kot elastično obešen togi obroč. Za preskušanje numeričnega modela je bil poleg običajne analize zaustavljanja analiziran tudi izklop v nuji po nenadnem povečanju neuravnoteženosti rotorja. Opravljena je bila numerična integracija po metodi Runge-Kutta četrtega reda. Analizirana sta bila dva različna modela normalnih sil za interakcijo rotorja in statorja. Tako pri linearnih in nelinearnih analiziranih parametrih (Hunt in Crossley) so modeli normalnih sil dali podoben odgovor rotorja in statorja. Za preverjanje nekaterih rezultatov in fino nastavitev numeričnega modula je bila opravljena eksperimentalna raziskava na preskusnem sistemu.

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Ključne besede: stik rotor- stator, poves gredi, numerično simuliranje, vibracije rotorjev, meritve



Slika 1. Razširjeni Jeffcottov model rotorja

MINLP optimiranje jeklene industrijske hale

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V prispevku je predstavljeno optimiranje topologije in standardnih prerezov jeklene industrijske hale, izdelane iz jeklenih standardnih vroče valjanih I prerezov. Konstrukcijo sestavljajo glavni okvirji na katere so pritrjene lege. Optimiranje je izvedeno z mešanim celoštevilskim nelinearnim programiranjem (MINLP). Pri MINLP poteka diskretno optimiranje topologije in standardnih izmer sočasno z računom zveznih parametrov. Za reševanje nekonveksnega, nelinearnega in kombiniranega diskretno-zveznega optimizacijskega problema jeklenih hal je uporabljen modificirani algoritem zunanje aproksimacije s sprostitvijo enačb (Modified OA/ER). Poleg optimalne mase konstrukcije so še izračunani optimalna topologija z optimalnim številom glavnih okvirjev in leg ter standardni prerezi vseh konstrukcijskih elementov. Prispevek vsebuje teoretični opis problema in praktični primer z rezultati optimiranja. © 2008 Strojniški vestnik. Vse pravice pridržane.

Ključne besede: industrijske hale, optimiranje topologije, optimiranje prerezov, nelinearno programiranje, MINLP



Slika 1: Industrijska hala

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Geometrijska natančnost z 2-D tiskalnim modelom

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Za izboljšanje natančnosti trirazsežnega tiskanja (3DT) smo izvedli vrsto eksperimentov na dvorazsežnih (2-D) tiskanih modelih. Pri vsakem eksperimentu je obravnava kombinacija treh glavnih dejavnikov: debeline plasti, usmeritve nanosa in vrste veziva. Rezultati eksperimentov dajejo ustrezen faktor anizotropnega skaliranja in dokazujejo primernost anizotropnega skaliranja. Analiza rezultatov kaže, da je boljšo natančnost vzorca mogoče doseči z manjšo debelino plasti, infiltracijo z epoksi-smolo in usmeritvijo vzorca v smeri nanosa po osi X.

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Ključne besede: hitra izdelava prototipov, trirazsežno tiskanje, geometrijska natančnost



Slika 3. 3D-tiskanje preskusnih vzorcev

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Strojniški vestnik - Journal of Mechanical Engineering 54(2008)10, SI 124 Navodila avtorjem

Navodila avtorjem

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Članki naj bodo kratki in naj obsegajo približno 8-12 strani.

Člankom so lahko priložene tudi dodatne računalniške simulacije ali predstavitve, pripravljene v primerni obliki, ki bodo bralcem dostopne na spletni strani revije.

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- Literatura, ki mora biti v besedilu oštevilčena zaporedno in označena z oglatimi oklepaji [1] ter na koncu članka zbrana v seznamu literature.

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Seznam literature

Vsa literatura mora biti navedena v seznamu na koncu članka v prikazani obliki po vrsti za revije, zbornike in knjige:

- [1] Wagner, A., Bajsić, I., Fajdiga, M. Measurement of the surface-temperature field in a fog lamp using resistance-based temperature detectors. Strojniški vestnik – Journal of Mechanical Engineering, February 2004, vol. 50, no. 2, p. 72-79.
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International Thermal Science Seminar II, Bled, June 13.-16., 2004.

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